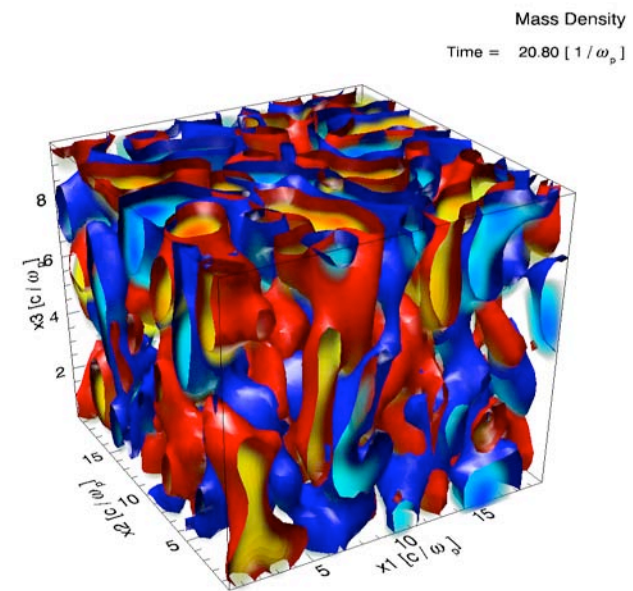


Physical problems (microphysics) in relativistic plasma flows

Luís O. Silva

GoLP/Centro de Física dos Plasmas
Instituto Superior Técnico, Lisbon, Portugal

<http://cfp.ist.utl.pt/golp/>



Collaborations



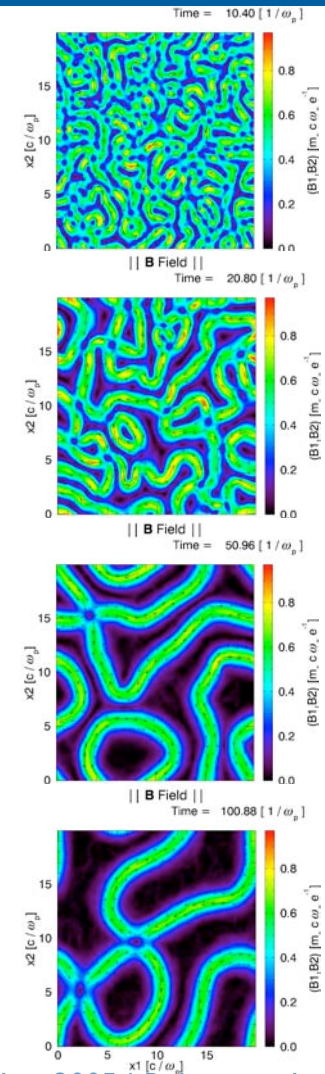
Massimiliano Fiore, Ricardo A. Fonseca,
Michael Marti, Gianfranco Sorasio



Chuang Ren (now at U. Rochester), John Tonge,
Michail Tzoufras, Warren B. Mori



Mikhail Medvedev

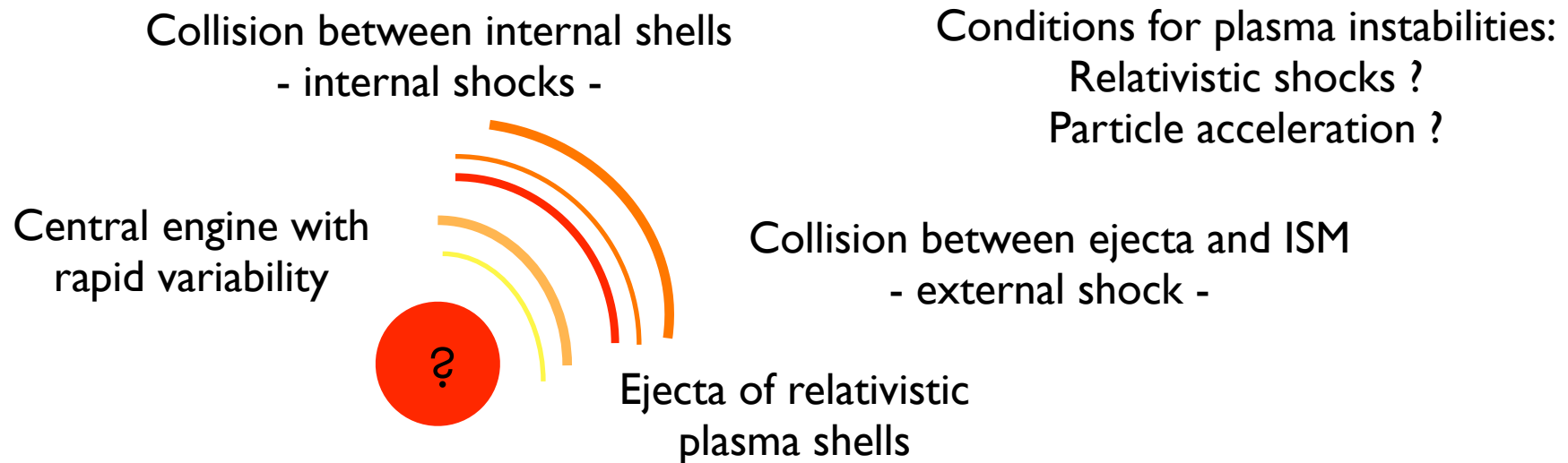


Outline

- Particle-in-cell simulations
 - ≡ probing the microphysics on the time scale of the electron collective dynamics ($f(\Upsilon)$, ϵ_B , ϵ_i , ϵ_e)
- Collisionless plasma instabilities
 - ≡ Weibel and two-stream: electromagnetic beam plasma instability
- Sub-equipartition B-field generation (ϵ_B)
- Particle acceleration ($f(\Upsilon)$)
- Summary

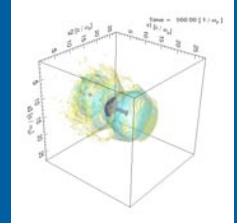
Plasma instabilities in GRBs

- In the fireball model of GRBs, relativistic plasma shells collide/overtake each other (Re'em Sari, yesterday; e.g. Piran, 05 and references therein)
- Synchrotron radiation indicates the presence of sub-equipartition magnetic fields: $\epsilon_{\text{B-field}}/\epsilon_{\text{particles}} \sim 10^{-3}$
- Weibel instability can be the mechanism responsible for B-field generation in this scenario (Medvedev and Loeb, 99)



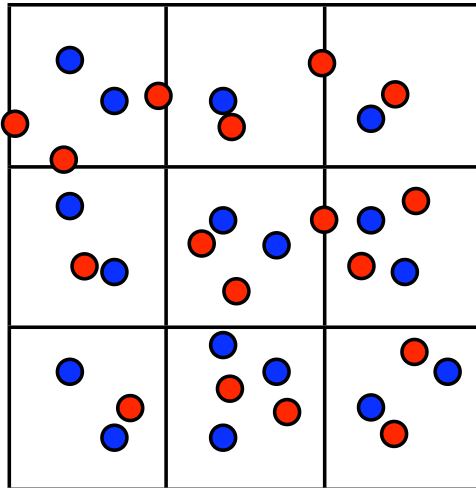


Numerical simulations



Particle-in-cell simulations

Solving Maxwell's equations on a grid with self-consistent charges and currents due to charged particle dynamics



State-of-the-art

$\sim 10^9$ particles
 $\sim (500)^3$ cells

RAM ~ 0.5 TByte

Run time: hours to months

Data/run ~ 1 TByte

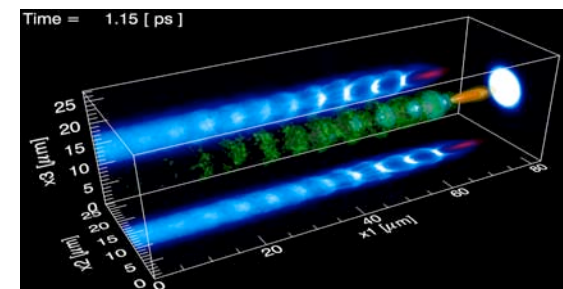
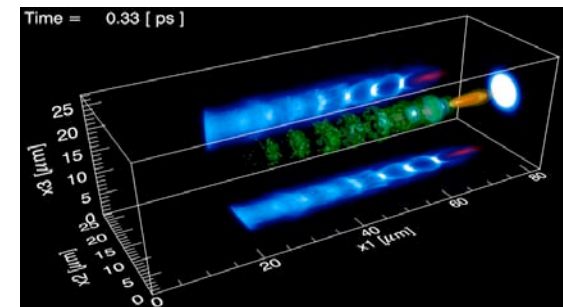
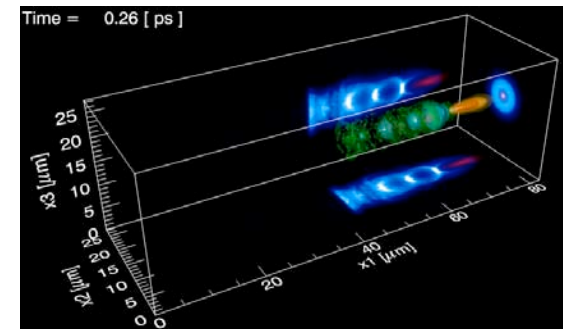
One-to-one simulations of plasma based accelerators & cluster dynamics

Weibel/two stream instability in fast igniton, astrophysics

Particle-in-cell (PIC) - (Dawson, Buneman, 1960's)

Maxwell's equation solved on simulation grid

Particles pushed with Lorentz force



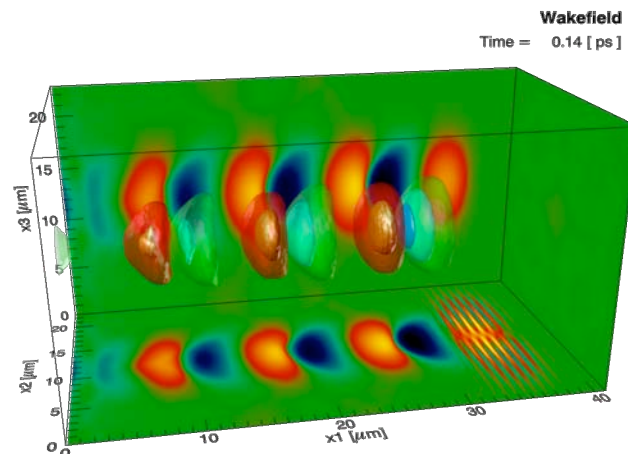
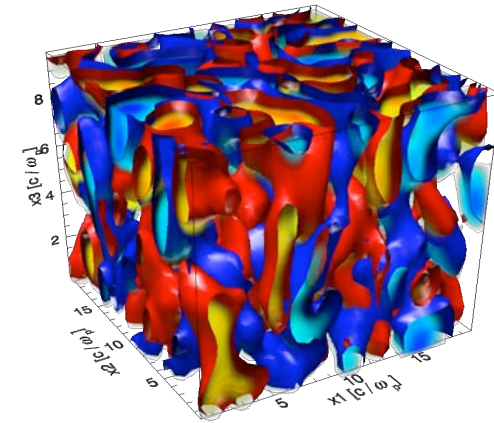
osiris 2.0

osiris
v2.0



osiris.framework

- Massively parallel particle-in-cell code
- Visualization and data analysis infrastructure
- Developed by the *osiris.consortium*
⇒ UCLA + IST + USC

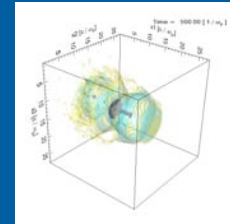


New in version 2.0

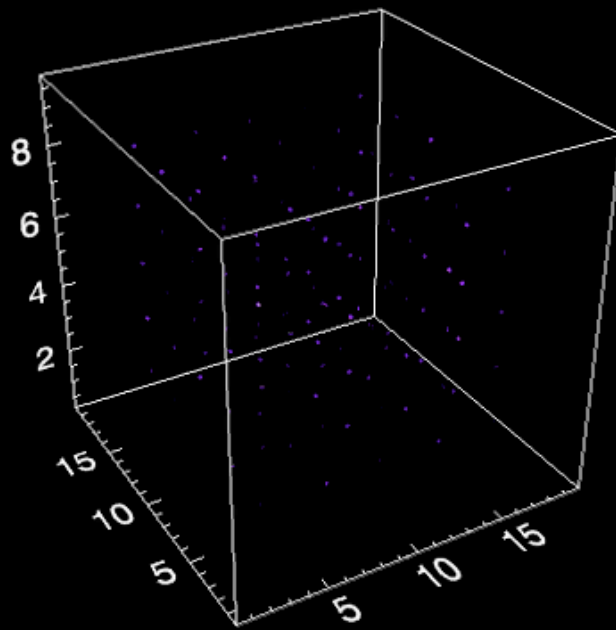
- Bessel beams
- Binary collisions module
- Impact and tunnel ionization (ADK model)
- Dynamic load balancing
- Parallel I/O



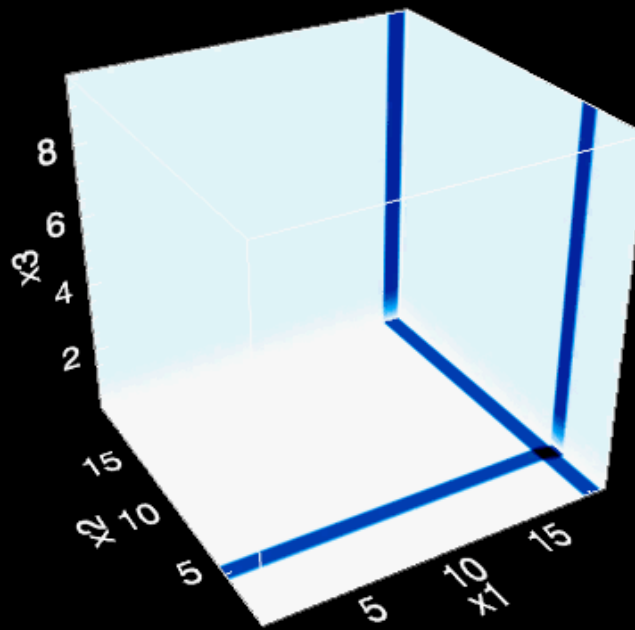
Collisionless plasma instabilities



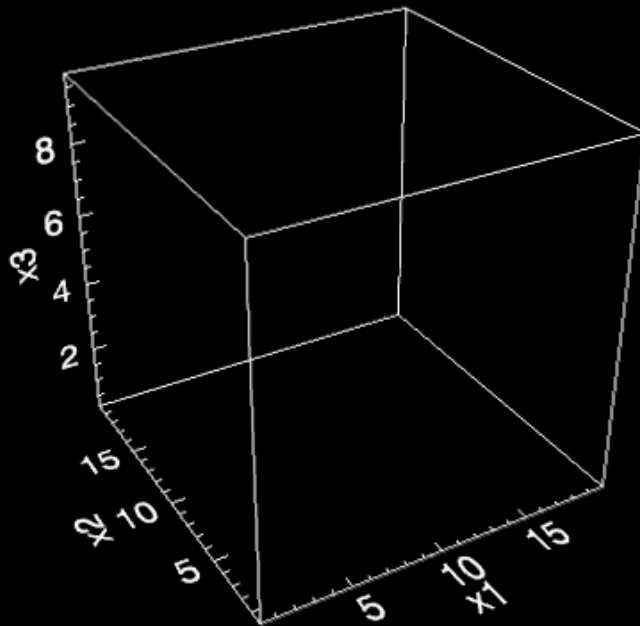
Magnetic Field



Electron Mass Density



EM Energy



3D PIC Simulations of the Weibel Instability

R. A. Fonseca

Run: epc003a.3d
Time = 0.00 [1 / ω_p]

Magnetic Field



Mass Density



EM Energy



Frame: 1/145



Collisionless instabilities of plasma streams

General theory: Watson, Bludman, Rosenbluth (1960)

Two-stream instability (longitudinal mode):

$$\Gamma \sim (n_b/n_0)^{1/3}$$

n_b, v_b - number density/velocity of plasma stream

$$k_{||} \sim \omega_{pe0}/v_b$$

n_0 - number density of background (stationary) plasma

Weibel/filamentation instability for e^-/e^+ (transverse mode):

$$\Gamma \sim (n_b/n_0)^{1/2}$$

$$k_{\perp} \sim \omega_{pe0}/c \quad (T \neq 0)$$

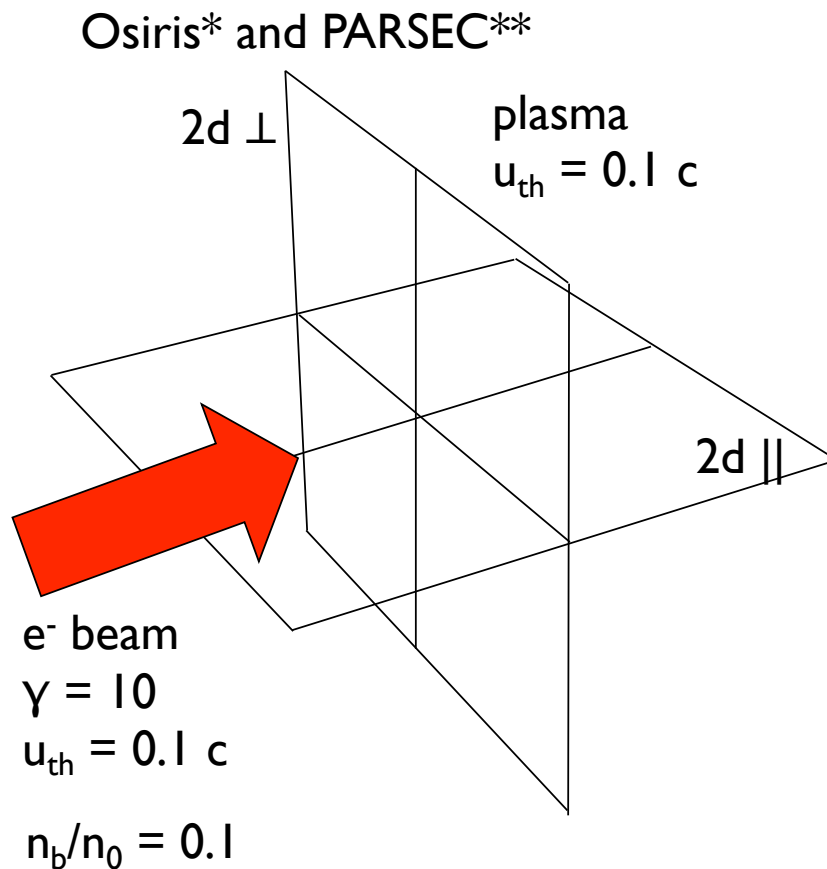
Relevant time and spatial scales

$$L \approx c/\omega_{p0} \approx 5 \text{ km} / \sqrt{n_0 [\text{cm}^{-3}]}$$

$$T \approx 1/\omega_{p0} \approx 20 \mu\text{s} / \sqrt{n_0 [\text{cm}^{-3}]}$$

Multidimensional electromagnetic beam-plasma instability will be a combination of transverse and longitudinal modes

Three dimensional PIC simulations I



Periodic system

Current and charge neutral

Moving ions

- 3D runs

$(128)^3$ cells

$(12.8 c/\omega_{p0})^3$

$\omega_{p0}t_{max} = 500$

16 particles/(species×cell)

- 2D runs (\parallel and \perp)

(512×128) cells

$(51.2 \times 12.8) (c/\omega_{p0})^2$

$\omega_{p0}t_{max} = 500$

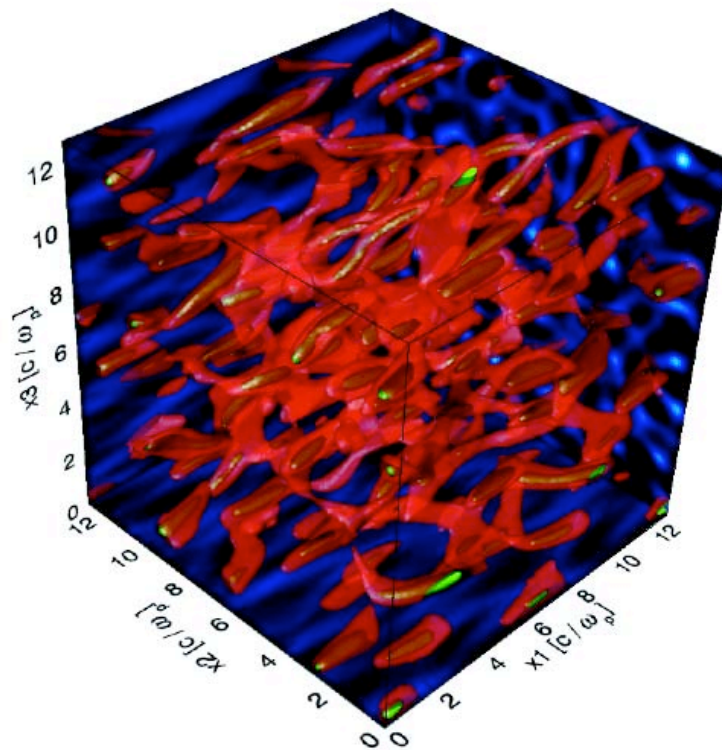
16 particles/(species×cell)

*R.A. Fonseca et al, LCNS 2002;

** J. Tonge, UCLA PhD Thesis 2002

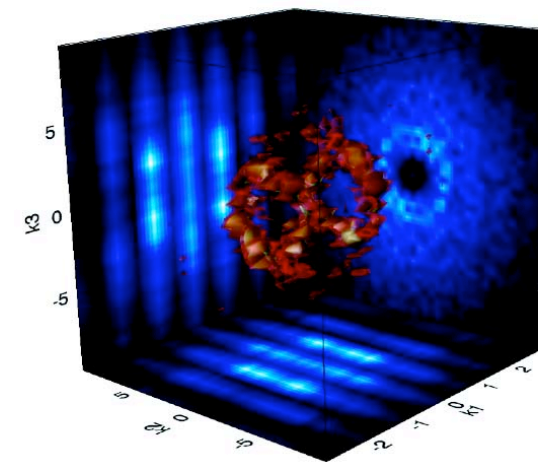
Three dimensional PIC simulations II

Beam density @ $t = 62.4/\omega_{pe0} = 7.5 \Gamma_{\text{weibel}}$



Green = $1.15 n_{e0}$; Red = $0.53 n_{e0}$

FFT



Green = 40%; Red = 27.5%

Fastest growing mode

$$k_{\parallel} \sim 1 \quad k_{\perp} \sim 3$$

$$\theta_{\text{tilt}} \sim 18.4^{\circ}$$

Covariant fluid theory I

$$\frac{\partial T_i^{\alpha\beta}}{\partial x^\beta} = m_i \bar{n}_i F^\alpha \quad \text{Covariant fluid equations} \quad F^\alpha \equiv \text{Lorentz 4 - force}$$

$$T^{\alpha\beta} = p \eta^{\alpha\beta} + (p + \varepsilon) U^\alpha U^\beta \quad \text{Energy-momentum stress tensor}$$

3D representation

$$(\partial_t + \mathbf{v} \cdot \nabla) \mathbf{v} = \frac{qnc^2}{\gamma^2(p + \varepsilon)} \left[\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} \right] - \frac{qn}{\gamma^2(p + \varepsilon)} (\mathbf{v} \cdot \mathbf{E}) \mathbf{v} - \frac{c^2}{\gamma^2(p + \varepsilon)} \left[\nabla p + \frac{\mathbf{v}}{c^2} \partial_t p \right]$$

+ Maxwell's equations

Key issue: relativistic equation of state $p(T,n)$ and $\varepsilon(T,n)$

Covariant fluid theory II

$$U^\alpha = \int u^\alpha \frac{f(u)}{(1 + \mathbf{u} \cdot \mathbf{u})^{1/2}} d\mathbf{u}$$

Fluid proper velocity

$$T^{\alpha\beta} = mc^2 \int u^\alpha u^\beta \frac{f(u)}{(1 + \mathbf{u} \cdot \mathbf{u})^{1/2}} d\mathbf{u}$$

Energy-momentum stress tensor

In the rest frame of the fluid:

$$T^{11} = p \quad T^{00} = \varepsilon$$

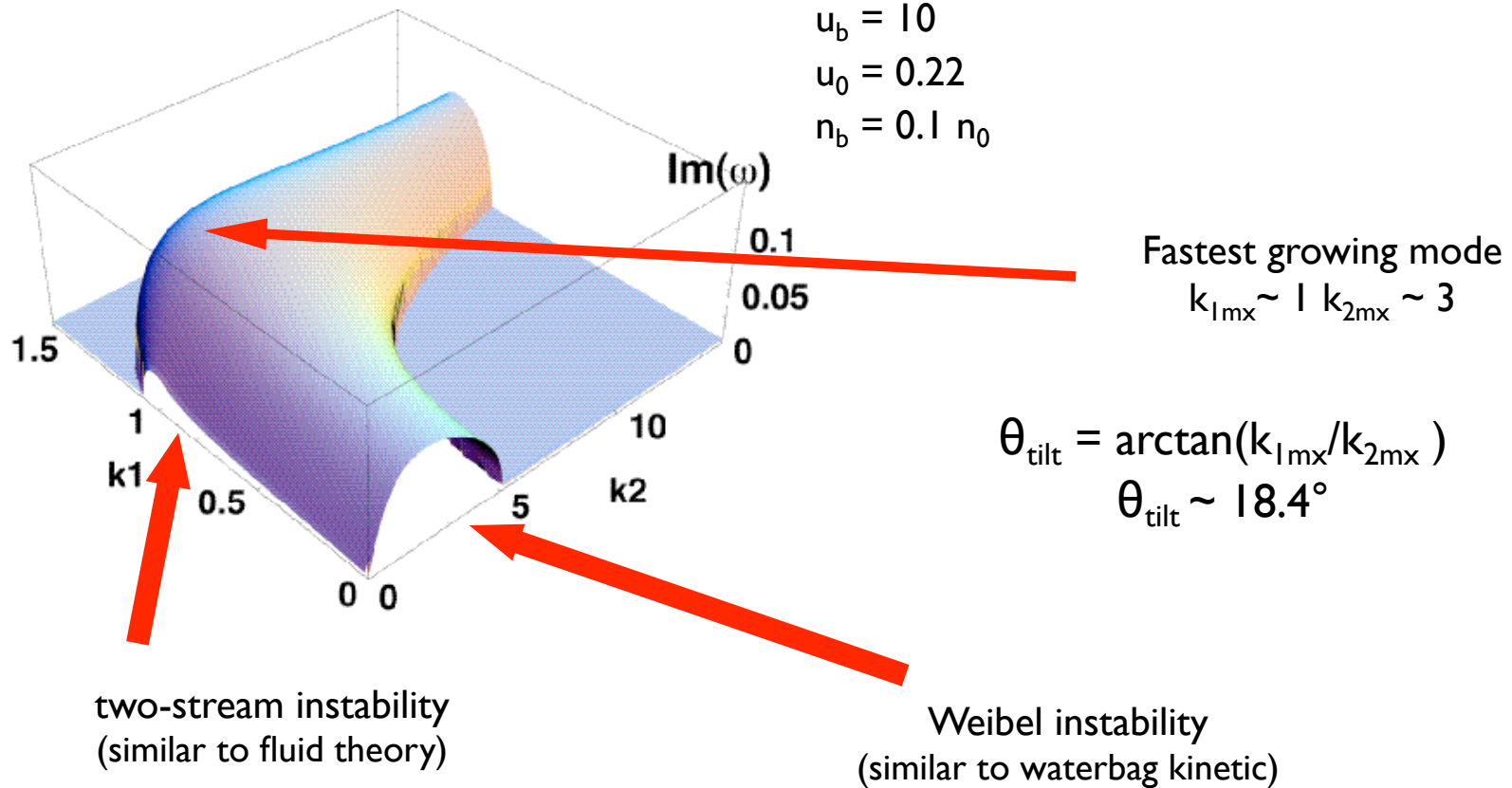
For a waterbag $f(u)$:

$$\frac{p}{n_0 mc^2} = \frac{1}{8u_0^3} \left\{ u_0 \tilde{n} (1 + u_0^2 \tilde{n}^2)^{1/2} (2u_0^2 \tilde{n}^2 - 3) + 3 \operatorname{arcsinh}(u_0 \tilde{n}) \right\} \quad \tilde{n} = (\bar{n}/n_0)^{1/3}$$

$$\varepsilon = n_0 mc^2 \left\{ \tilde{n} + \frac{1}{2} u_0^2 \tilde{n}^{5/3} \right\} - p$$

Tilted filamentation of relativistic electron streams I*

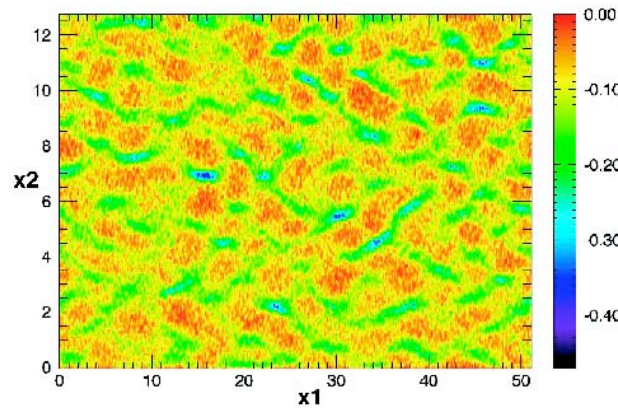
Multidimensional (2D) electromagnetic beam-plasma instability ($T \neq 0$)



*Silva et al, Bull. Am. Phys. Soc., 2001; Bret and Deutsch, Phys. Rev. Lett, 2004

Comparison theory-simulations I

Electron beam spatial structure

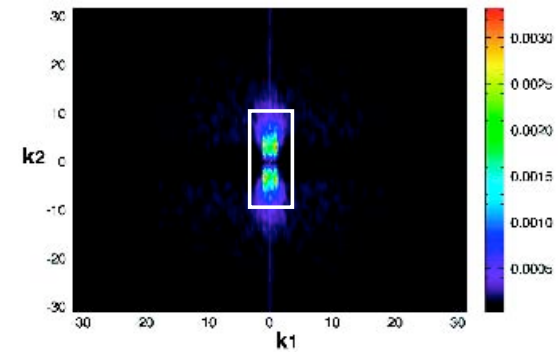


$$k_{1mx} \sim 1, k_{2mx} \sim 3.0$$

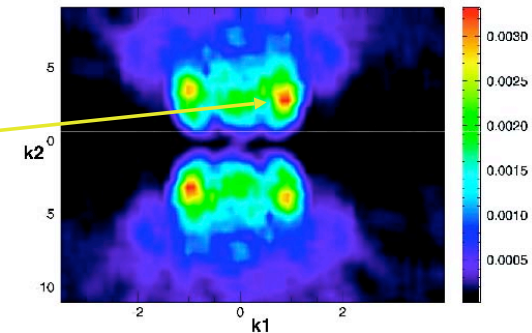
Typical spraying angle $\sim 18^\circ$

2D || runs
(512 × 128) cells
(51.2 × 12.8) $(c/\omega_{p0})^2$
 $\omega_{p0}t_{\max} = 500$
16 particles/(species × cell)

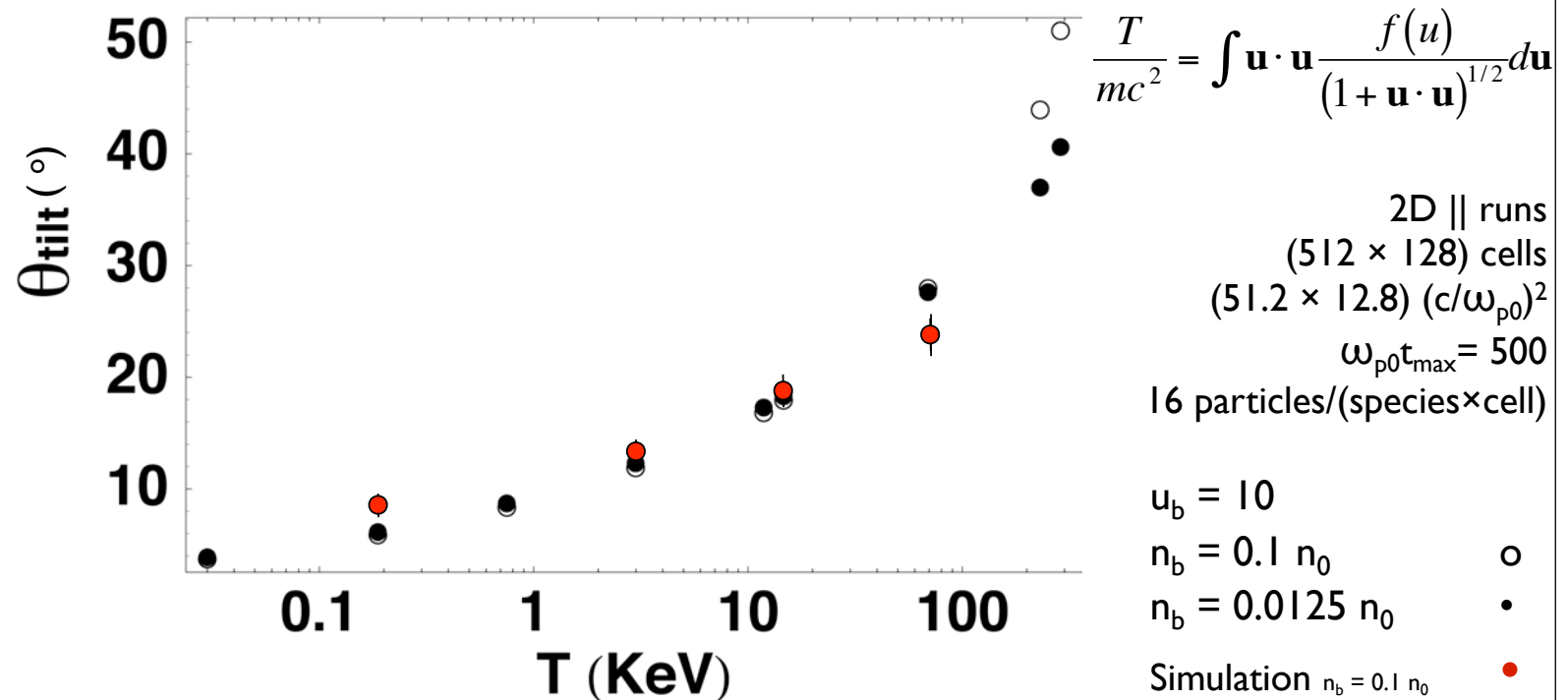
FFT



FFT (zoom)



Comparison theory-simulations II

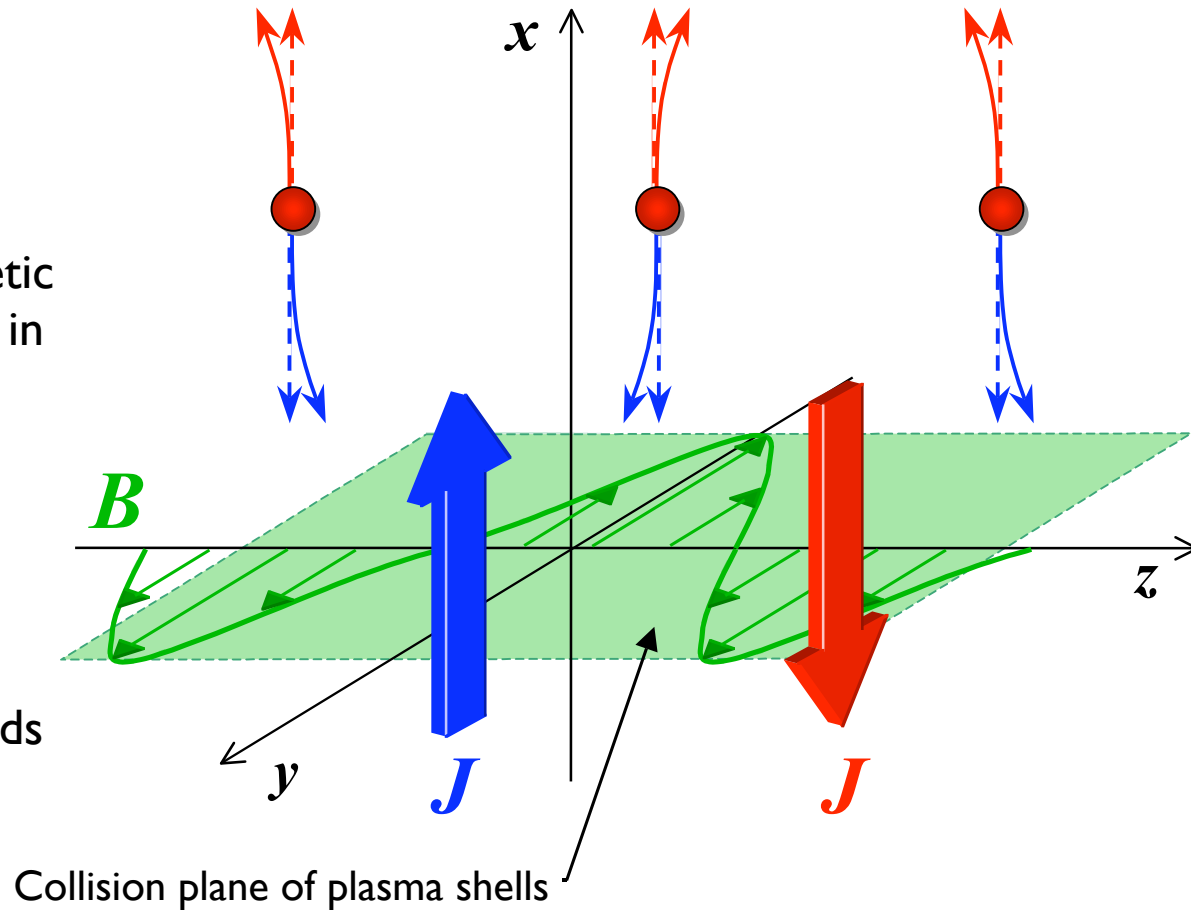


Main consequence of tilted filamentation (for hot beams)
 nonlinear evolution/interaction of filaments
 beam spraying

Weibel instability*

Fluctuation in the magnetic field leads to fluctuation in current density which in turn enhances magnetic field fluctuation

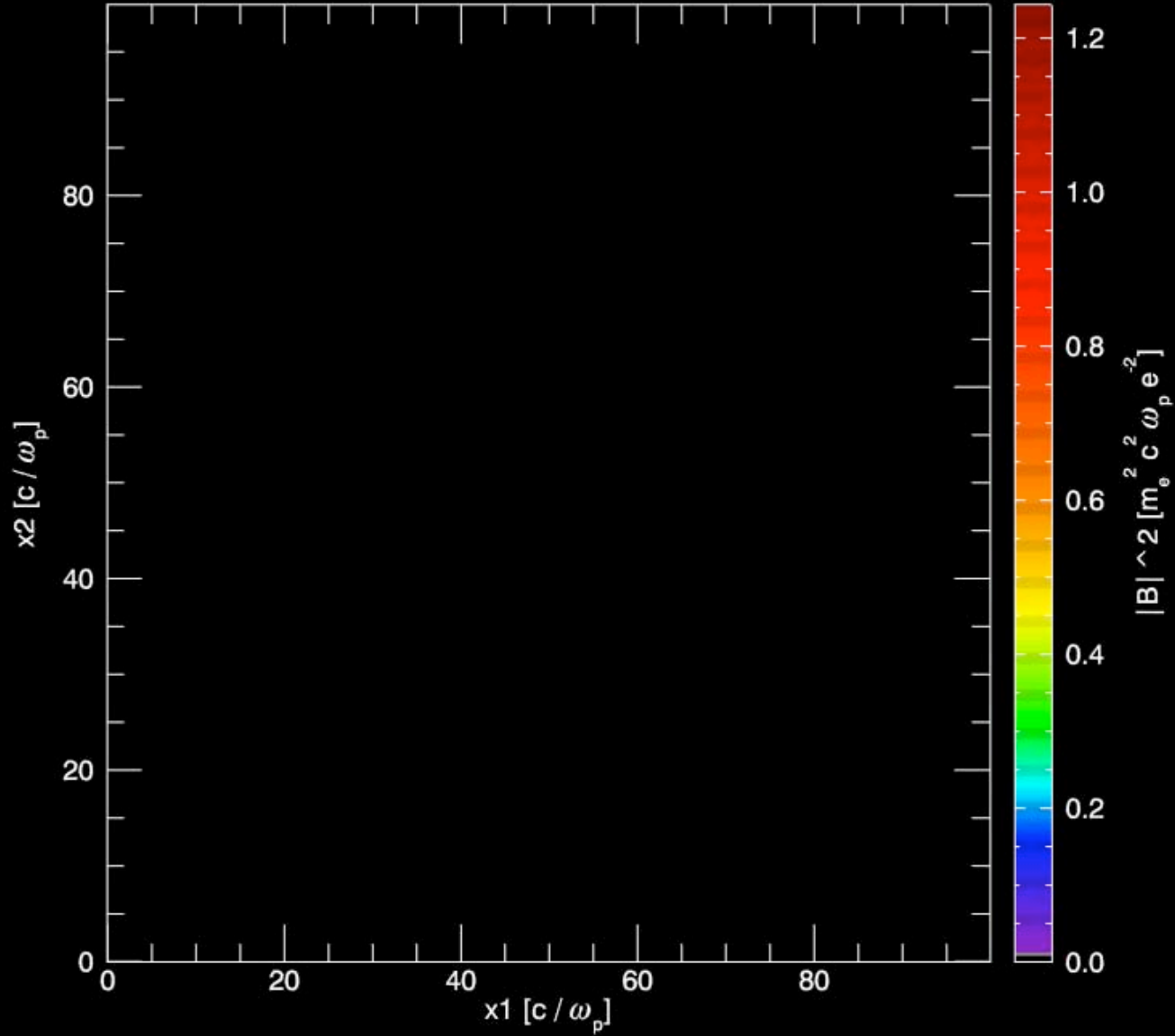
Free energy of particles converted to (mainly) quasi-static magnetic fields



*E.S. Weibel, Phys. Rev. Lett., 1959

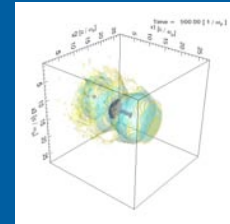
B^2

Time = 0.00 [1 / ω_p]





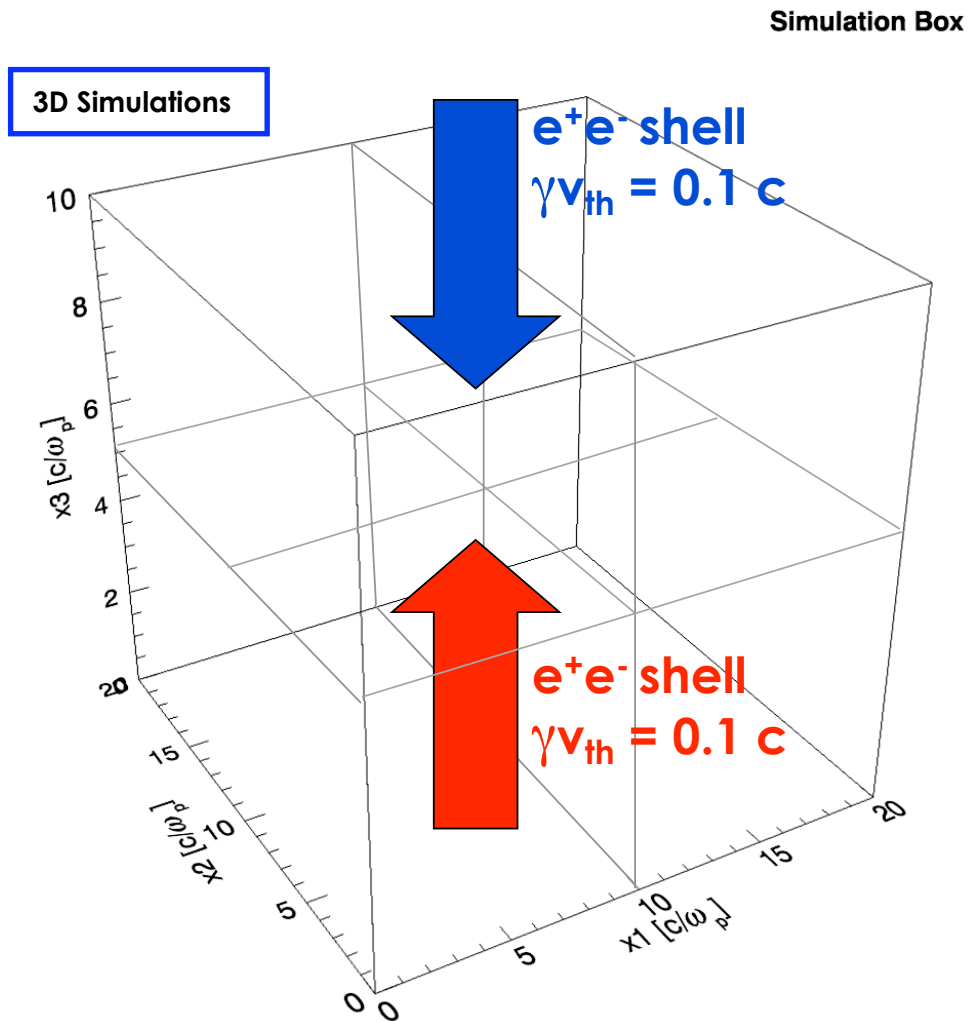
PIC Simulations



What have we learned so far from PIC simulations?

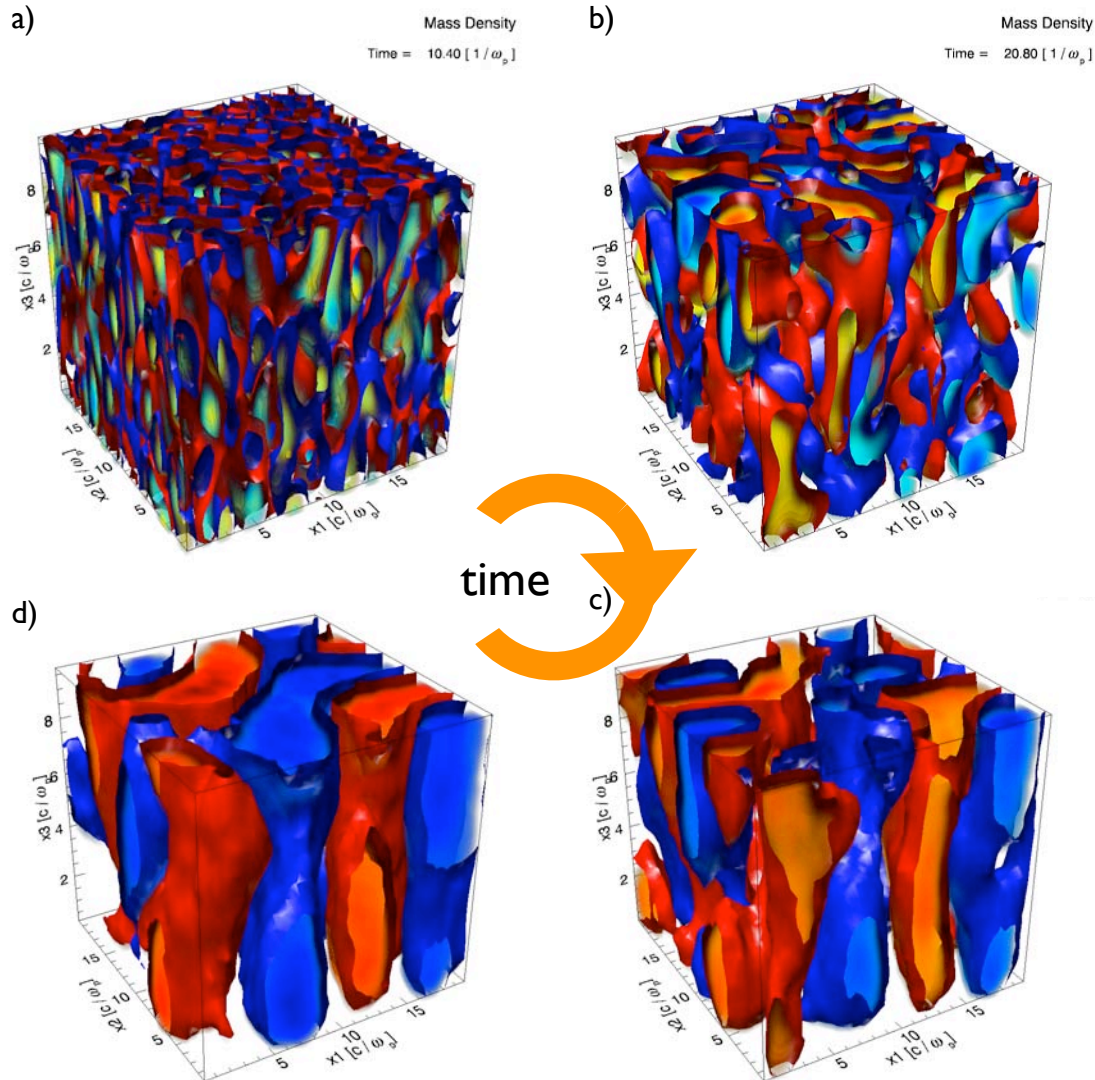
- ⊕ Kazimura et al (1998); Gruzinov (2001); Silva et al (2003); Nishikawa et al (2003); Frederiksen et al (2004); Hededal and Nordlund (2005); Milosavljevic et al (2005)
 - ≡ Weibel seems to work (sub-equipartition B-fields) (e-e+ & very light ions)
 - ≡ Nonlinear dynamics of filaments leads to B-field evolution on a very slow time scale (much slower than $1/\omega_{p0}$)
 - ≡ Reconnection in nonlinear stage of instability leads to heating/acceleration
 - ≡ Radial space-charge fields associated with filaments play a role in particle acceleration/heating
 - ≡ Recent works are already analyzing synchrotron spectra measured in the simulations

Numerical set-up



- **3D simulations**
 - 200 x 200 x 100 cells
(20 x 20 x 10 c^3/ω_p^3)
 - 2 particles per cell species

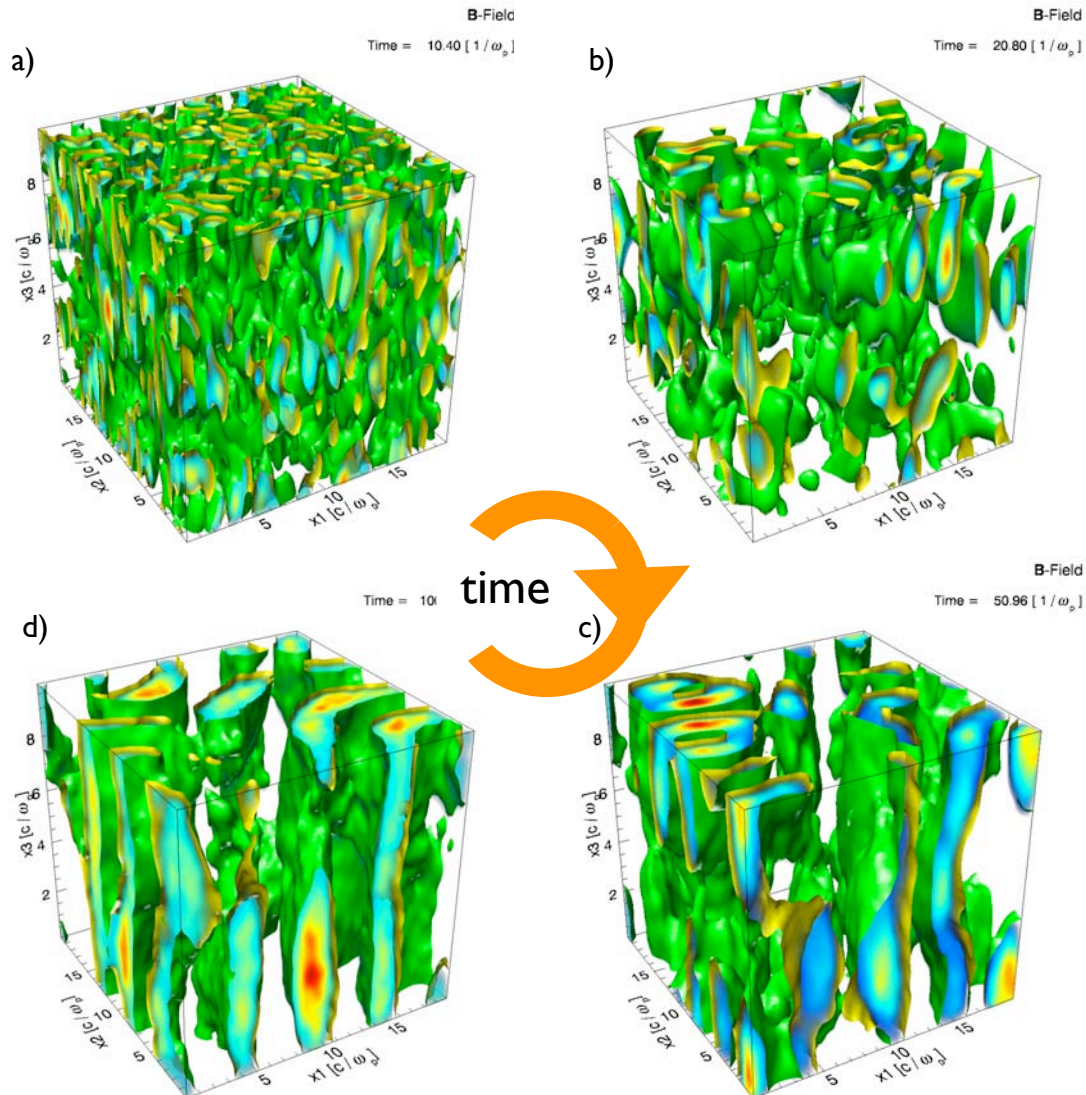
3D results - mass density



- **Red** Iso-surfaces: species with initial positive j_z
- **Blue** Iso-surfaces: species with initial negative j_z
- All isosurfaces drawn at a density value of 1.1 (1.0 being the initial density)
- a) $t = 10.1 \omega_p^{-1}$, b) $t = 20.8 \omega_p^{-1}$, c) $t = 50.96 \omega_p^{-1}$, d) $t = 100.88 \omega_p^{-1}$

Species with opposite initial current are separated by regions of intense magnetic field.

3D results - B-field generation



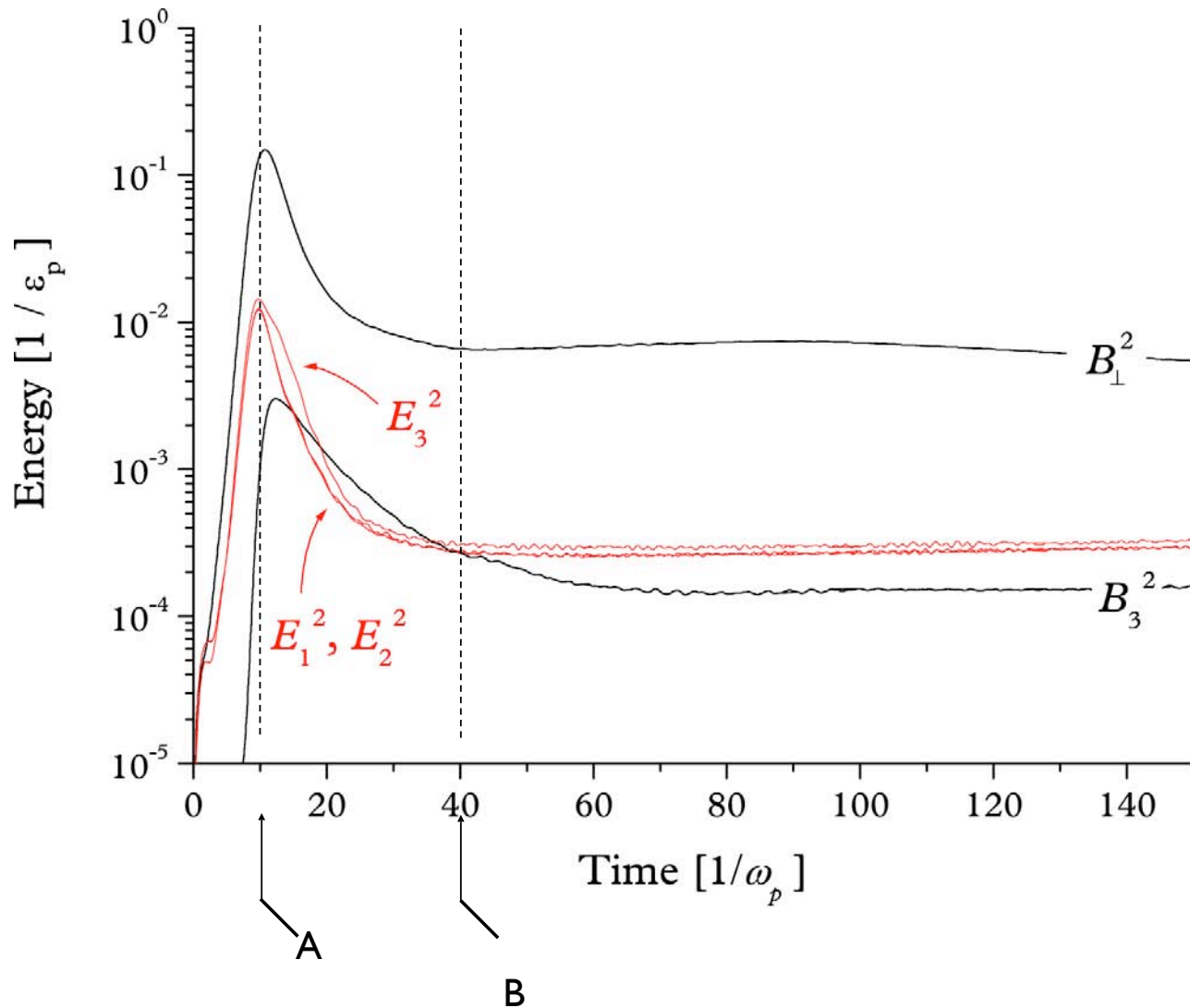
- Isosurfaces (Green - regions of lower values, Yellow regions of higher values) of the magnitude of the magnetic field

- Isosurfaces drawn at a) 0.1, b) 0.025, c) 0.01 and d) 0.006

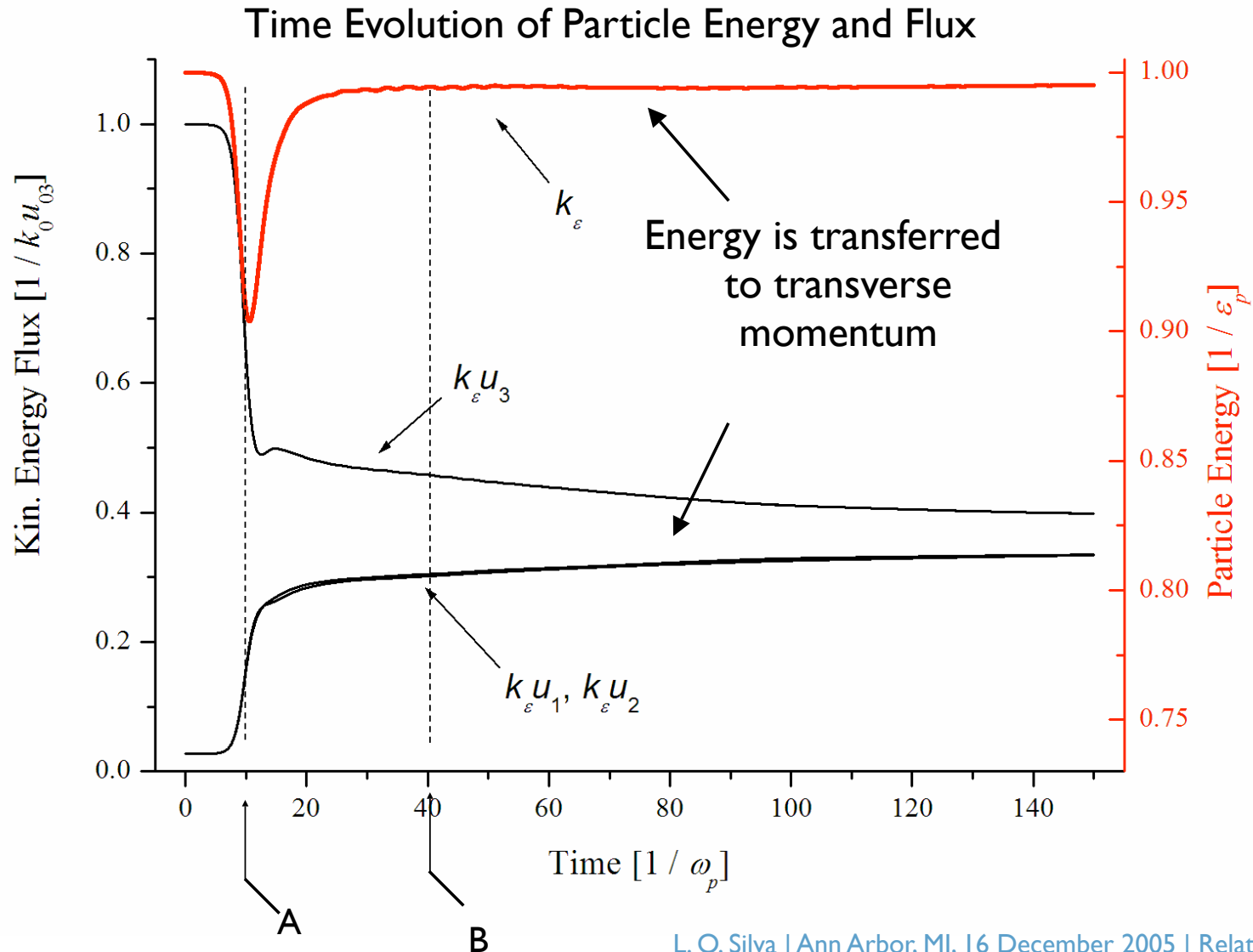
- a) $t = 10.1 \omega_p^{-1}$, b) $t = 20.8 \omega_p^{-1}$, c) $t = 50.96 \omega_p^{-1}$, d) $t = 100.88 \omega_p^{-1}$

The self generated magnetic field exists in the region that separates species with initial opposite current

Field energy evolution - 3D Simulations

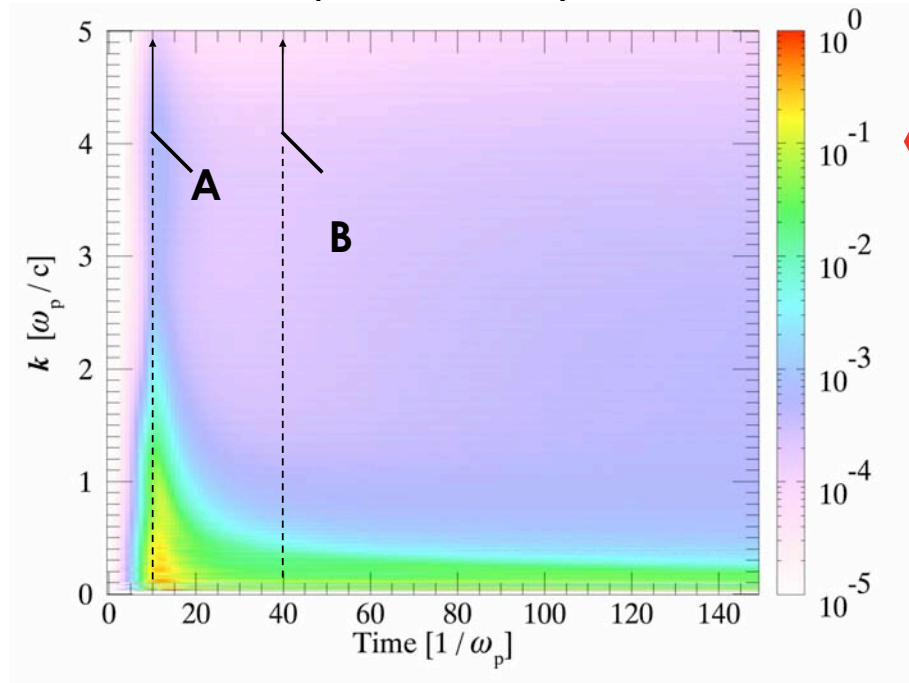


Particle energy evolution - 3D Simulations



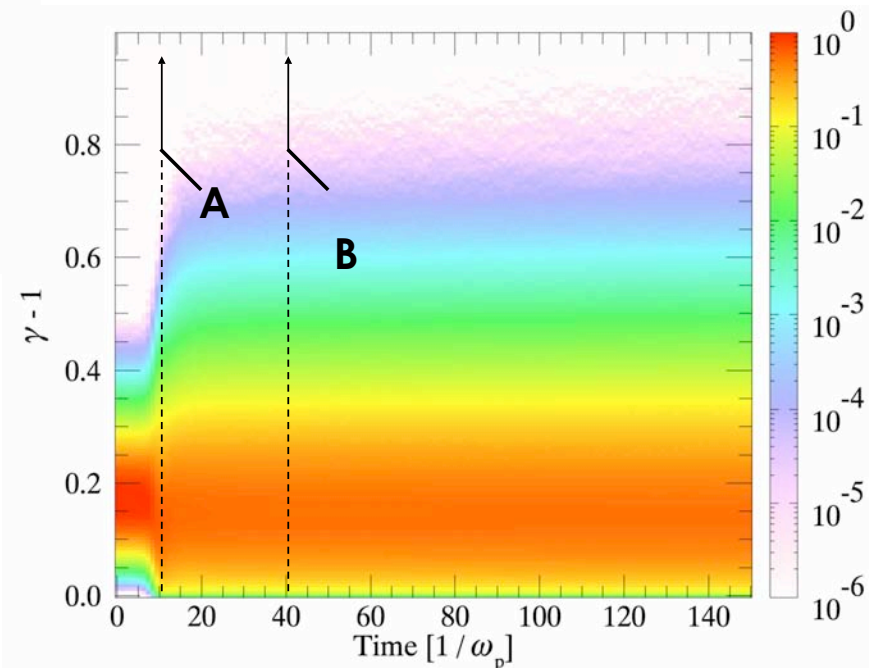
Field structure and particle energy distribution

B-Field Spectral Density Evolution



Magnetic field evolves from small wavelengths to large wavelengths (exponential decay)

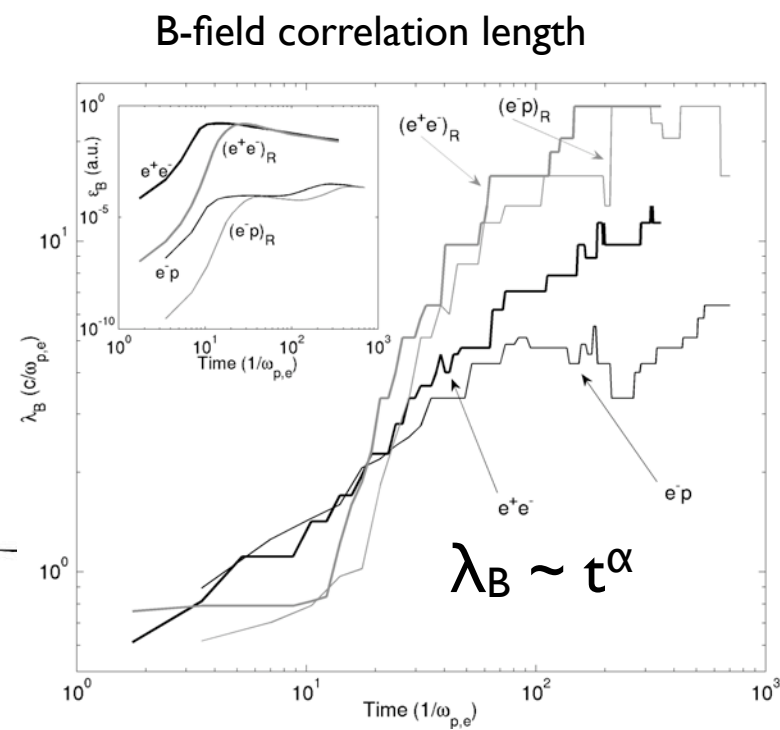
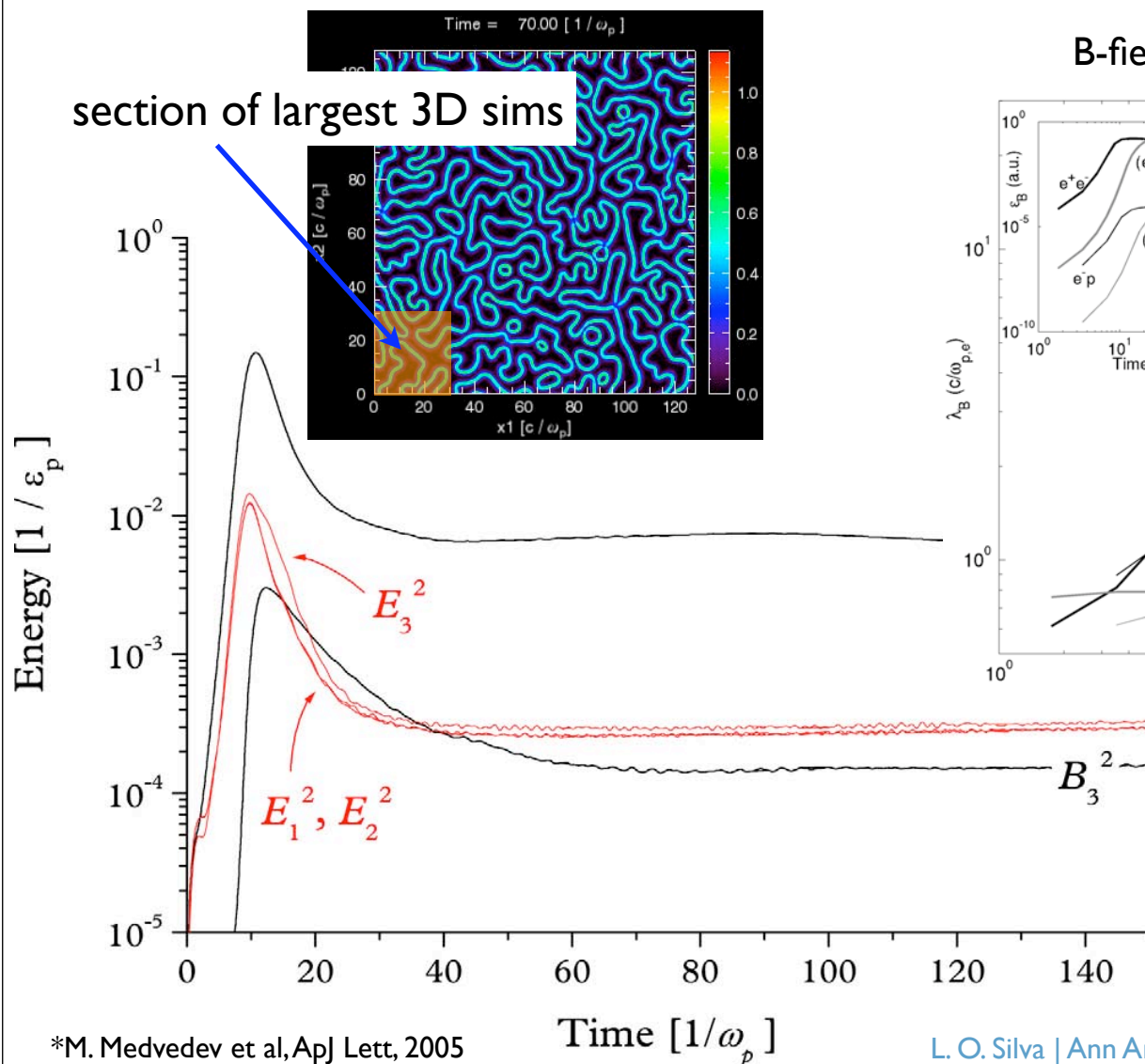
Part. Energy Distribution Evolution



Strong transverse temperature increase and particle acceleration



Long time evolution of field structure*

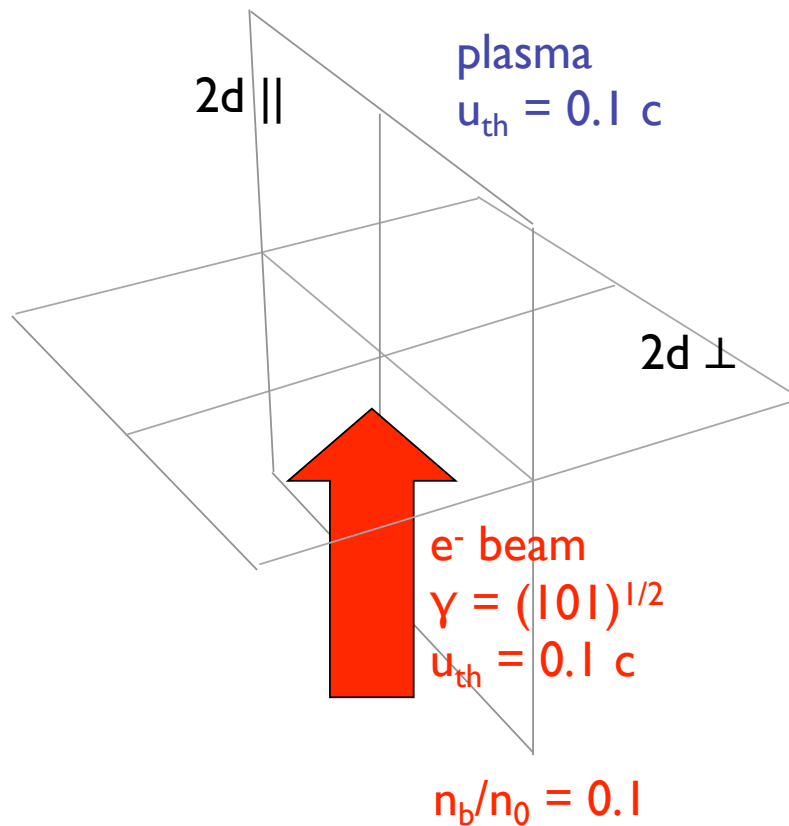


What are the open questions/problems?

- Long time evolution of B-field for more realistic geometries
- Configuration of simulations not appropriate to check particle acceleration and relativistic shock formation:
 - ≡ either periodic systems or systems where plasma injected from side wall
- Collision of electron-”proton” shells
 - ≡ typical length scale and time scale increase by $(m_i/m_e)^{1/2}$
 - ≡ grand challenge for 3D simulations

Increase simulation box size and interaction time by reducing dimensionality of the simulations and using moving window

Simulation parameters (PARSEC)*



Periodic system
Current and charge neutral
Moving ions

- 3D runs

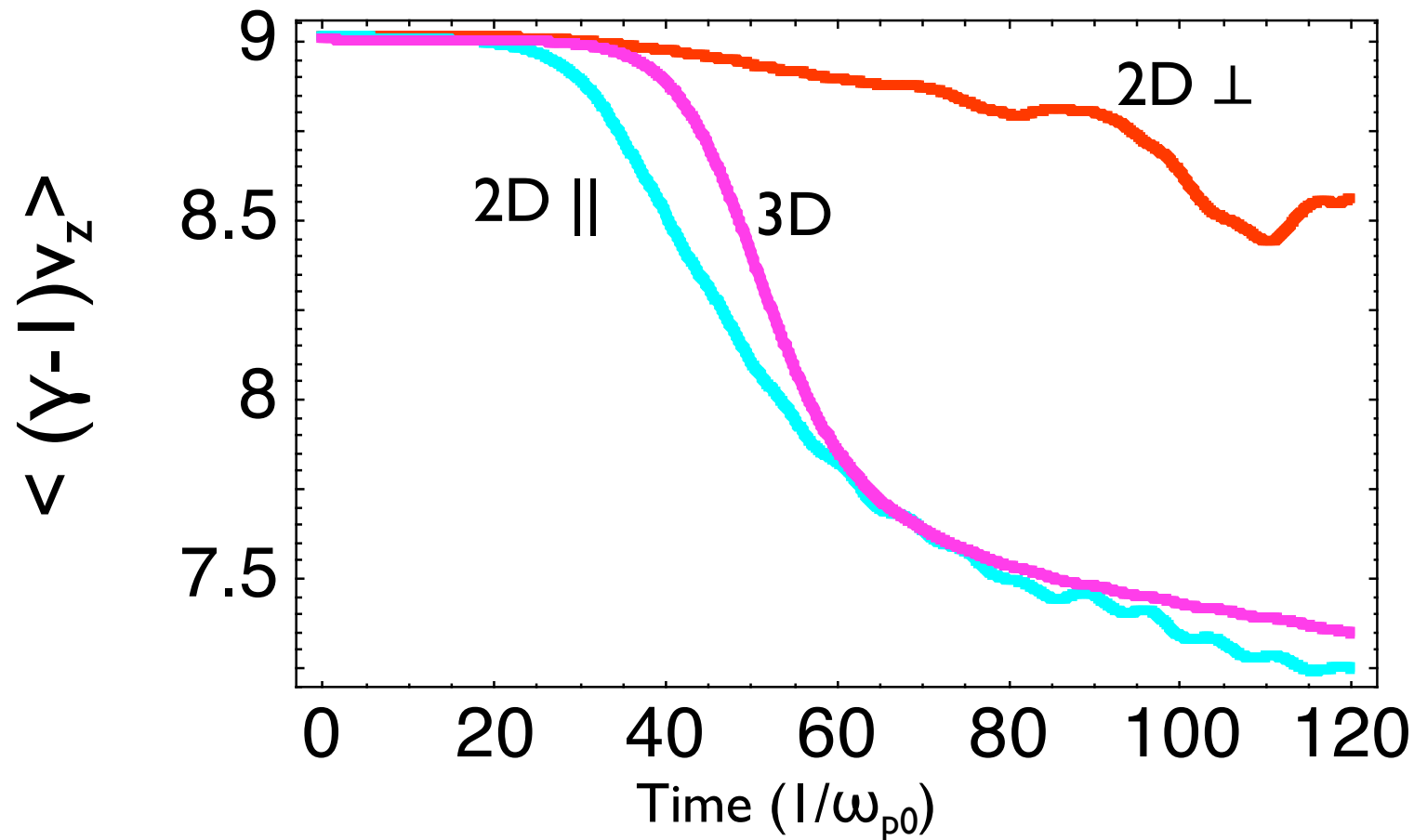
$(128)^3$ cells
 $(12.8 c/\omega_{p0})^3$
 $\omega_{p0}t_{max} = 120$
16 particles/cell

- 2D runs (\parallel and \perp)

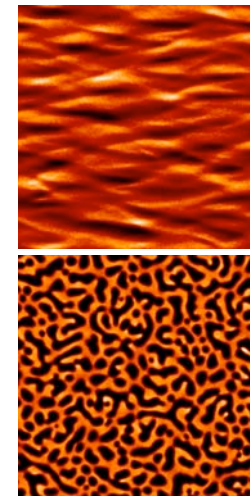
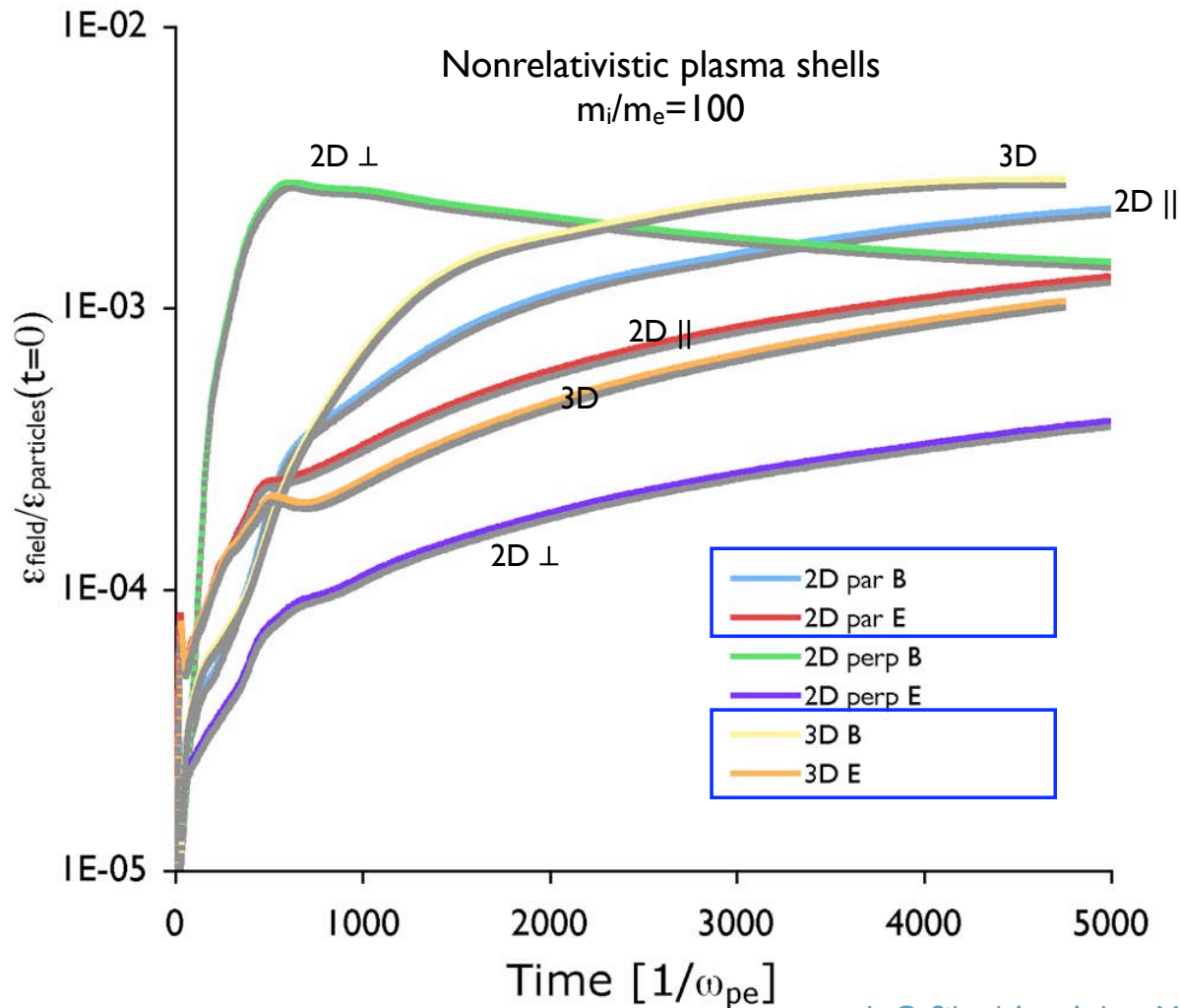
$(128)^2$ cells
 $(12.8 c/\omega_{p0})^2$
 $\omega_{p0}t_{max} = 120$
16 particles/cell

*J. Tonge, UCLA PhD Thesis, 2002

Kinetic energy flux per particle



Can lower dimensional simulations capture the [relevant] physics?



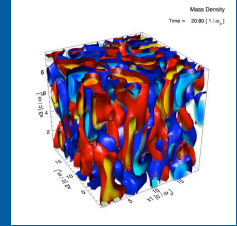
2D ||

2D ⊥

3D simulation box
 $(25.6 c/\omega_{p0})^3$
 2D simulation box
 $(256.0 c/\omega_{p0})^2$



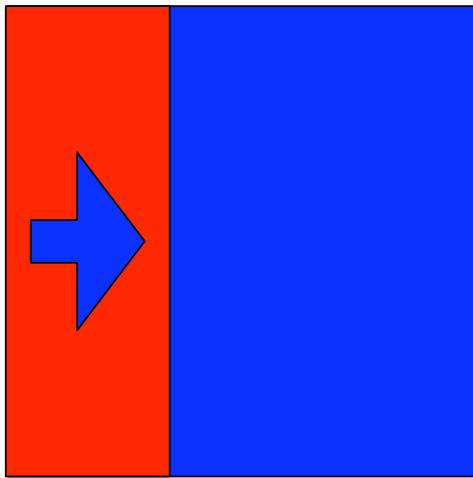
Magnetic field generation



B-field evolution on large scales/long times

Relativistic ($\gamma = 10, 5$) plasma shell
injected from 2D simulation wall

$m_i/m_e = 100$
16 particles/(cell species)

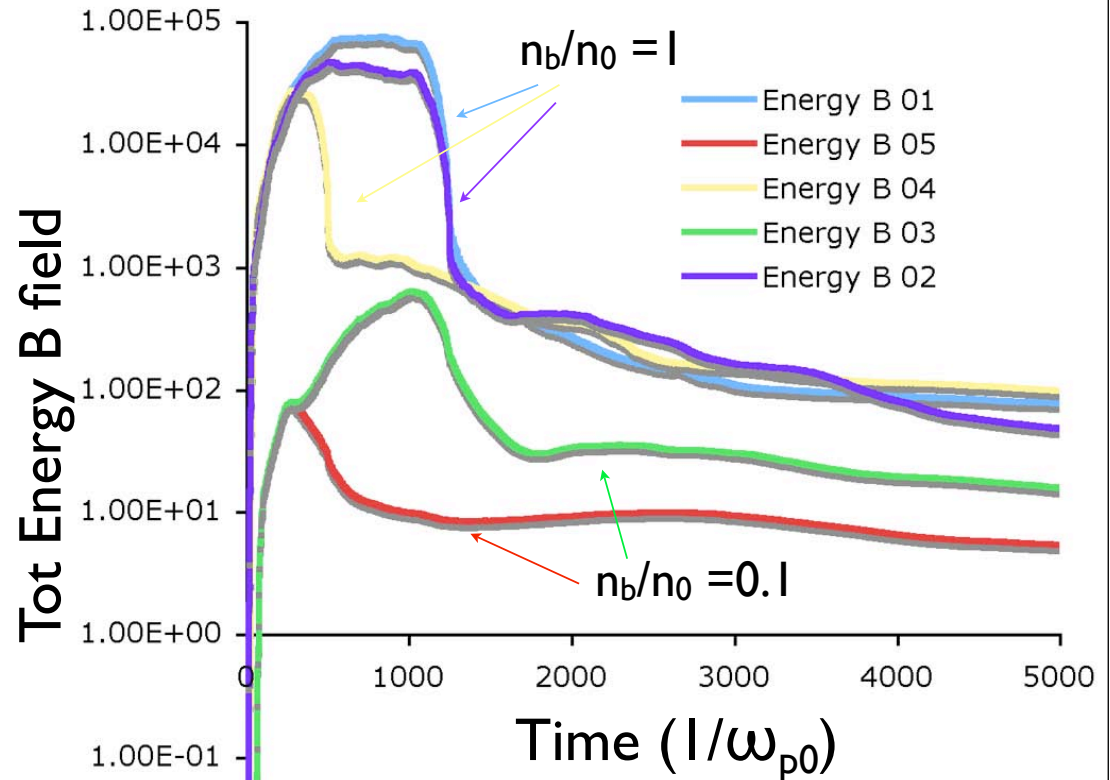


$256 c/\omega_{p0} \times 256 c/\omega_{p0}$

$t_{on} = 0$

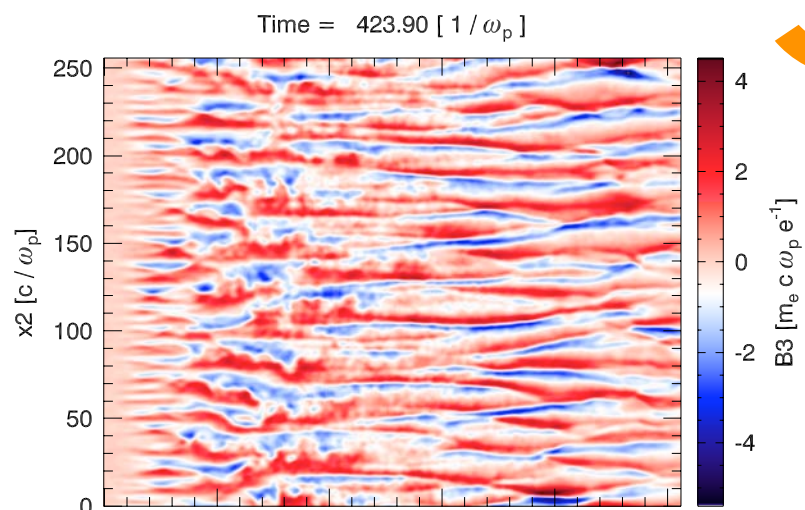
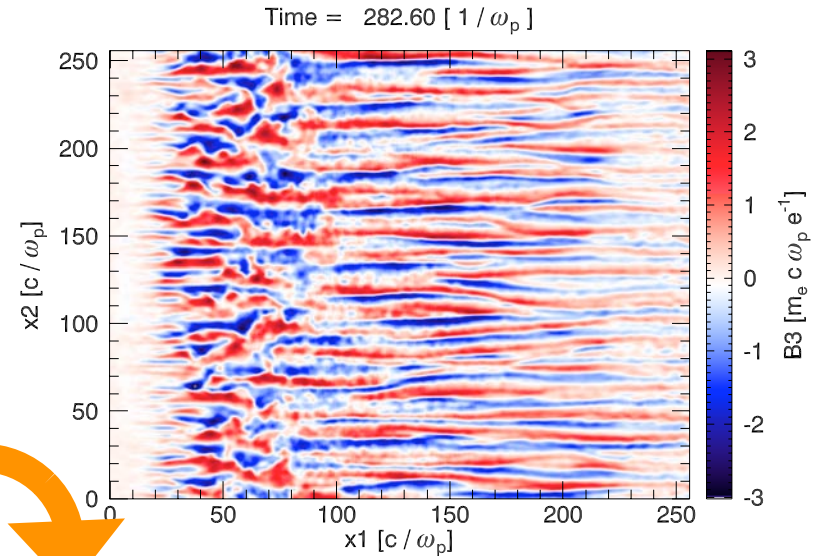
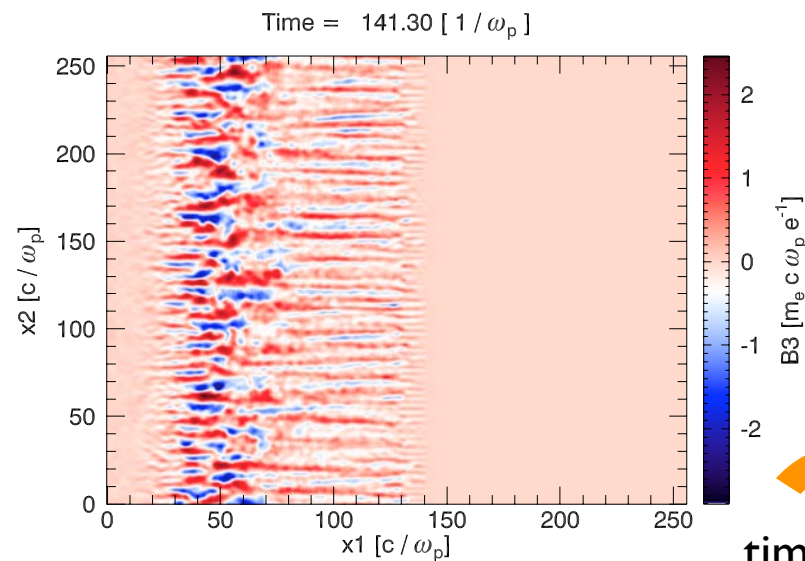
$t_{off} = 1000/\omega_{p0}, 256/\omega_{p0}$

$n_b/n_0 = 1, 0.1$



@ $n_0 = 10 \text{ cm}^{-3}$
 $256 c/\omega_{p0} \sim 400 \text{ km}$
 $1000/\omega_{p0} \sim 5 \text{ ms}$

Weibel instability on electrons - B field



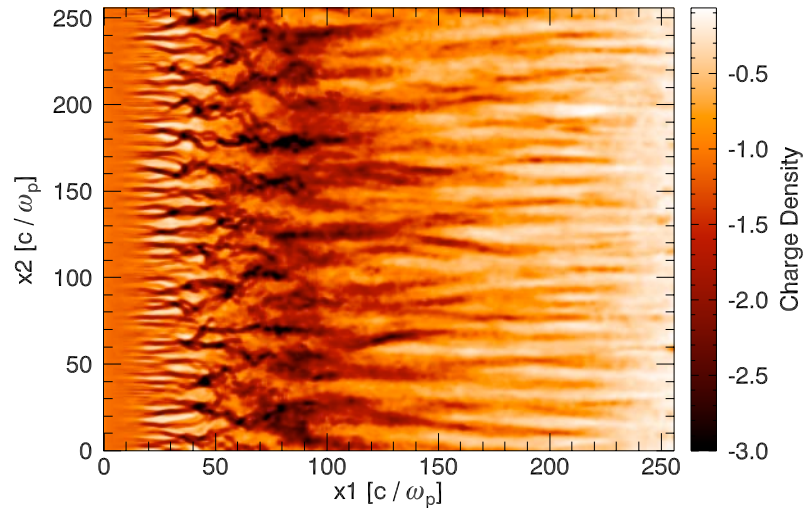
time

Spatio-temporal growth is evident with clear dynamics of tilted filament coalescence to large scales

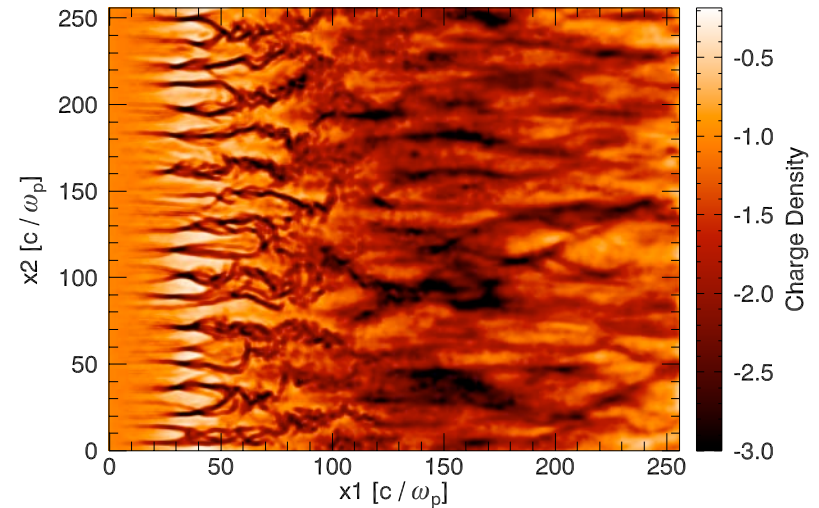
Spatio-temporal dynamics seeds Weibel on ions

Weibel instability on electrons

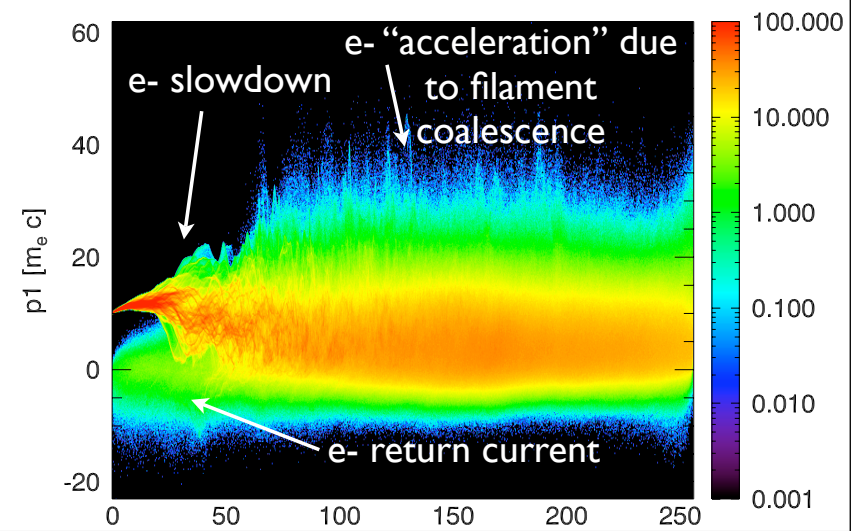
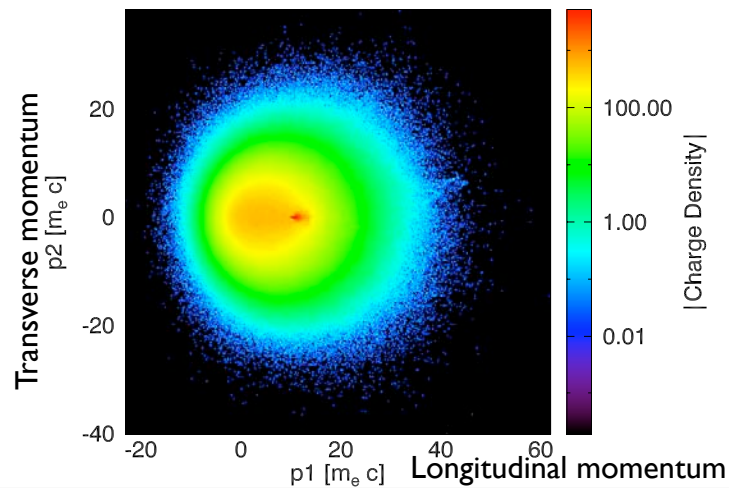
Time = 282.60 [$1/\omega_p$]



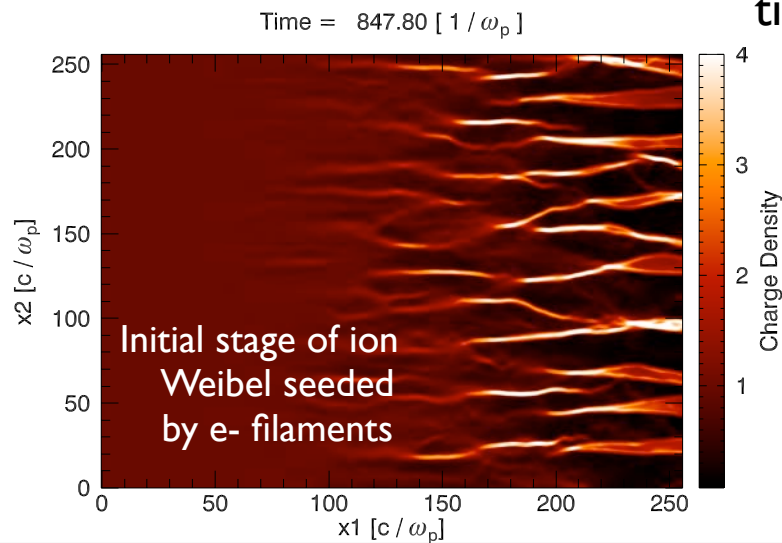
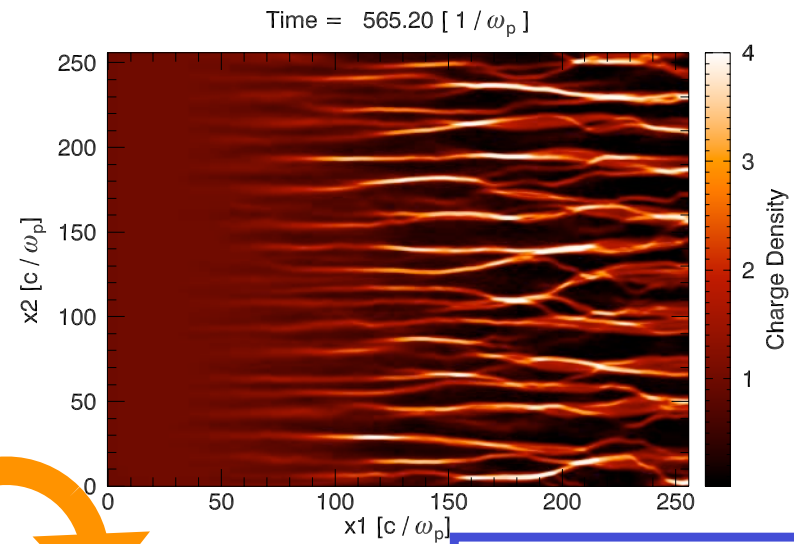
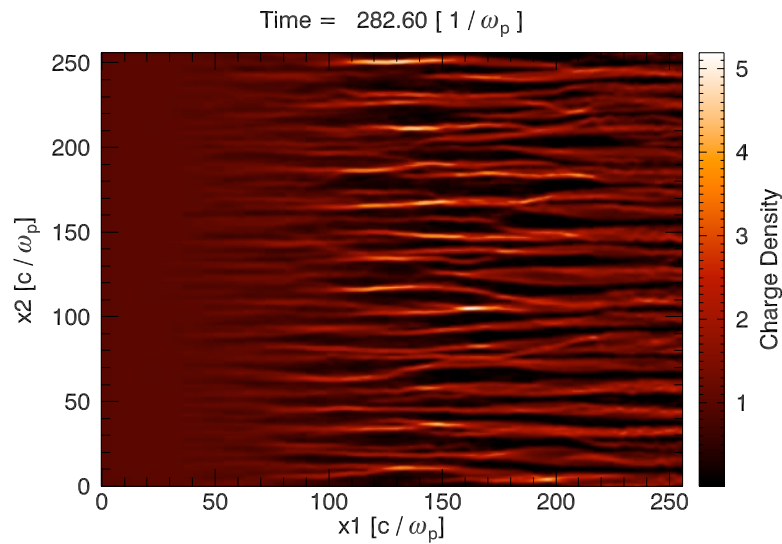
Time = 565.20 [$1/\omega_p$]



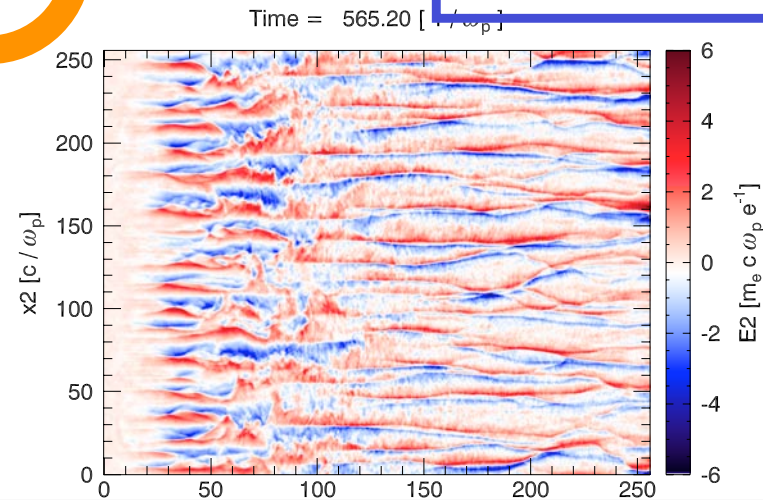
Time = 565.20 [$1/\omega_p$]



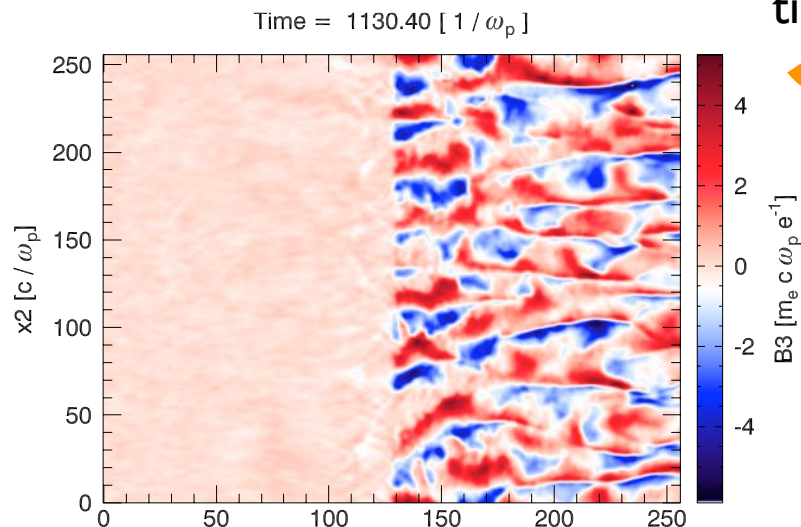
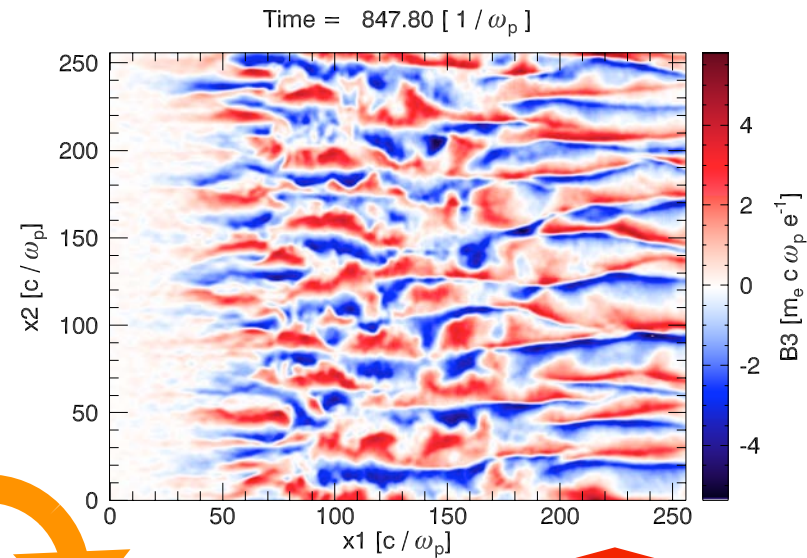
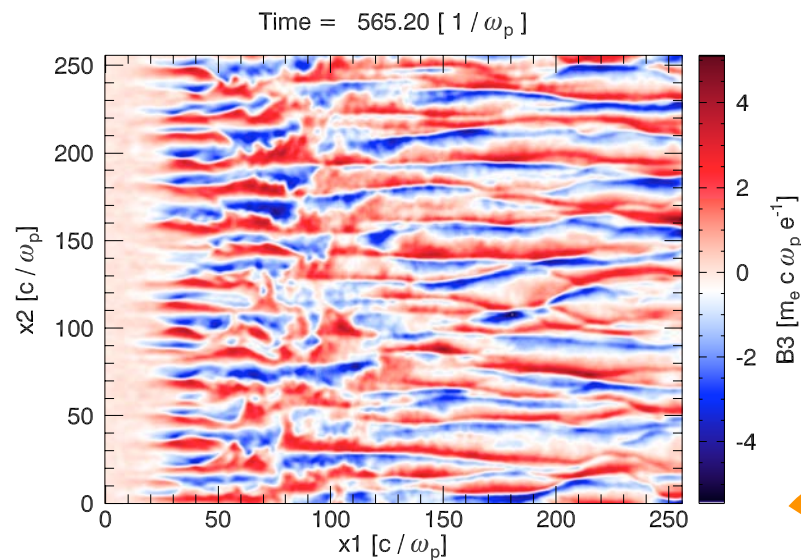
Weibel instability - ion density



E2 field $\gamma_{\max} e^- \sim E_{\max} L_{\text{filaments}} \sim 60$



Weibel instability on ions - B field



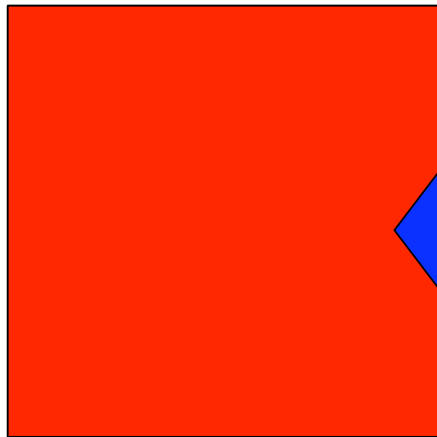
time

Large scale evolution can only be captured with long and wide box

For $n_b \sim n_0$ the self-generated field is transported by the relativistic plasma slab

B-field evolution on the head of relativistic shell I

$m_i/m_e = 100$
16 particles/(cell species)



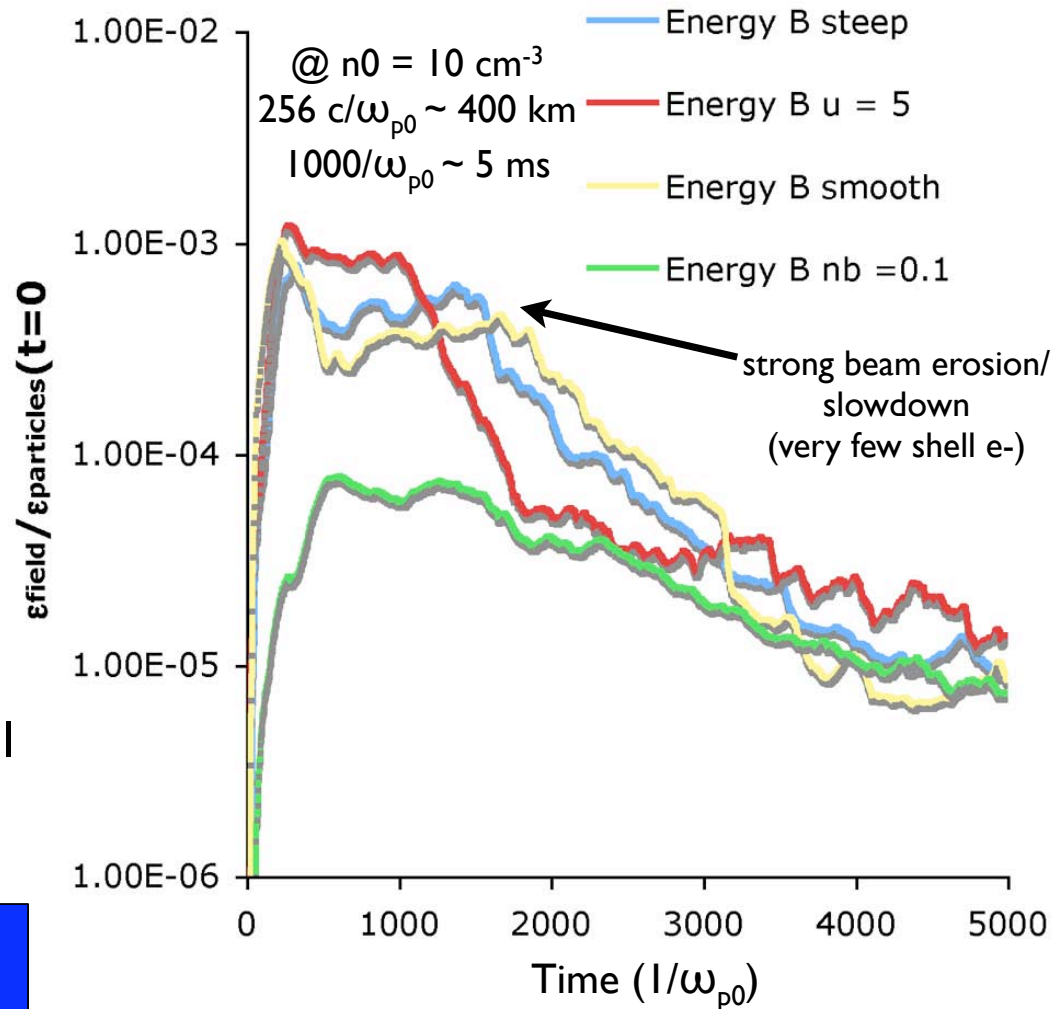
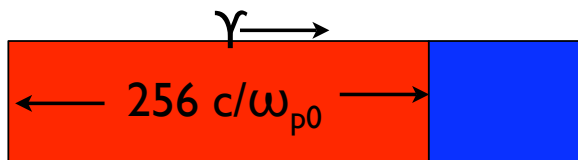
Incoming stationary plasma

$256 c/\omega_{p0} \times 256 c/\omega_{p0}$

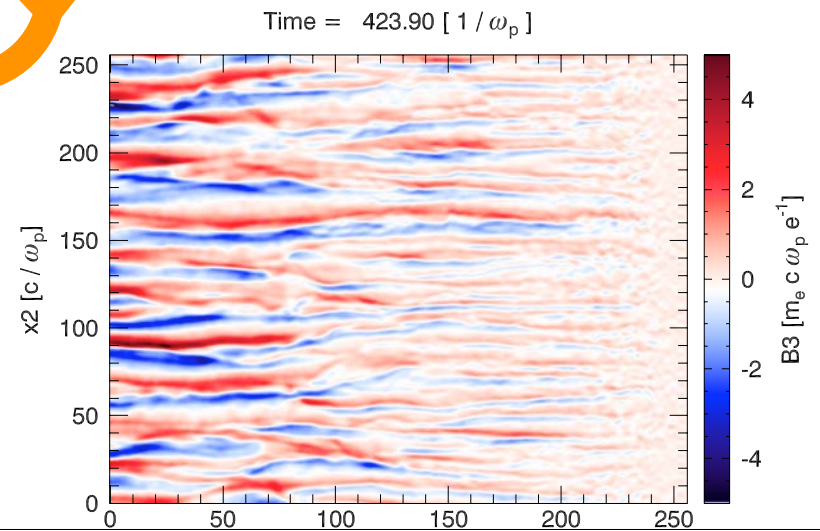
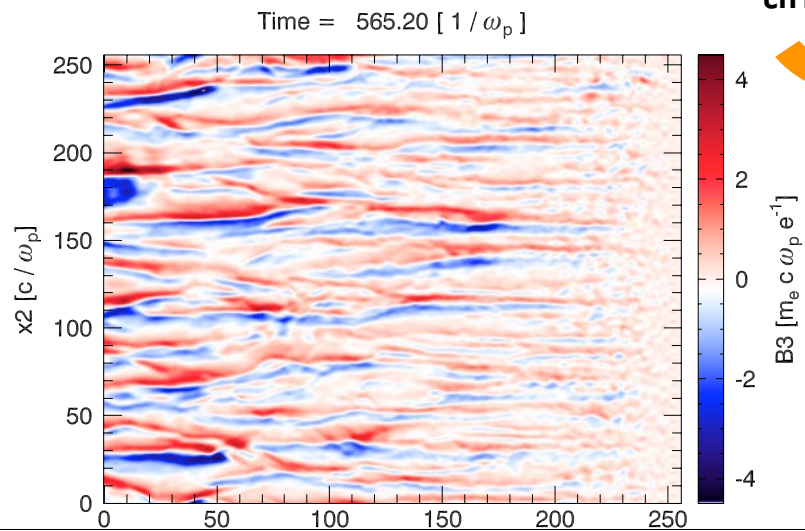
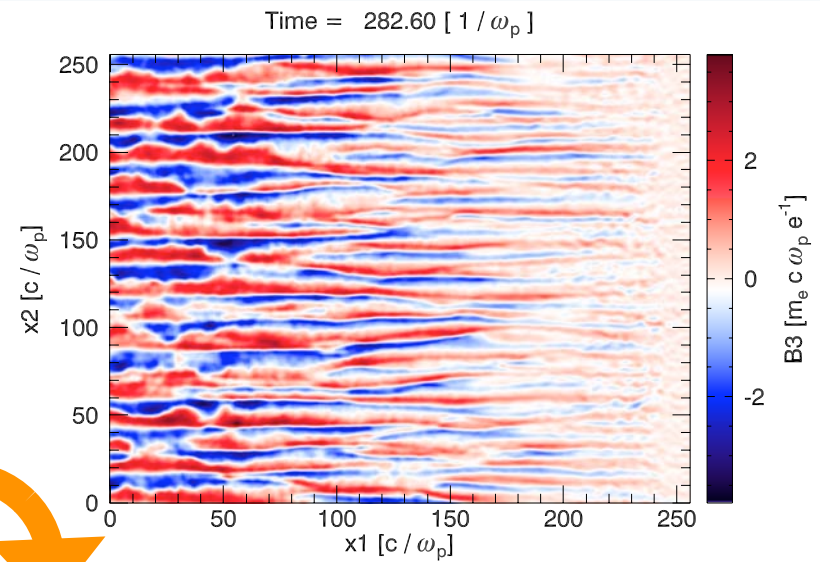
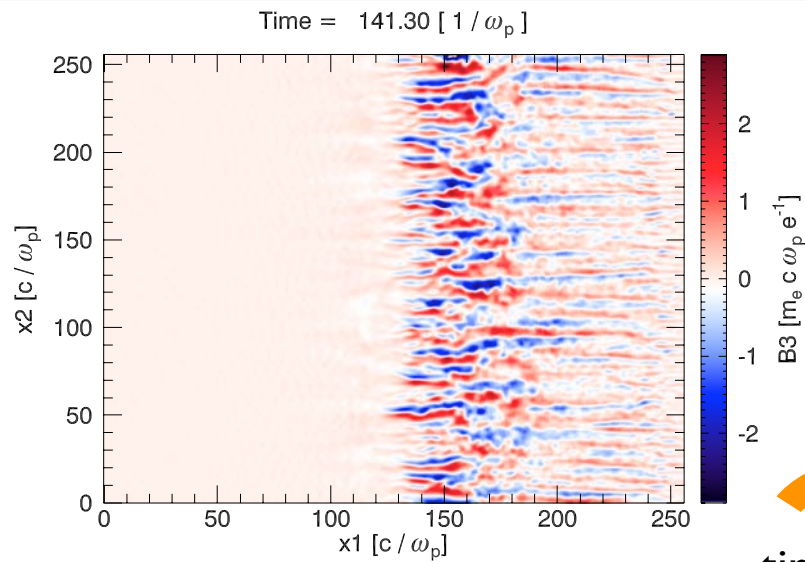
Moving window @ c

Relativistic ($\gamma = 10, 5$) plasma shell

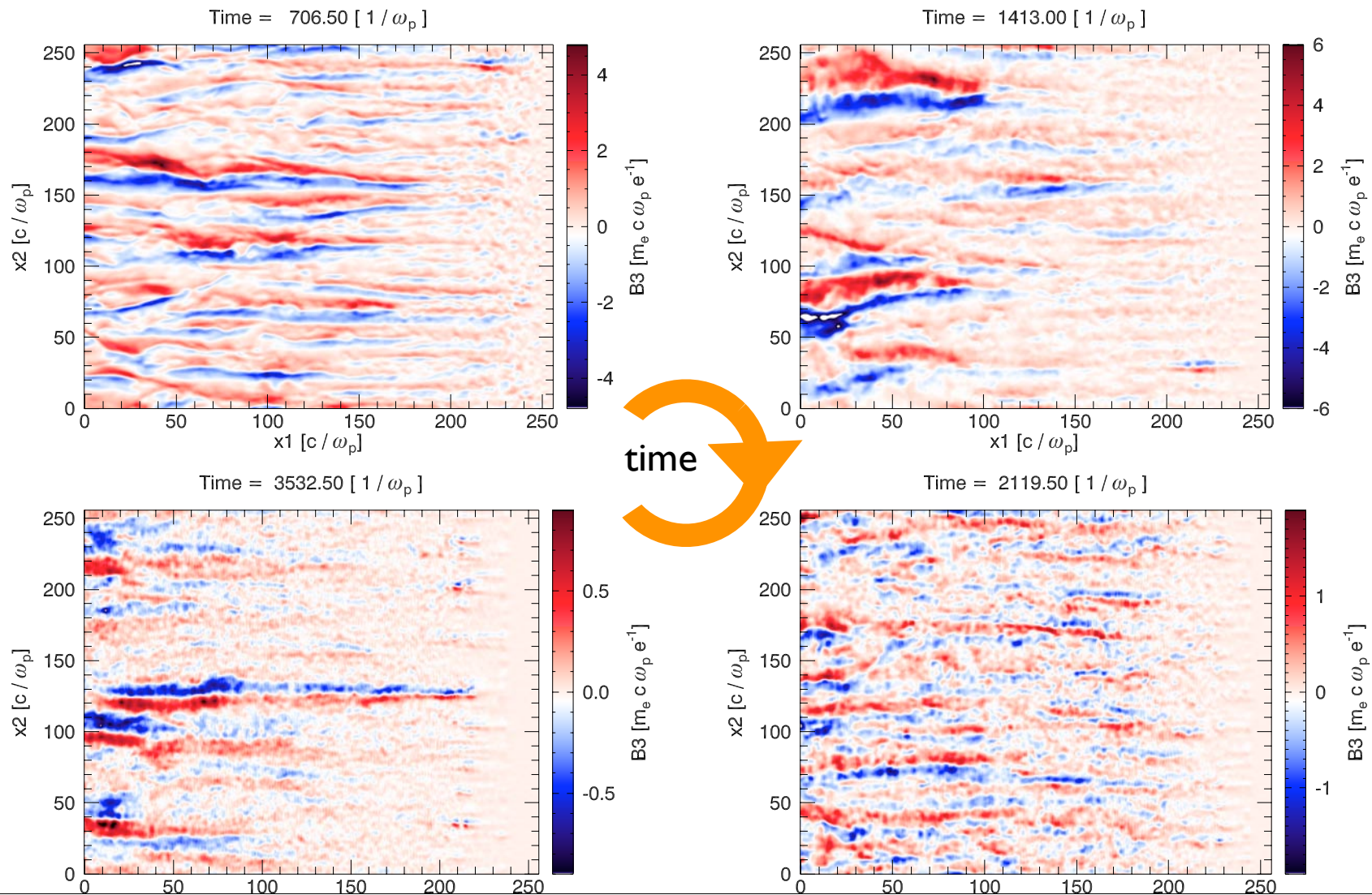
$n_b/n_0 = 1, 0.1$



B-field on the head of relativistic shell - Weibel on e-

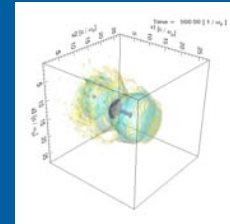


B-field on the head of relativistic shell - Weibel on i



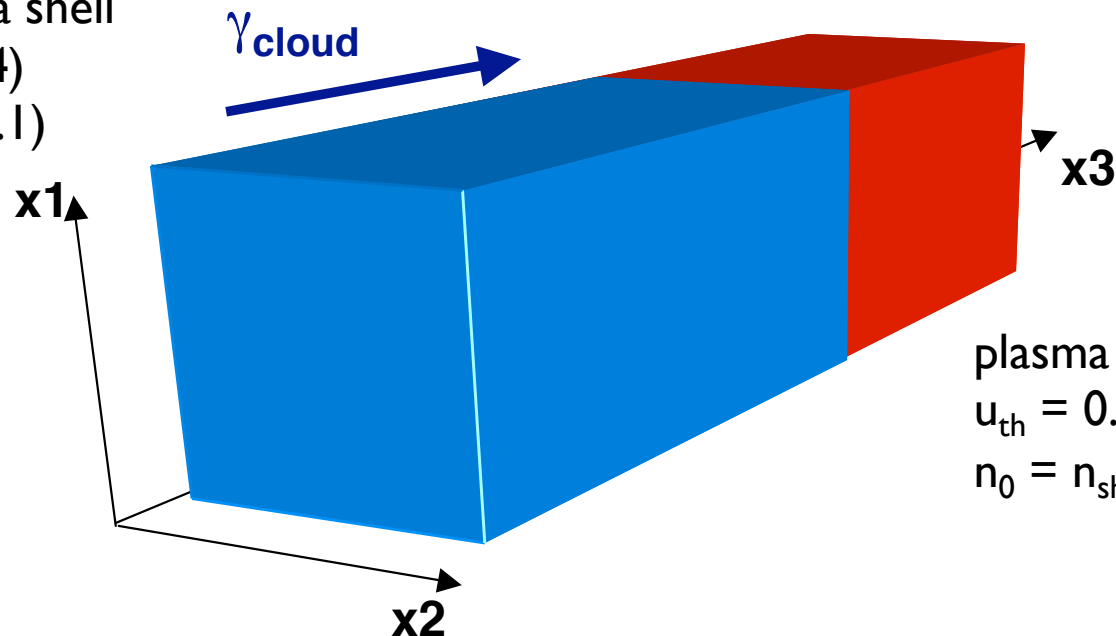


Particle acceleration in plasma collisions



3D PIC simulations of e-e⁺ shell collisions with moving window

e⁺e⁻ plasma shell
 $\gamma = 10$ (1.4)
 $u_{\text{th}}/c = 1$ (0.1)



plasma
 $u_{\text{th}} = 0.1 c$
 $n_0 = n_{\text{shell}}$

3D runs

64×64×1600 cells

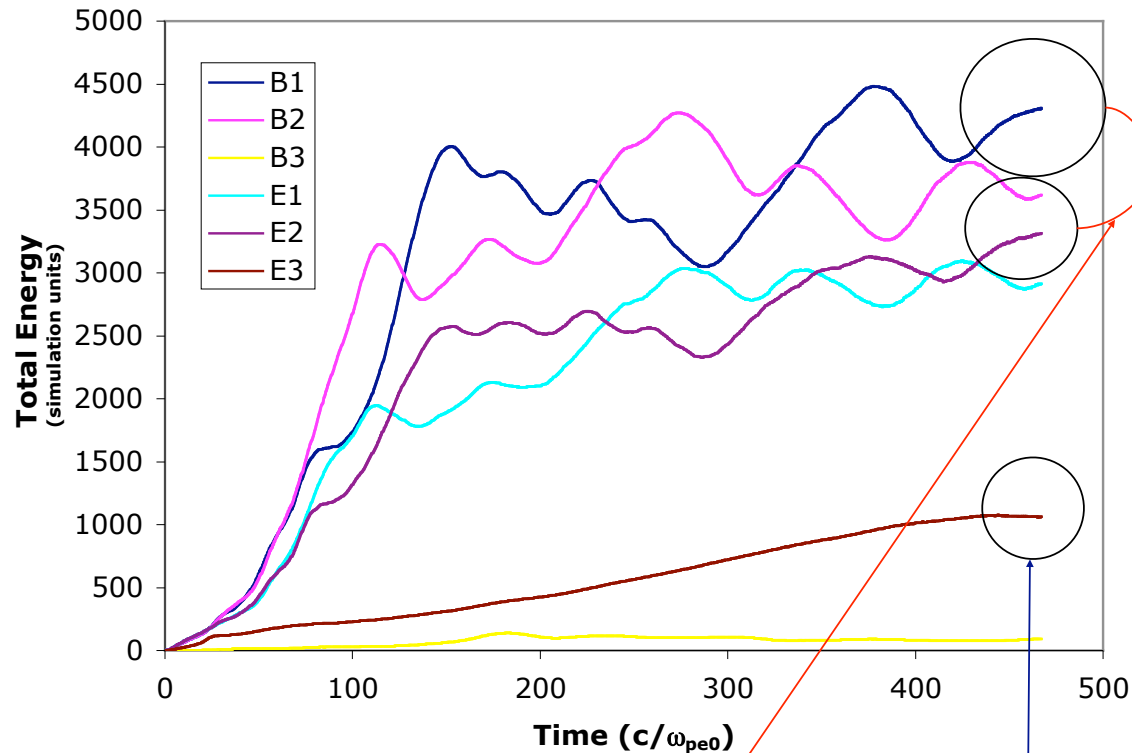
12.8×12.8×320.0 (c/ω_{p0})³

ω_{p0}t_{max} = 500 (250)

8 particles/(cloud×cell)

Moving window (@ c) follows shock front

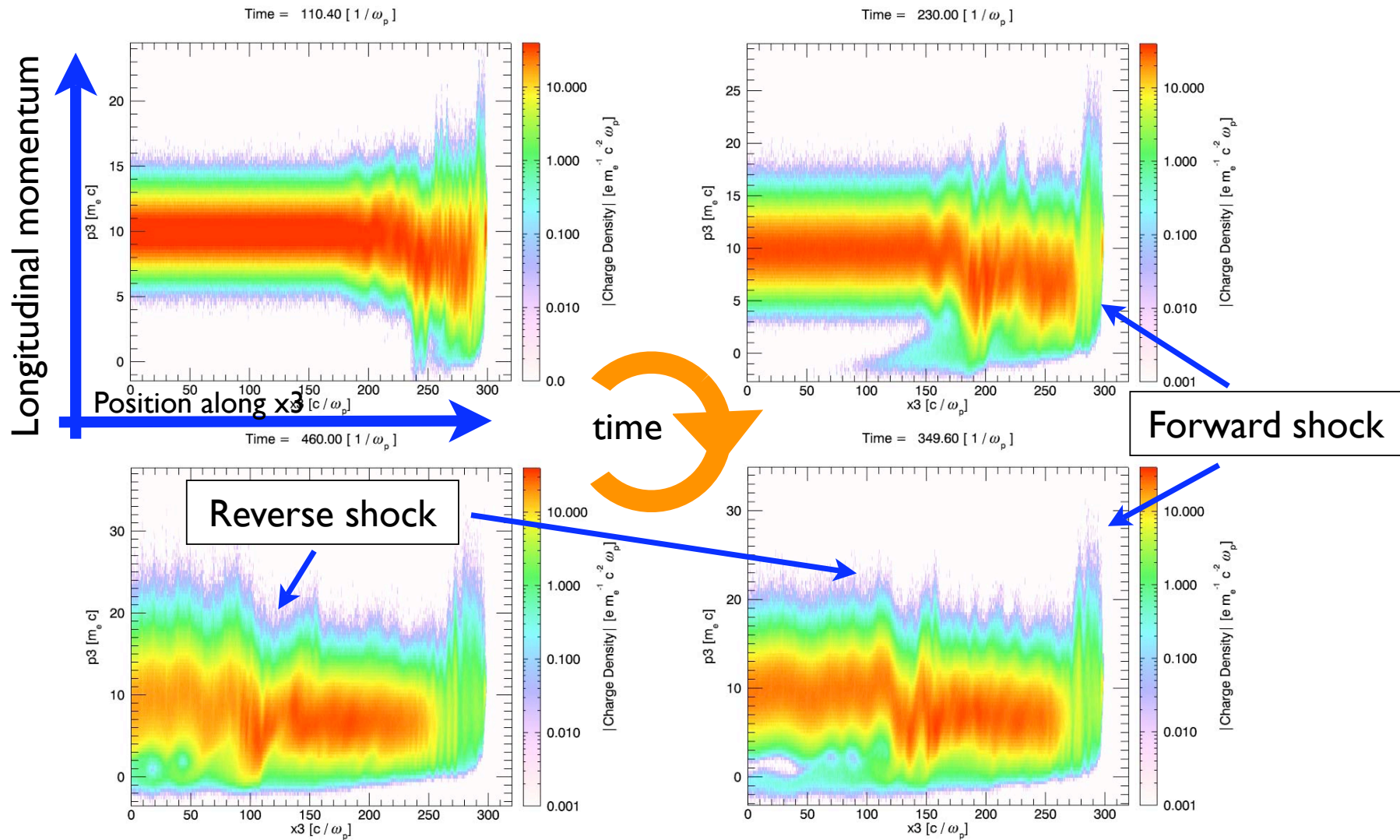
Time history for fields ($\gamma_{\text{shell}} = 10$)



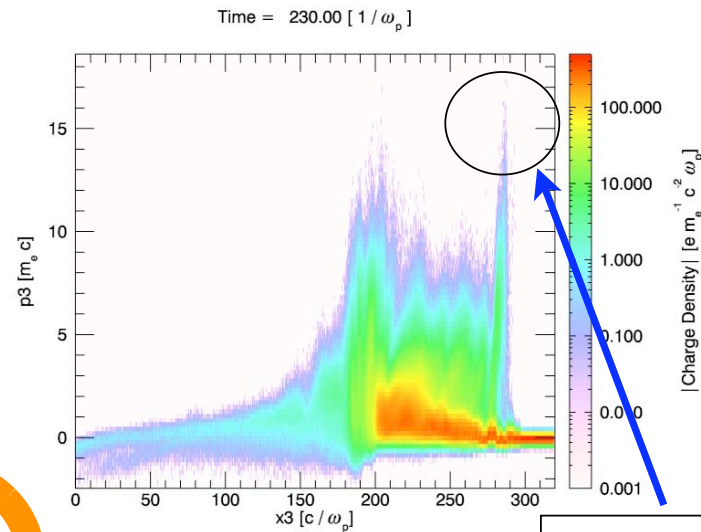
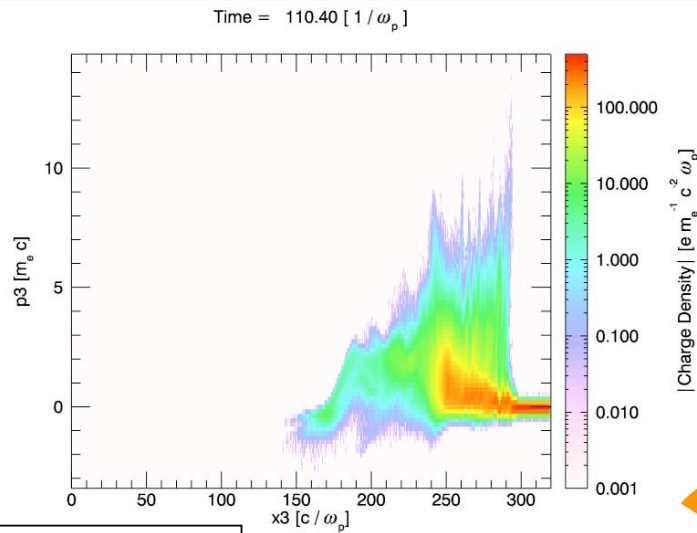
Quasi-steady
state in the fields

Total energy in B-field $\sim 10^{-4}$ total initial kinetic energy
Comparable energy in B and E (crucial for particle acceleration)
E₁ and E₂ comparable to B₁ and B₂ (e.m. field)
E₃ is also significant (crucial for particle acceleration)

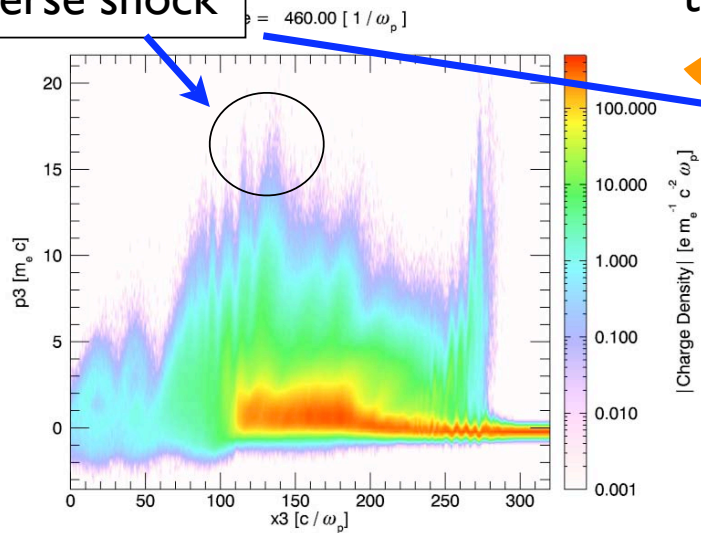
Shock formation ($\Upsilon_{\text{cloud}} = 10$) - shell particles



Shock formation ($\gamma_{\text{shell}} = 10$) - background plasma



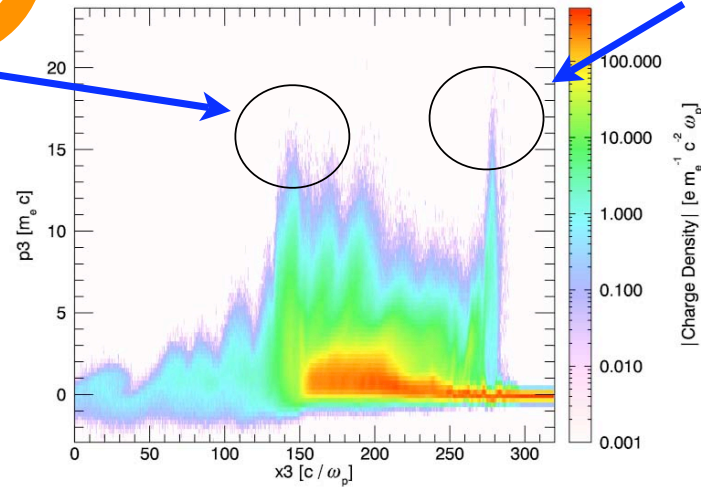
Reverse shock



time

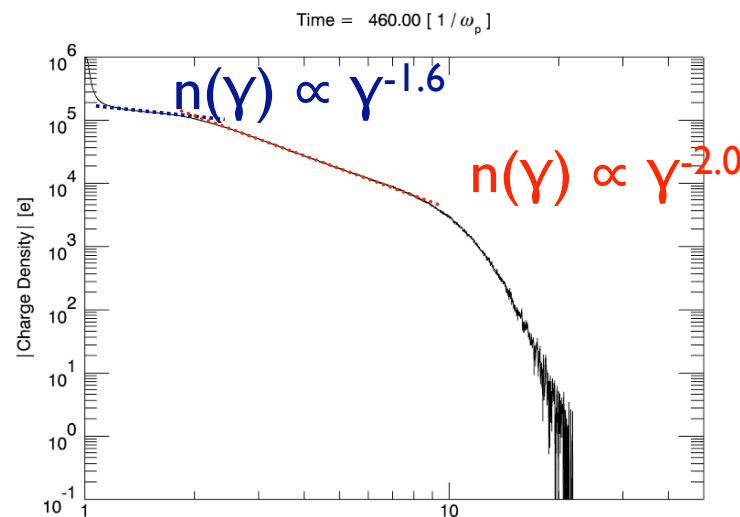
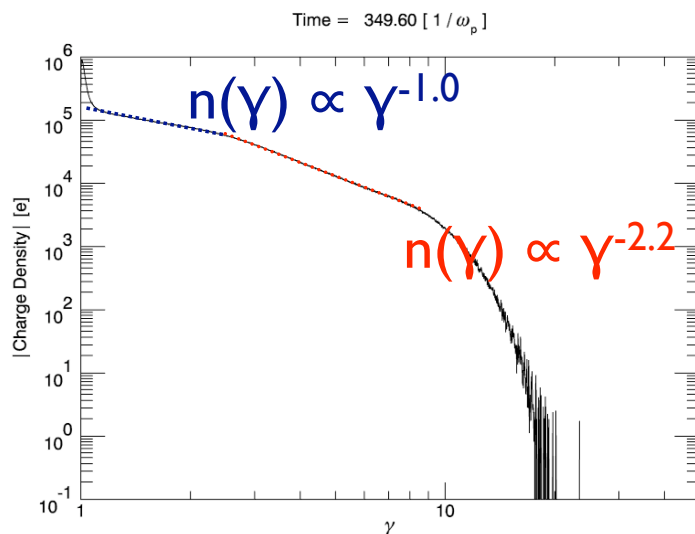
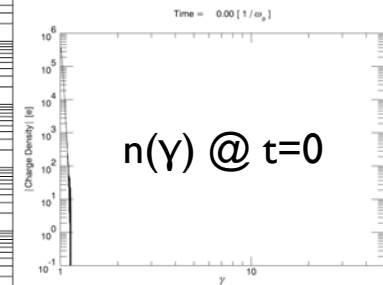
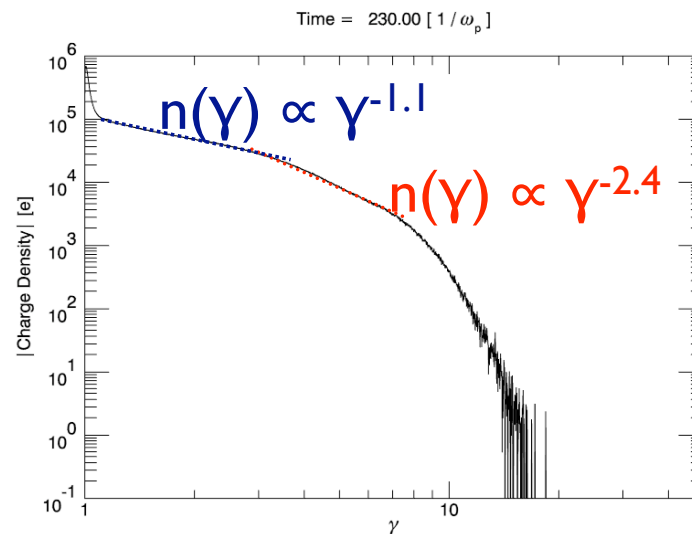
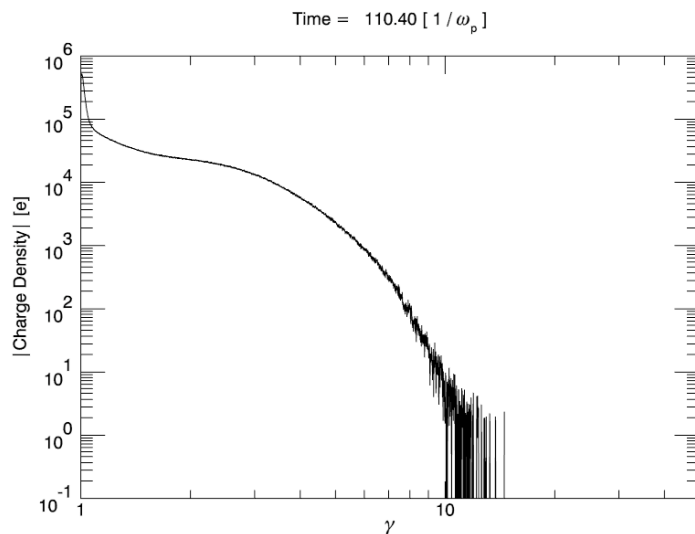


Time = 349.60 [1 / ω_p]



Forward shock

Particle acceleration ($\gamma_{\text{shell}} = 10$) - background plasma



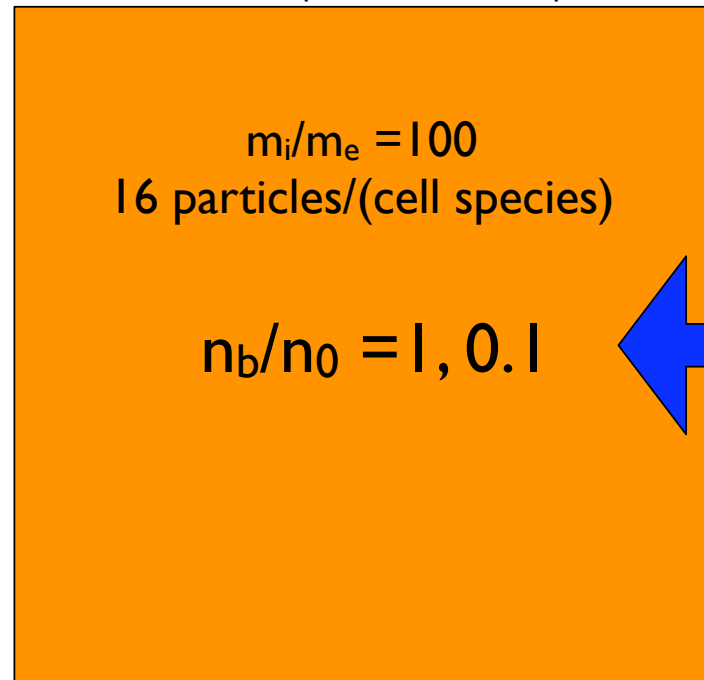
Power Law
 $n(\gamma) \propto \gamma^{-\alpha}$

2D PIC e⁻ i shells and particle acceleration ($\gamma_{\text{shell}} = 10$)

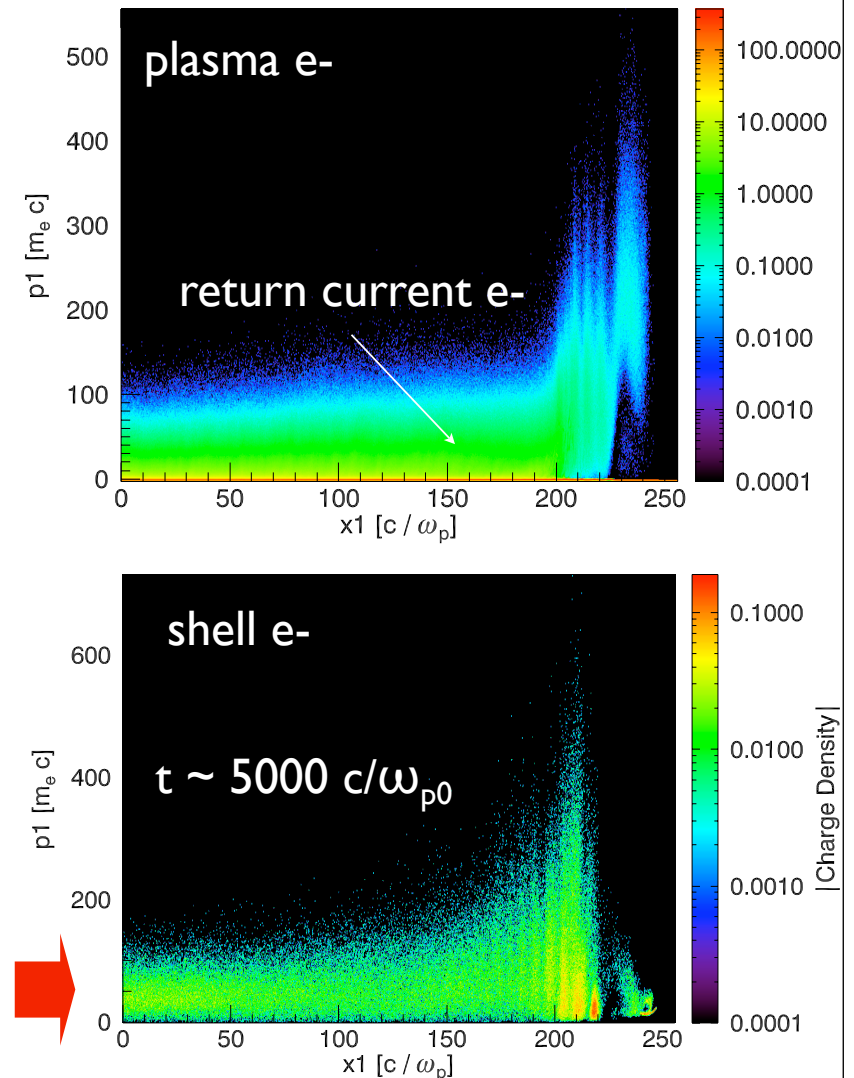
Moving window @ c

Relativistic ($\gamma = 10, 5$) plasma shell

$256 c/\omega_{p0} \times 256 c/\omega_{p0}$

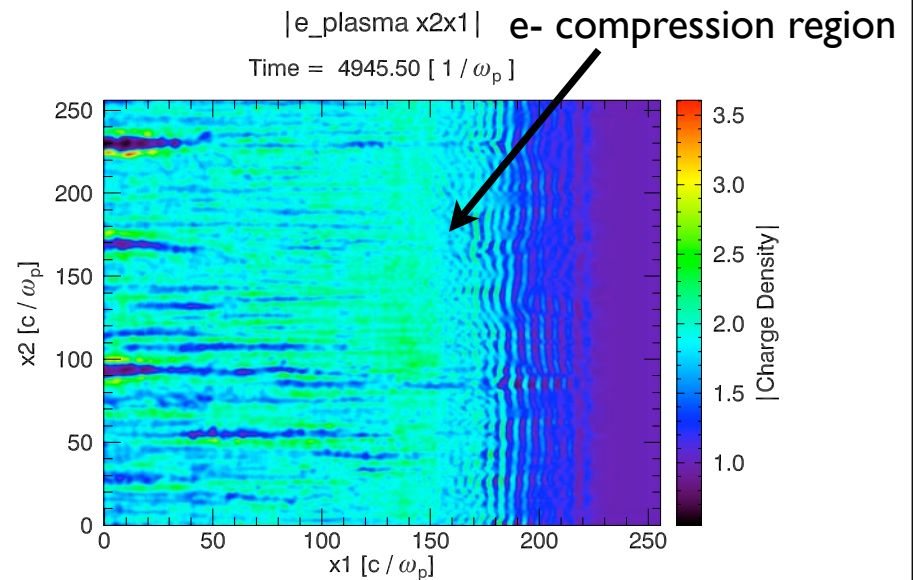
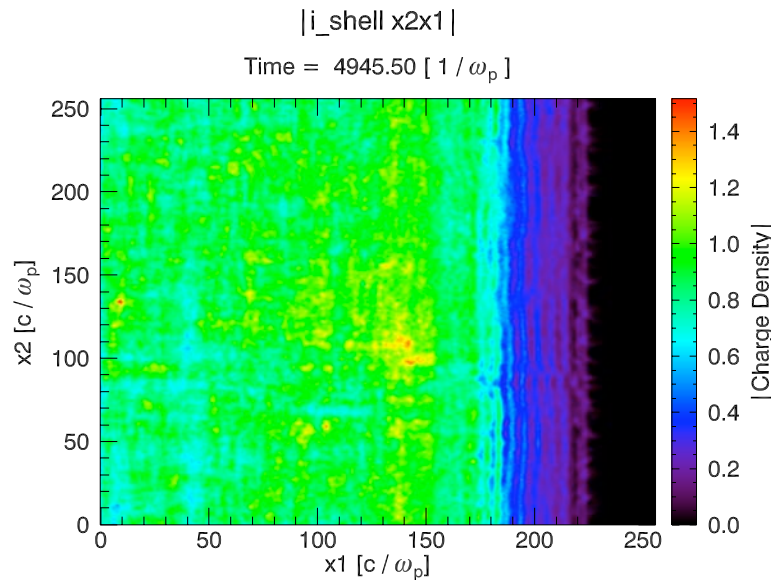


Strong erosion of shell e⁻ @ later times

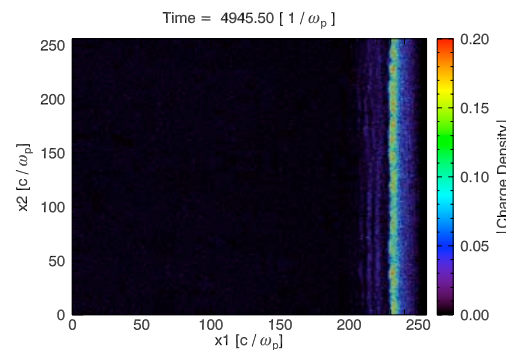


(Relativistic) Buneman responsible for accelerating structure

Strong electric field due to Buneman instability of ion shell with ambient background plasma



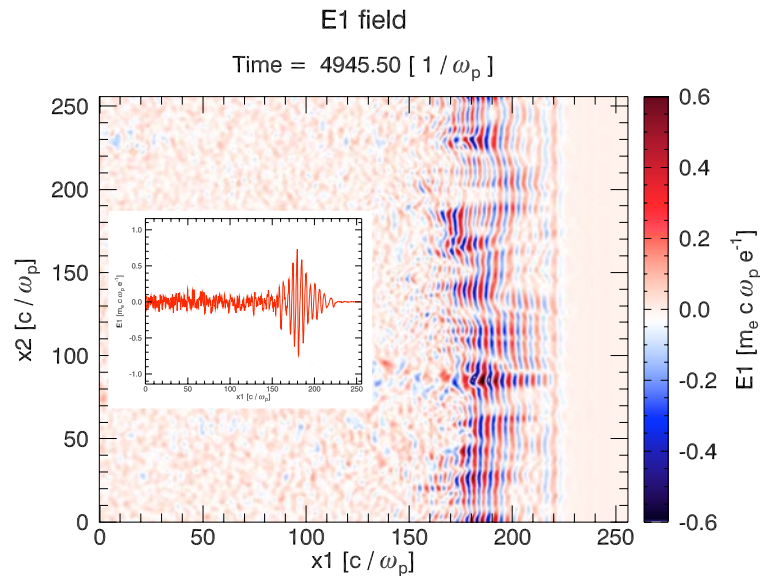
Small fraction of trapped electrons from the shell



$$k_{||} \approx \omega_{p0}/c$$

$$\Gamma \propto \sqrt{\frac{n_b}{n_0}} \sqrt{\frac{m_e}{m_i}}$$

Strong acceleration by parallel E field in leading edge of shell

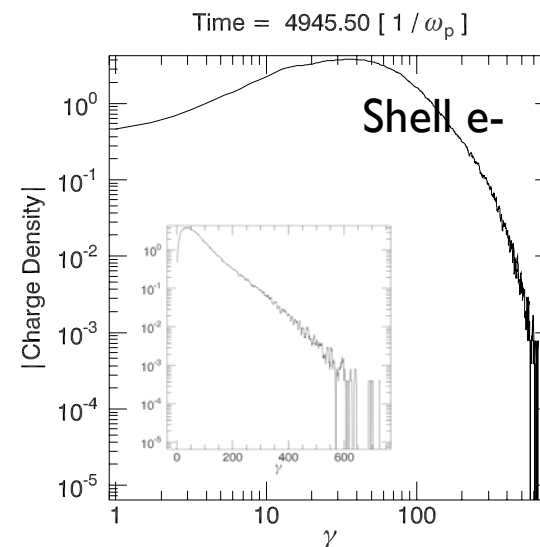
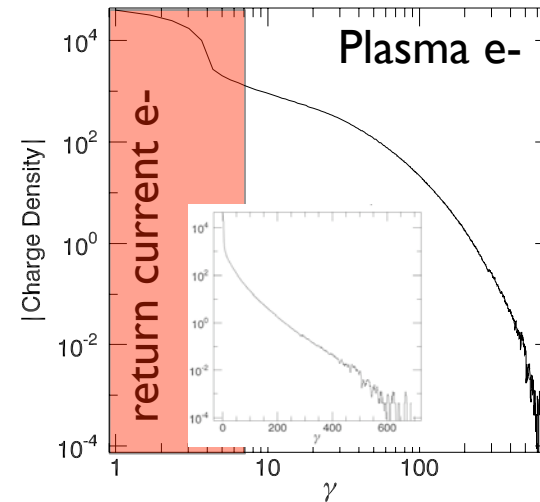


Accelerating structure
moving with $\Upsilon_\phi \sim \Upsilon_{\text{shell}} \sim 10$

$$\gamma_{\text{max}} \approx 4 \gamma_\phi^2 \tilde{E}_{\text{max}}$$

$\tilde{E}_{\text{max}} \approx 1 - 2$ (as measured in simulations, waves close to wavebreaking)

$$\tilde{E}_{\text{max}} = E_{\text{max}}/E_0 \approx E_{\text{max}}[V/cm]/\sqrt{n_0[\text{cm}^{-3}]}$$



Summary

- Particle-in-cell simulations as a tool to probe the microphysics of relativistic flows
- 3D simulations are very expensive but lower dimensional simulations can capture some of the physics
- Very large 2D simulations seem to confirm sub-equipartition B-field generation for very long times with ions ($\epsilon_B \sim 10^{-3}$, up to $2000/\omega_{p0}$) - in progress largest 2D simulations (4 x size)
- Dynamics of leading edge of relativistic shell
 - ≡ strong erosion of e- in shell
 - ≡ formation of hot return current with plasma e-
 - ≡ onset of Buneman instability leads to e- acceleration up to $\sim \gamma_{\text{shell}}^2$