

## Homework Assignment #3 — Due Thursday, January 24

Textbook problems: Ch. 8: 8.18, 8.20

Ch. 9: 9.1, 9.3

- 8.18 a) From the use of Green's theorem in two dimensions show that the TM and TE modes in a waveguide defined by the boundary-value problems (8.34) and (8.36) are orthogonal in the sense that

$$\int_A E_{z\lambda} E_{z\mu} da = 0 \quad \text{for } \lambda \neq \mu$$

for TM modes, and a corresponding relation for  $H_z$  for TE modes.

- b) Prove that the relations (8.131)–(8.134) form a consistent set of normalization conditions for the fields, including the circumstances when  $\lambda$  is a TM mode and  $\mu$  is a TE mode.
- 8.20 An infinitely long rectangular waveguide has a coaxial line terminating in the short side of the guide with the thin central conductor forming a semicircular loop of radius  $R$  whose center is a height  $h$  above the floor of the guide, as shown in the accompanying cross-sectional view. The half-loop is in the plane  $z = 0$  and its radius  $R$  is sufficiently small that the current can be taken as having a constant value  $I_0$  everywhere on the loop.
- a) Prove that to the extent that the current is constant around the half-loop, the TM modes are not excited. Give a physical explanation of this lack of excitation.
- b) Determine the amplitude for the lowest TE mode in the guide and show that its value is independent of the height  $h$ .
- c) Show that the power radiated in either direction in the lowest TE mode is

$$P = \frac{I_0^2}{16} Z \frac{a}{b} \left( \frac{\pi R}{a} \right)^4$$

where  $Z$  is the wave impedance of the  $\text{TE}_{10}$  mode. Here assume  $R \ll a, b$ .

- 9.1 A common textbook example of a radiating system (see Problem 9.2) is a configuration of charges fixed relative to each other but in rotation. The charge density is obviously a function of time, but it is not in the form of (9.1).

- a) Show that for rotating charges one alternative is to calculate *real* time-dependent multipole moments using  $\rho(\vec{x}, t)$  directly and then compute the multipole moments for a given harmonic frequency with the convention of (9.1) by inspection or Fourier decomposition of the time-dependent moments. Note that care must

be taken when calculating  $q_{lm}(t)$  to form linear combinations that are real before making the connection.

- b) Consider a charge density  $\rho(\vec{x}, t)$  that is periodic in time with period  $T = 2\pi/\omega_0$ . By making a Fourier *series* expansion, show that it can be written as

$$\rho(\vec{x}, t) = \rho_0(\vec{x}) + \sum_{n=1}^{\infty} \Re[2\rho_n(\vec{x})e^{-in\omega_0 t}]$$

where

$$\rho_n(\vec{x}) = \frac{1}{T} \int_0^T \rho(\vec{x}, t) e^{in\omega_0 t} dt$$

This shows explicitly how to establish connection with (9.1).

- c) For a single charge  $q$  rotating about the origin in the  $x$ - $y$  plane in a circle of radius  $R$  at constant angular speed  $\omega_0$ , calculate the  $l = 0$  and  $l = 1$  multipole moments by the methods of parts a and b and compare. In method b express the charge density  $\rho_n(\vec{x})$  in cylindrical coordinates. Are there higher multipoles, for example, quadrupole? At what frequencies?

9.3 Two halves of a spherical metallic shell of radius  $R$  and infinite conductivity are separated by a very small insulating gap. An alternating potential is applied between the two halves of the sphere so that the potentials are  $\pm V \cos \omega t$ . In the long-wavelength limit, find the radiation fields, the angular distribution of radiated power, and the total radiated power from the sphere.