

## Homework Assignment #6 — Due Tuesday, February 21

Textbook problems: Ch. 10: 10.2, 10.3, 10.7, 10.10

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- 10.2 Electromagnetic radiation with elliptic polarization, described (in the notation of Section 7.2) by the polarization vector,

$$\vec{\epsilon} = \frac{1}{\sqrt{1+r^2}}(\vec{\epsilon}_+ + r e^{i\alpha} \vec{\epsilon}_-)$$

is scattered by a perfectly conducting sphere of radius  $a$ . Generalize the amplitude in the scattering cross section (10.71), which applies for  $r = 0$  or  $r = \infty$ , and calculate the cross section for scattering in the long-wavelength limit. Show that

$$\frac{d\sigma}{d\Omega} = k^4 a^6 \left[ \frac{5}{8}(1 + \cos^2 \theta) - \cos \theta - \frac{3}{4} \left( \frac{r}{1+r^2} \right) \sin^2 \theta \cos(2\phi - \alpha) \right]$$

Compare with Problem 10.1.

- 10.3 A solid uniform sphere of radius  $R$  and conductivity  $\sigma$  acts as a scatterer of a plane-wave beam of unpolarized radiation of frequency  $\omega$ , with  $\omega R/c \ll 1$ . The conductivity is large enough that the skin depth  $\delta$  is small compared to  $R$ .
- Justify and use a magnetostatic scalar potential to determine the magnetic field around the sphere, assuming the conductivity is infinite. (Remember that  $\omega \neq 0$ .)
  - Use the technique of Section 8.1 to determine the absorption cross section of the sphere. Show that it varies as  $(\omega)^{1/2}$  provided  $\sigma$  is independent of frequency.
- 10.7 Discuss the scattering of a plane wave of electromagnetic radiation by a nonpermeable, dielectric sphere of radius  $a$  and dielectric constant  $\epsilon_r$ .
- By finding the fields inside the sphere and matching to the incident plus scattered wave outside the sphere, determine without any restriction on  $ka$  the multipole coefficients in the scattered wave. Define suitable phase shifts for the problem.
  - Consider the long-wavelength limit ( $ka \ll 1$ ) and determine explicitly the differential and total scattering cross sections. Compare your results with those of Section 10.1.B.
  - In the limit  $\epsilon_r \rightarrow \infty$  compare your results to those for the perfectly conducting sphere.

10.10 The aperture or apertures in a perfectly conducting plane screen can be viewed as the location of effective sources that produce radiation (the diffracted fields). An aperture whose dimensions are small compared with a wavelength acts as a source of dipole radiation with the contributions of other multipoles being negligible.

- a) Beginning with (10.101) show that the effective electric and magnetic dipole moments can be expressed in terms of integrals of the tangential electric field in the aperture as follows:

$$\vec{p} = \epsilon \hat{n} \int (\vec{x} \cdot \vec{E}_{\text{tan}}) da$$
$$\vec{m} = \frac{2}{i\omega\mu} \int (\hat{n} \times \vec{E}_{\text{tan}}) da$$

where  $\vec{E}_{\text{tan}}$  is the *exact* tangential electric field in the aperture,  $\hat{n}$  is the normal to the plane screen, directed into the region of interest, and the integration is over the area of the openings.

- b) Show that the expression for the magnetic moment can be transformed into

$$\vec{m} = \frac{2}{\mu} \int \vec{x} (\hat{n} \cdot \vec{B}) da$$

Be careful about possible contributions from the edge of the aperture where some components of the fields are singular if the screen is infinitesimally thick.