

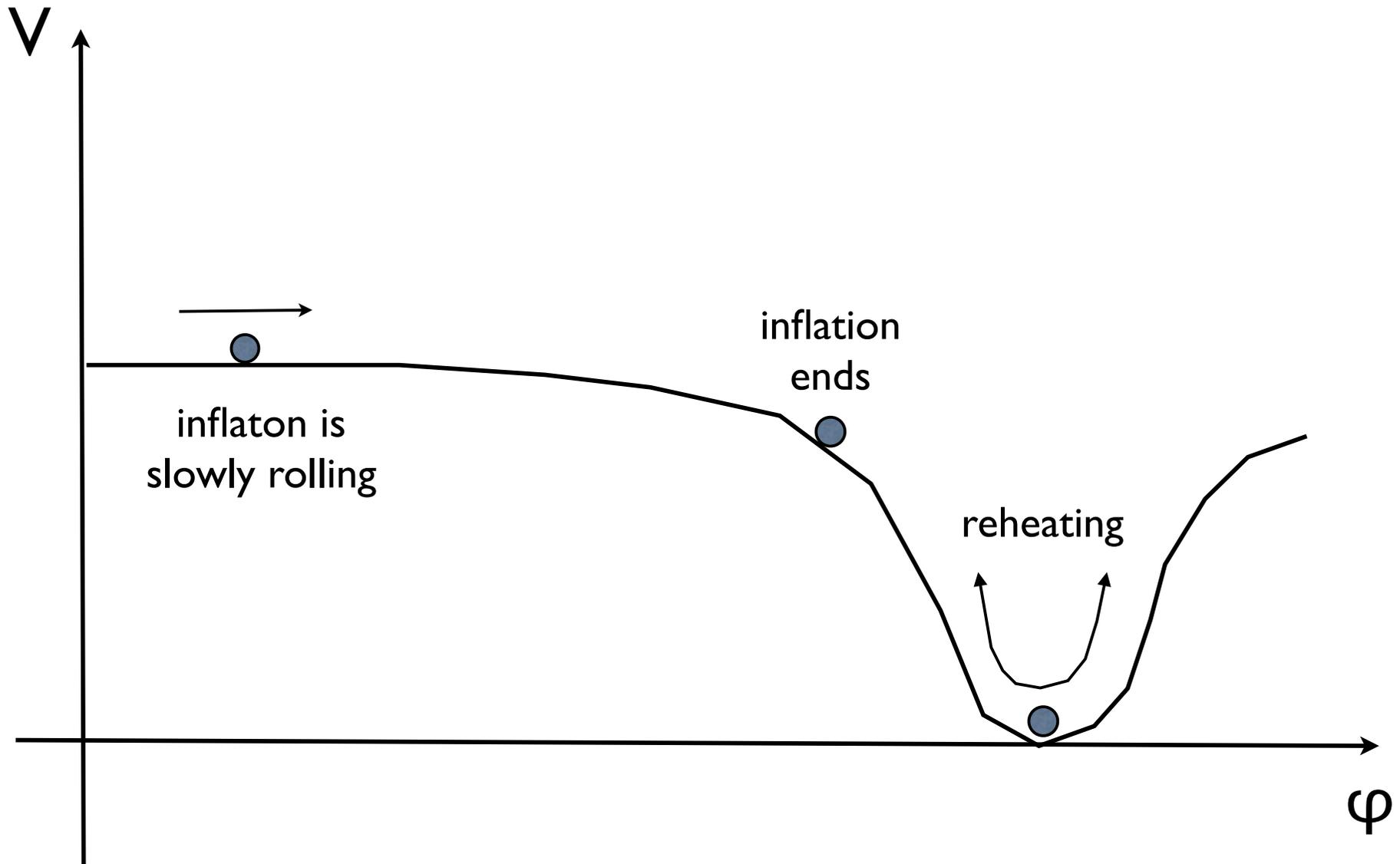
# The Quest for Primordial Non-Gaussianity

Overview and some recent developments

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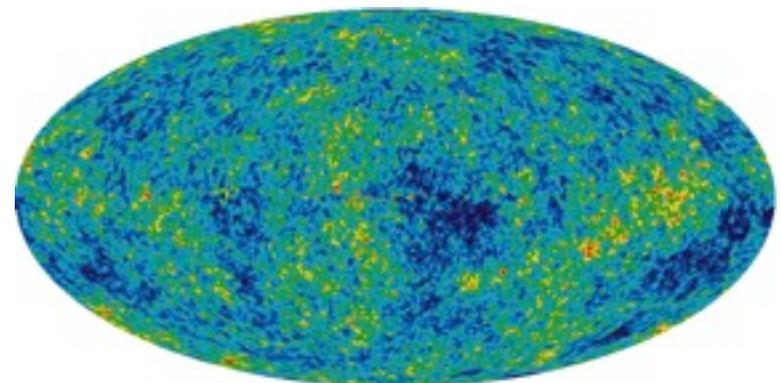
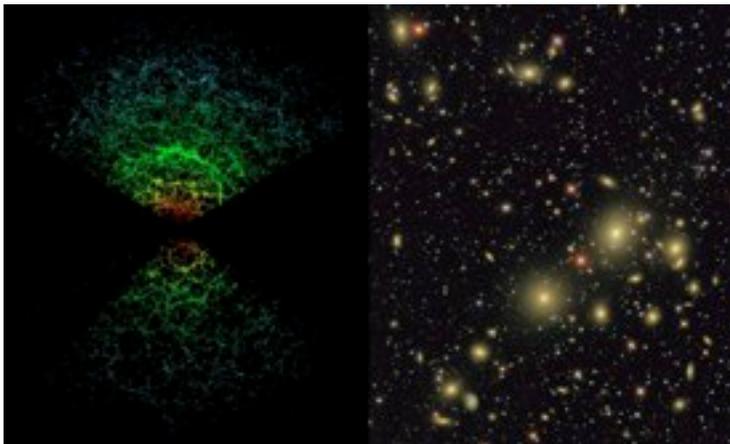
[On sabbatical at MPA and Excellence Cluster, Jan-Aug 2015]

# Motivation: testing Inflation



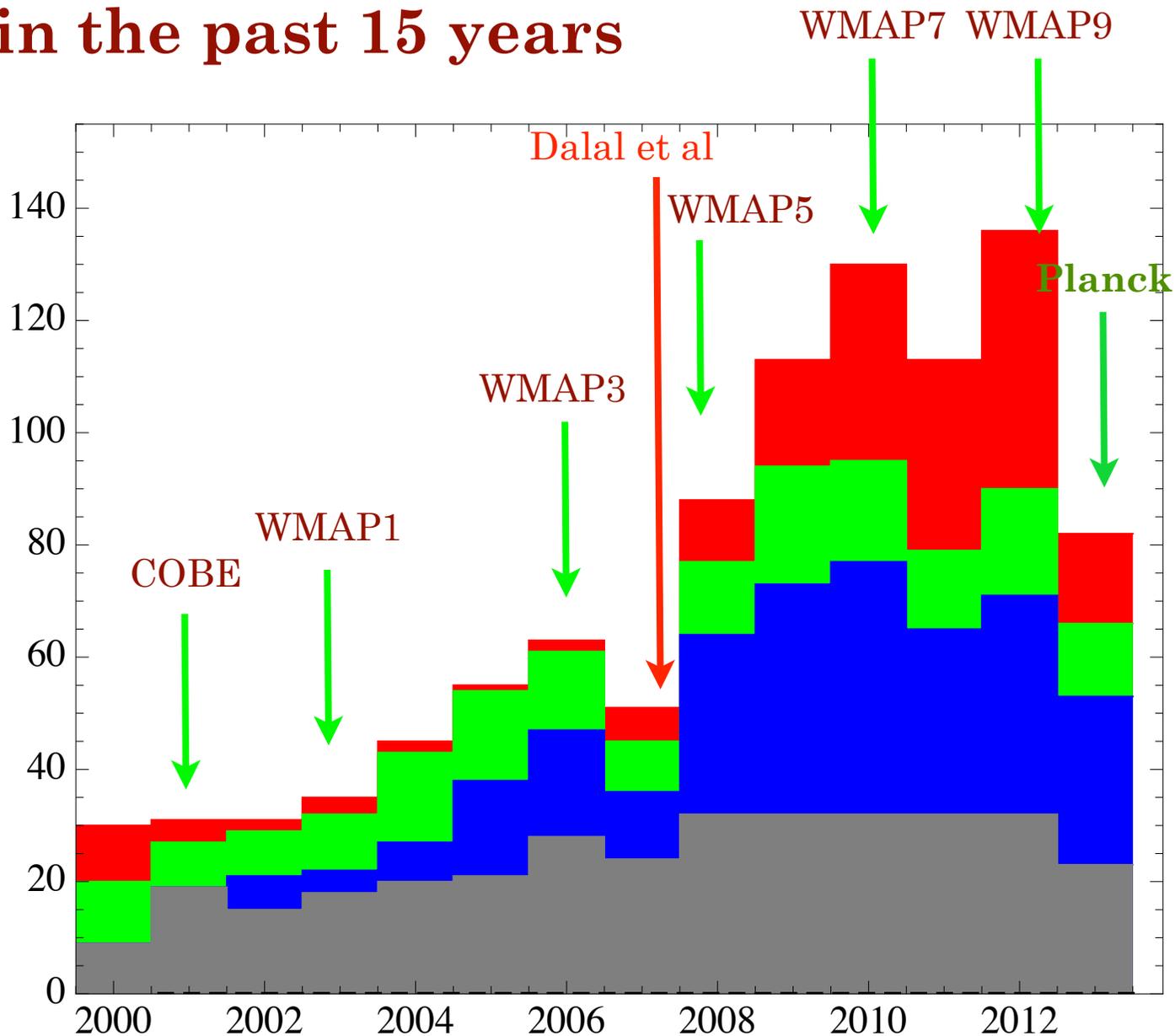
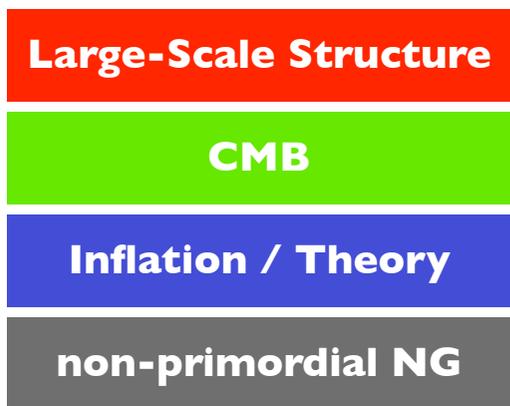
# Why study non-Gaussianity (NG)?

1. NG presents a window to the very early universe. For example, NG can distinguish between physically distinct models of inflation.
2. Conveniently, NG can be constrained/measured using CMB anisotropy maps and LSS. In particular, there is a rich set of observable quantities that are sensitive to primordial NG.



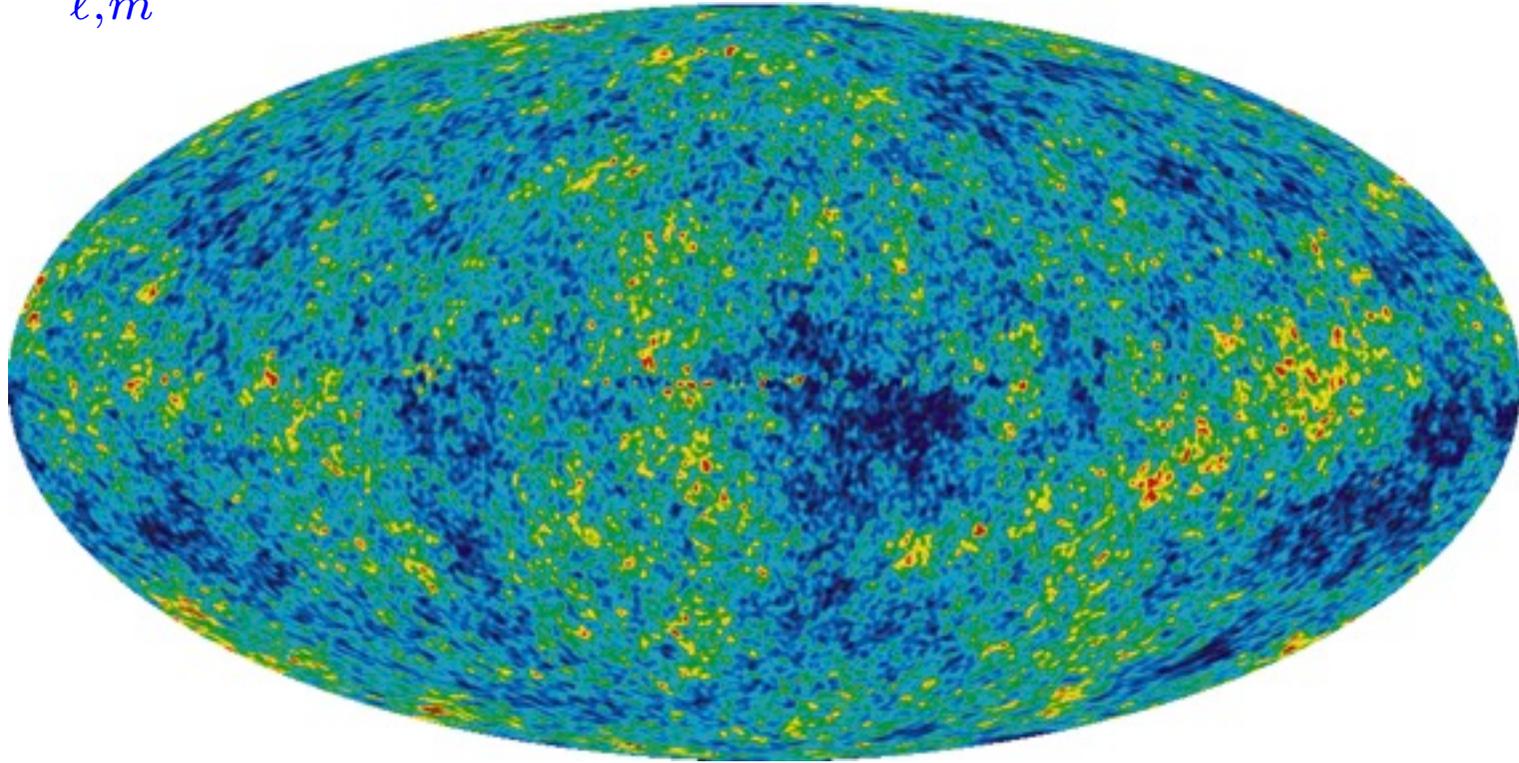
# Non-Gaussianity papers in the past 15 years

# of articles with  
“Non-Gaussian”  
in the title  
on the ADS data base



# Initial conditions in the universe

$$\frac{\delta T}{T}(\theta, \phi) = \sum_{\ell, m} a_{\ell m} Y_{\ell m}(\theta, \phi) \quad \ell \simeq \frac{180^\circ}{\theta}$$



**Generic** inflationary predictions: **Statistical Isotropy:**

- Flat geometry

$$\langle a_{\ell m} a_{\ell' m'} \rangle \equiv C_{\ell \ell' m m'} = C_\ell \delta_{\ell \ell'} \delta_{m m'}$$

- Nearly scale-invariant spectrum of density perturbations
- Background of gravity waves

**Gaussianity:**

- (Very nearly) gaussian initial conditions:  $\langle a_{\ell m} a_{\ell' m'} a_{\ell'' m''} \rangle = 0$

## Standard Inflation, with...

1. a single scalar field
2. the canonical kinetic term
3. always slow rolls
4. in Bunch-Davies vacuum
5. in Einstein gravity

produces **unobservable** NG

Therefore, measurement of nonzero NG would point to a **violation** of one of the assumptions above

# Recall: power spectrum

Define Fourier transform of density fluctuation:  $\delta(\vec{r}) = \int \frac{d^3 k}{(2\pi)^3} e^{-i\vec{k}\vec{r}} \delta_{\vec{k}}$

Then the power spectrum  $P(k)$  is defined via

$$\langle \delta_{\vec{k}_1} \delta_{\vec{k}_2}^* \rangle = (2\pi)^3 \delta^{(3)}(\vec{k}_1 + \vec{k}_2) P(k)$$

Sometimes it's nice to work in harmonic space

$$a_{\ell m} = 4\pi(-i)^\ell \int \frac{d^3 k}{(2\pi)^3} T_\ell(k) \delta(\vec{k}) Y_{\ell m}(\hat{k})$$

Then the angular power spectrum is defined as:

$$\langle a_{\ell m} a_{\ell' m'} \rangle = \delta_{\ell\ell'} \delta_{mm'} C_\ell$$

The bispectrum:  
similar, but for 3-pt function

Fourier space:

$$\langle \delta_{\vec{k}_1} \delta_{\vec{k}_2} \delta_{\vec{k}_3} \rangle = (2\pi)^3 \delta^{(3)}(\vec{k}_1 + \vec{k}_2 + \vec{k}_3) B(\vec{k}_1, \vec{k}_2, \vec{k}_3)$$

Harmonic space:

$$\langle a_{\ell_1 m_1} a_{\ell_2 m_2} a_{\ell_3 m_3} \rangle \equiv B_{\ell_1 \ell_2 \ell_3}^{m_1 m_2 m_3}$$

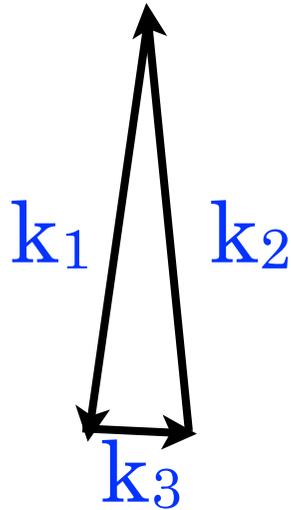
and the angle-averaged bispectrum is

$$B_{\ell_1 \ell_2 \ell_3} \equiv \sqrt{\frac{(2\ell_1 + 1)(2\ell_2 + 1)(2\ell_3 + 1)}{4\pi}} \begin{pmatrix} \ell_1 & \ell_2 & \ell_3 \\ 0 & 0 & 0 \end{pmatrix} \sum_{m_1 m_2 m_3} \begin{pmatrix} \ell_1 & \ell_2 & \ell_3 \\ m_1 & m_2 & m_3 \end{pmatrix} B_{\ell_1 \ell_2 \ell_3}^{m_1 m_2 m_3}$$

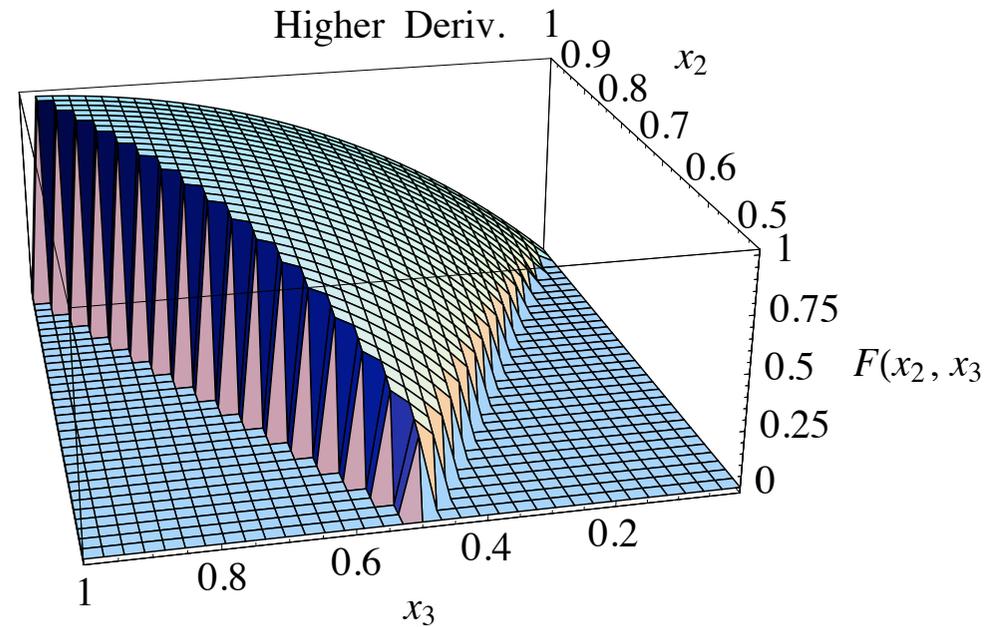
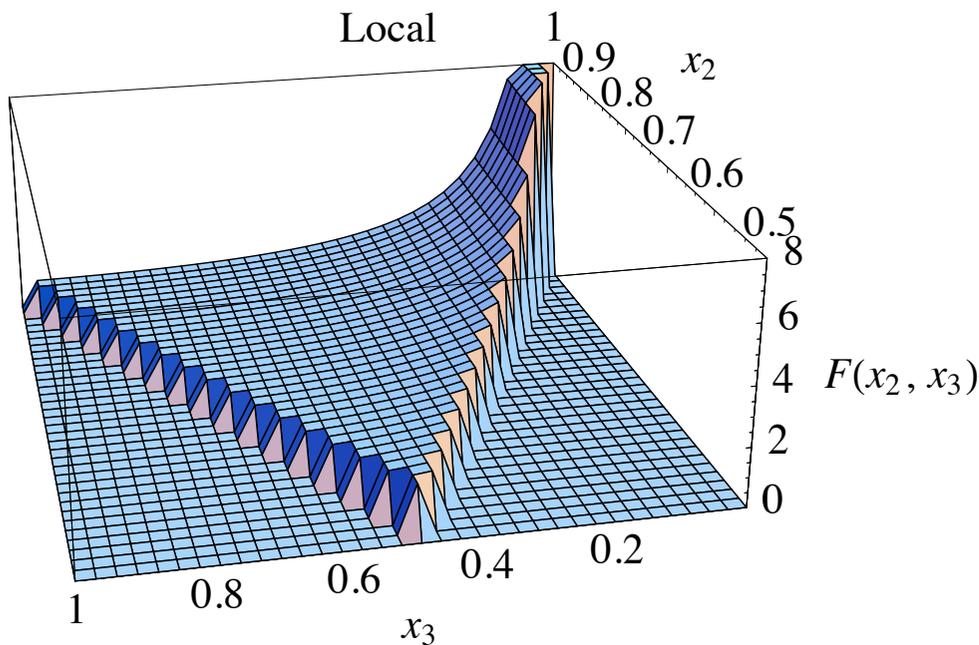
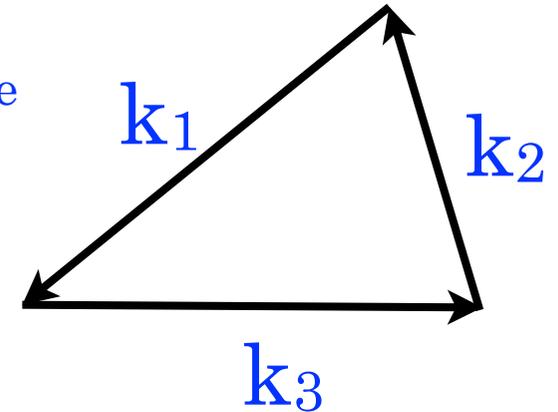
# 3-pt correlation function of CMB anisotropy ⇒ direct window into inflation

e.g. Luo & Schramm 1993

“local”  
(eg. multi-field)

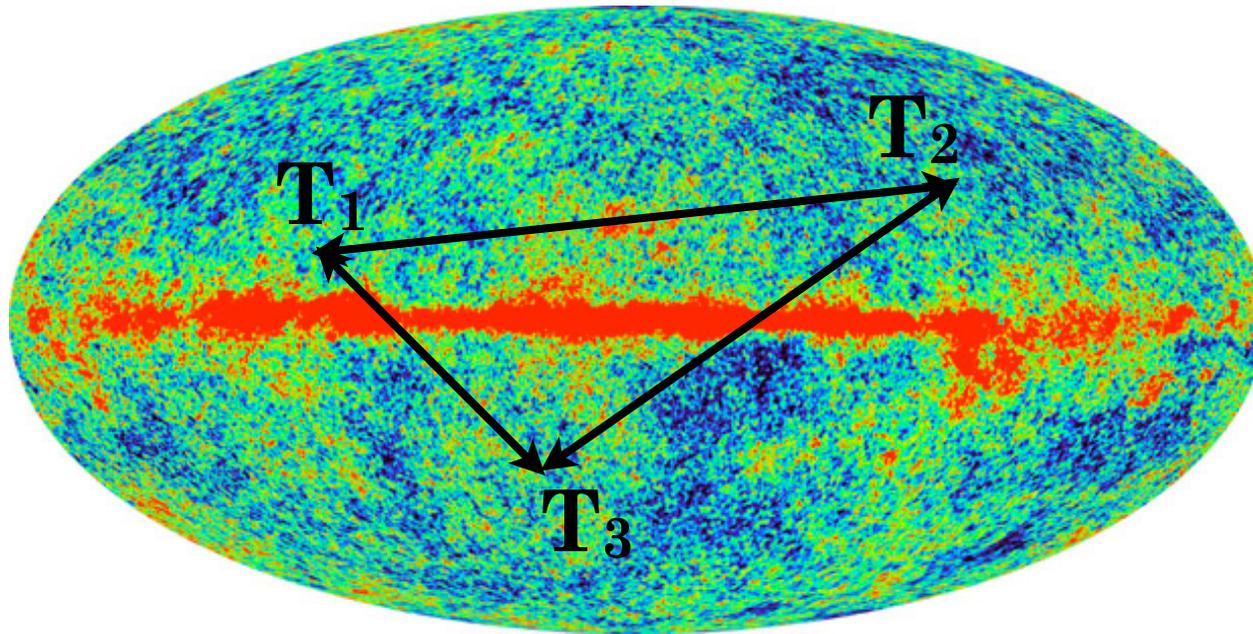


“equilateral”  
(eg. higher-derivative  
action; interactions)



Babich, Creminelli & Zaldarriaga (2004)

# NG from 3-point correlation function

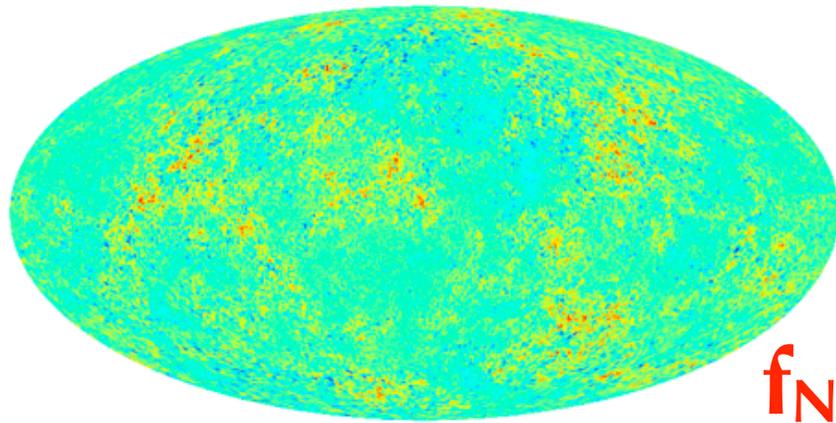


Local NG (squeezed triangles) - tests # inflationary fields

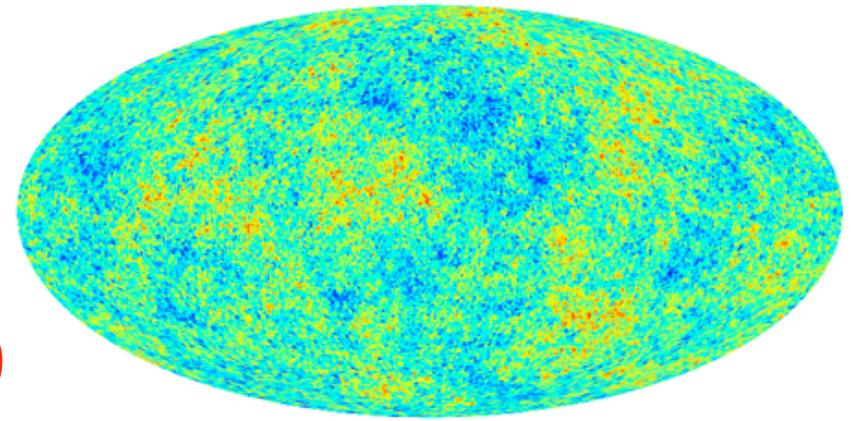
$$\Phi = \Phi_G + f_{\text{NL}} (\Phi_G^2 - \langle \Phi_G^2 \rangle)$$

“Equilateral”, “orthogonal” NG- tests inflationary interactions  
tests interactions; parameter  $f_{\text{NL}}^{\text{eq}}$ ,  $f_{\text{NL}}^{\text{orth}}$

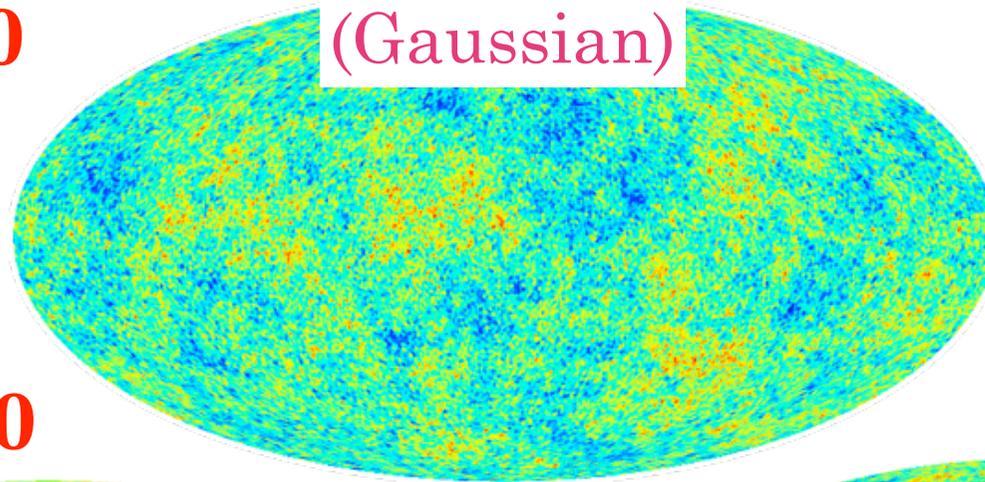
Threshold for new physics:  $f_{\text{NL}}^{\text{any kind}} \gtrsim \text{O}(1)$



$f_{\text{NL}} = -5000$

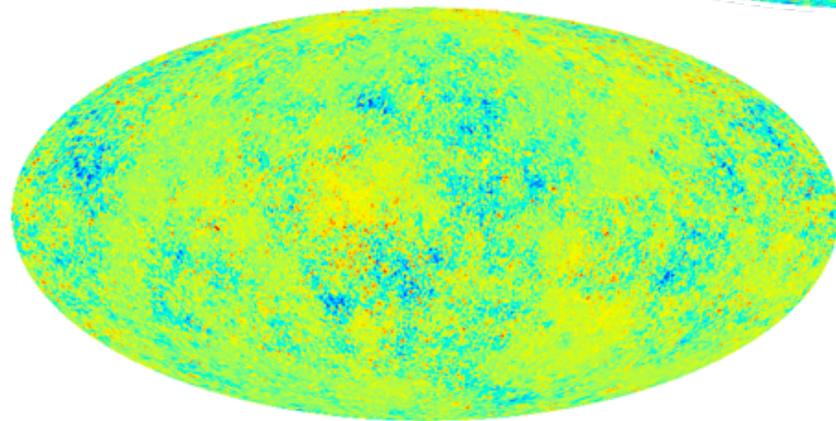


$f_{\text{NL}} = -500$

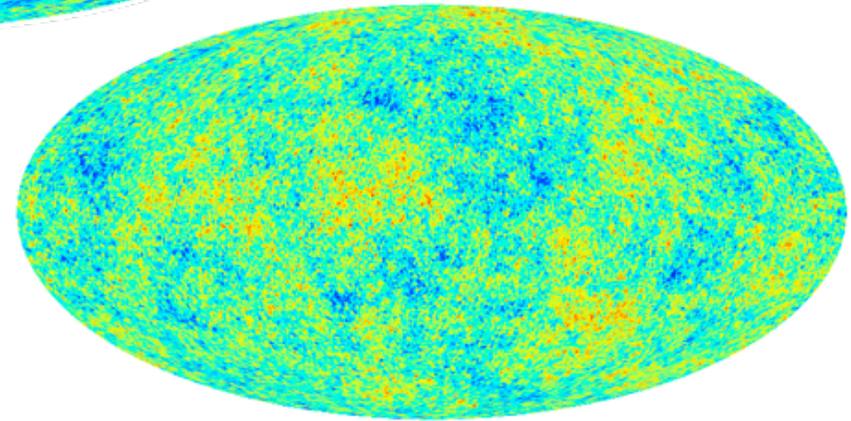


$f_{\text{NL}} = 0$   
(Gaussian)

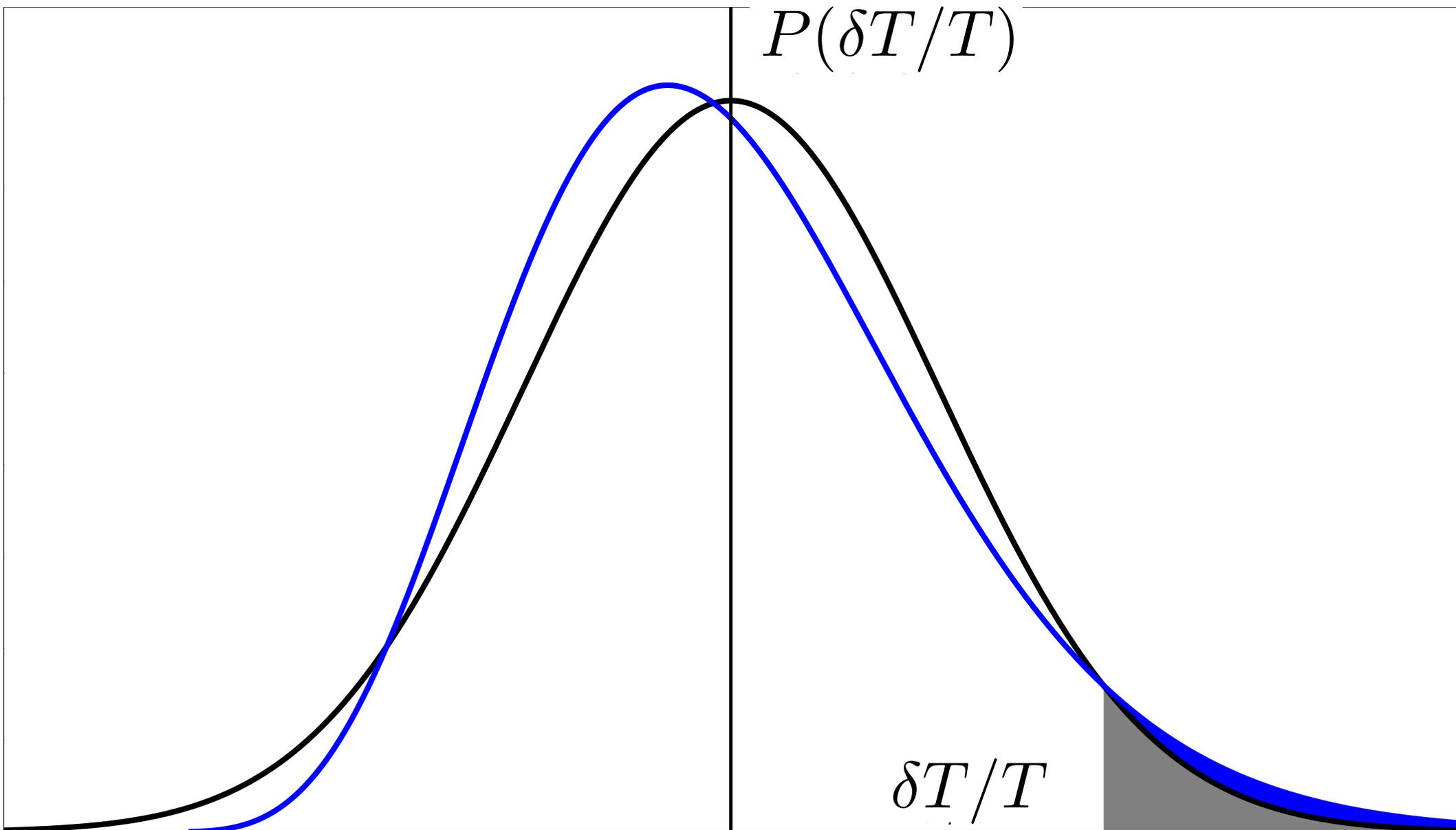
$f_{\text{NL}} = +5000$



$f_{\text{NL}} = +500$



Current upper bound on NG is  
~1000 times smaller than **this**:



# Brief history of NG measurements: 1990's

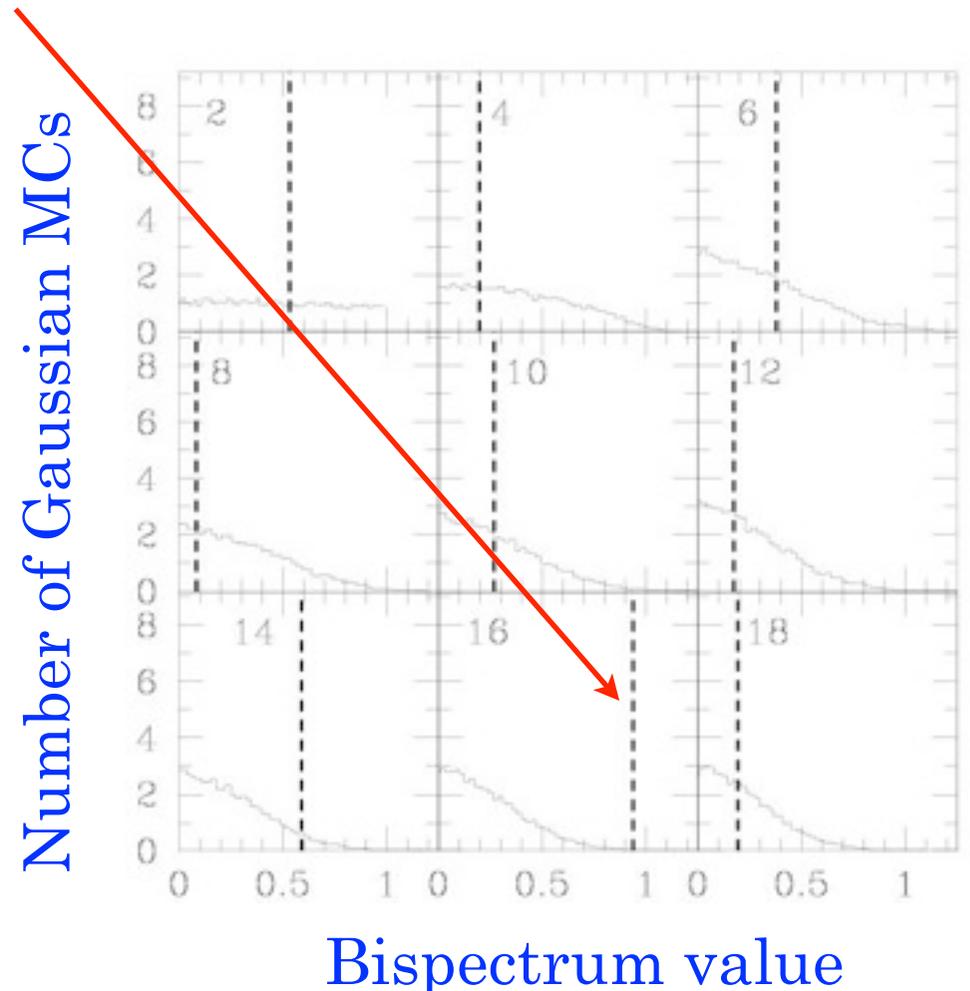
Early 1990s; COBE: Gaussian CMB sky (Kogut et al 1996)

$|f_{\text{NL}}| \approx 3000$  (Komatsu 2002)

1998; COBE: claim of NG at  $l=16$  equilateral bispectrum (Ferreira, Magueijo & Gorski 1998)

but explained by a known systematic effect! (Banday, Zaroubi & Gorski 1999)

(and anyway isn't unexpected given all bispectrum configurations you can measure; Komatsu 2002)



# Brief history of NG measurements: 2000's

Pre-WMAP CMB: all is gaussian (e.g. MAXIMA; Wu et al 2001)

WMAP pre-2008: all is gaussian

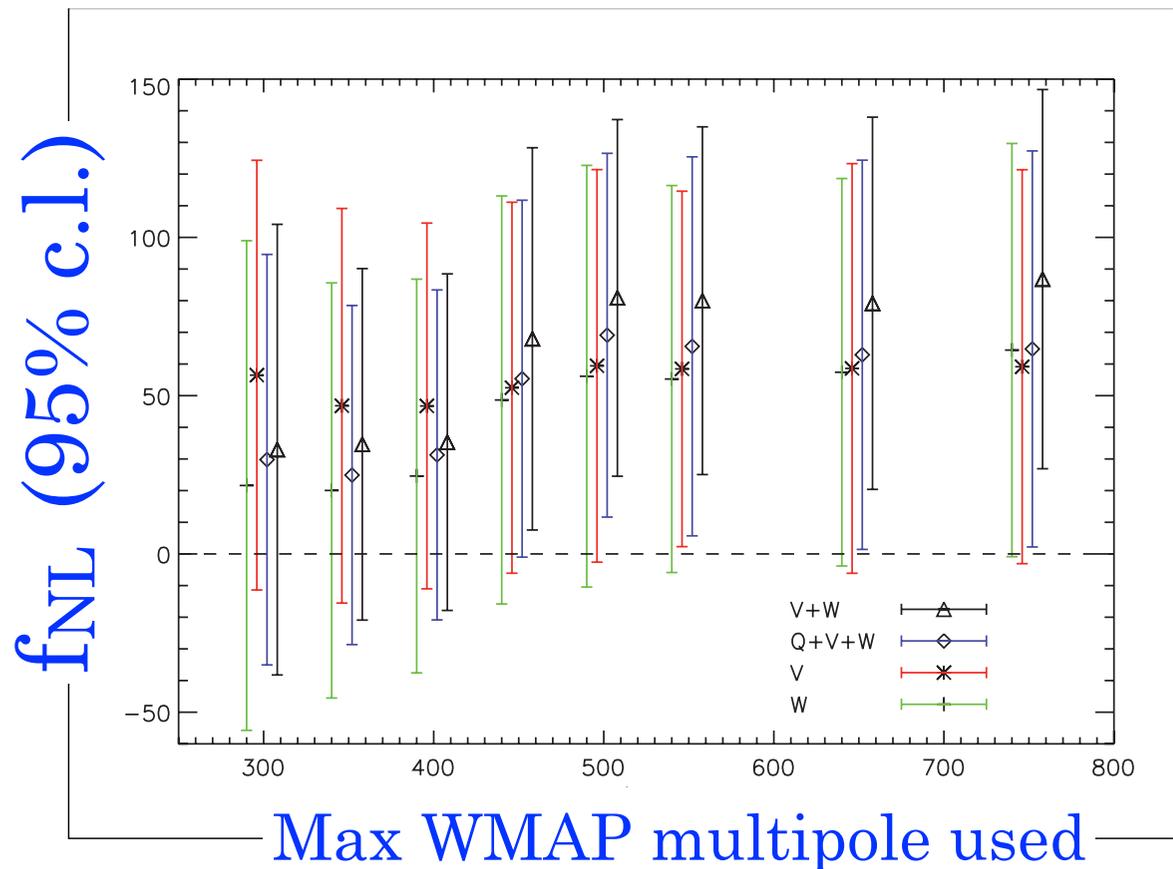
(Komatsu et al. 2003; Creminelli, Senatore, Zaldarriaga & Tegmark 2007)

$$-36 < f_{\text{NL}} < 100 \quad (95\% \text{ CL})$$

Dec 2007, claim of NG in WMAP

(Yadav & Wandelt arXiv:0712.1148)

$$27 < f_{\text{NL}} < 147 \quad (95\% \text{ CL})$$



# Constraints from WMAP (7-yr)

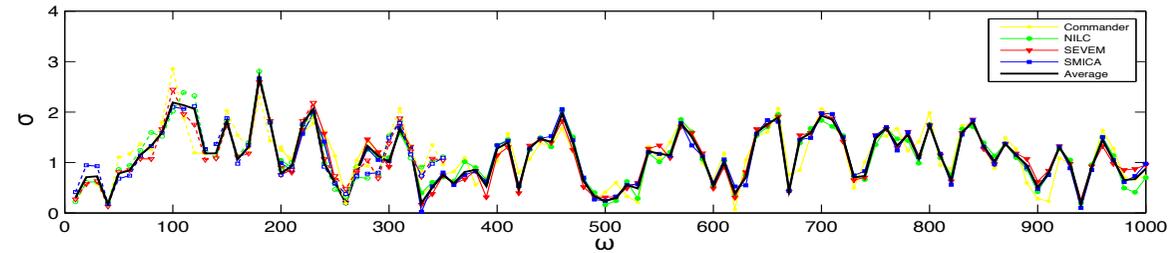
Band	Foreground <sup>b</sup>	$f_{NL}^{\text{local}}$	$f_{NL}^{\text{equil}}$	$f_{NL}^{\text{orthog}}$	$b_{src}$
V+W	Raw	$59 \pm 21$	$33 \pm 140$	$-199 \pm 104$	N/A
V+W	Clean	$42 \pm 21$	$29 \pm 140$	$-198 \pm 104$	N/A
V+W	Marg. <sup>c</sup>	$32 \pm 21$	$26 \pm 140$	$-202 \pm 104$	$-0.08 \pm 0.12$
V	Marg.	$43 \pm 24$	$64 \pm 150$	$-98 \pm 115$	$0.32 \pm 0.23$
W	Marg.	$39 \pm 24$	$36 \pm 154$	$-257 \pm 117$	$-0.13 \pm 0.19$

Komatsu et al. 2011

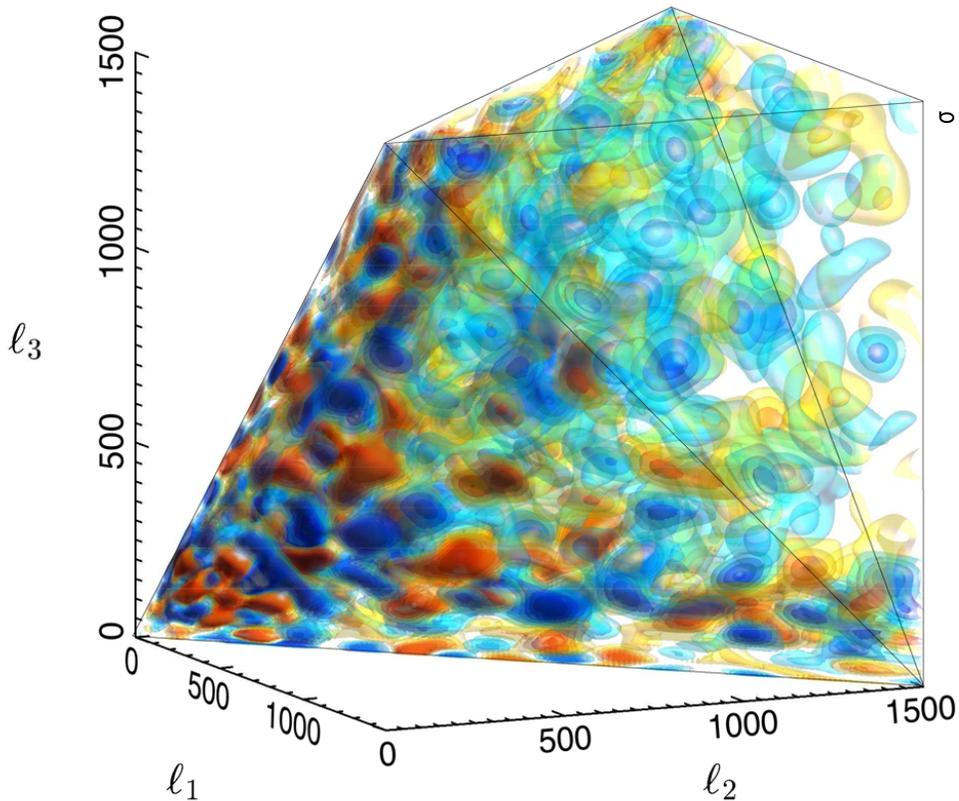
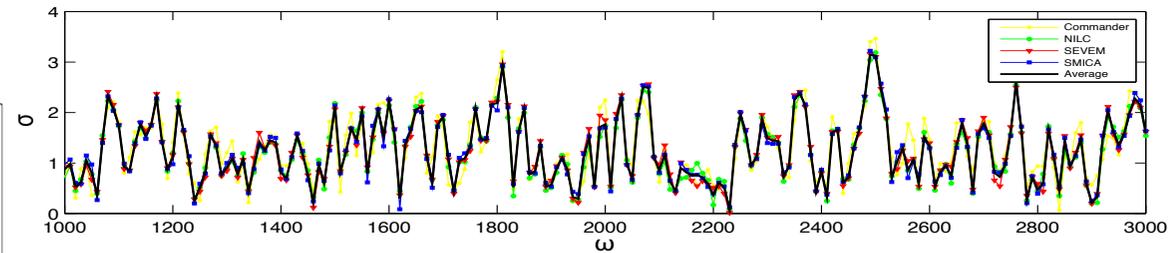
# Constraints from Planck

Shape and method	$f_{\text{NL}}(\text{KSW})$	
	Independent	ISW-lensing subtracted
SMICA ( $T$ )		
Local . . . . .	10.2 $\pm$ 5.7	<b>2.5 <math>\pm</math> 5.7</b>
Equilateral . . . . .	-13 $\pm$ 70	<b>-16 <math>\pm</math> 70</b>
Orthogonal . . . . .	-56 $\pm$ 33	<b>-34 <math>\pm</math> 33</b>
SMICA ( $T+E$ )		
Local . . . . .	6.5 $\pm$ 5.0	<b>0.8 <math>\pm</math> 5.0</b>
Equilateral . . . . .	3 $\pm$ 43	<b>-4 <math>\pm</math> 43</b>
Orthogonal . . . . .	-36 $\pm$ 21	<b>-26 <math>\pm</math> 21</b>

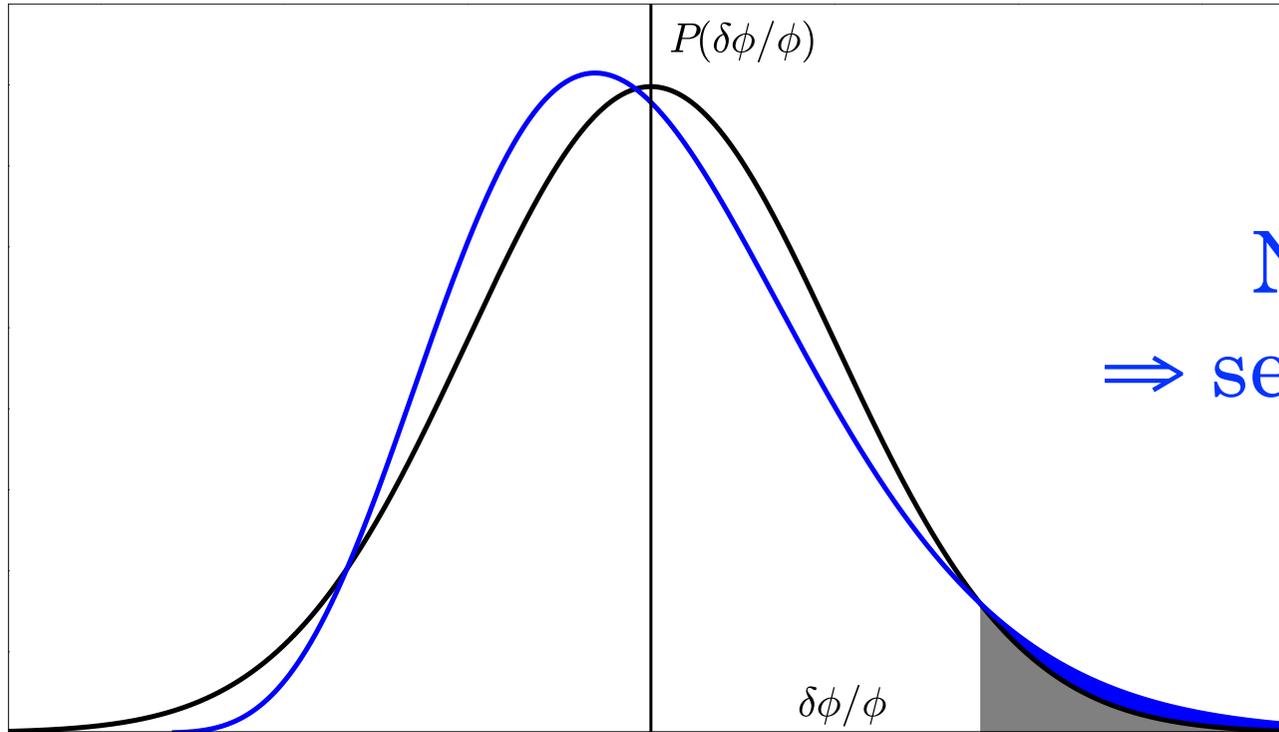
# Constraints from Planck: modal expansion



$$B(k_1, k_2, k_3) = \sum_{p,r,s} \alpha_{prs} q_p(k_1) q_r(k_2) q_s(k_3)$$



# Galaxy cluster counts' sensitivity to NG



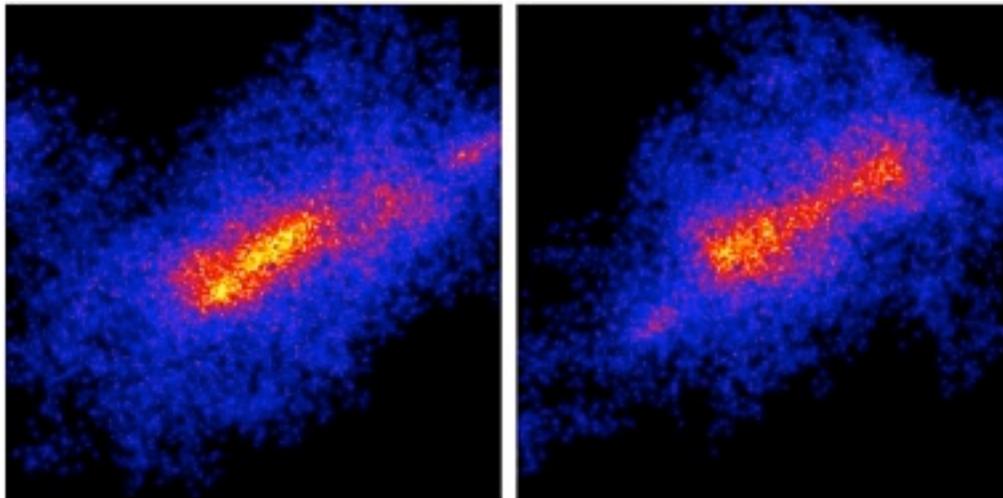
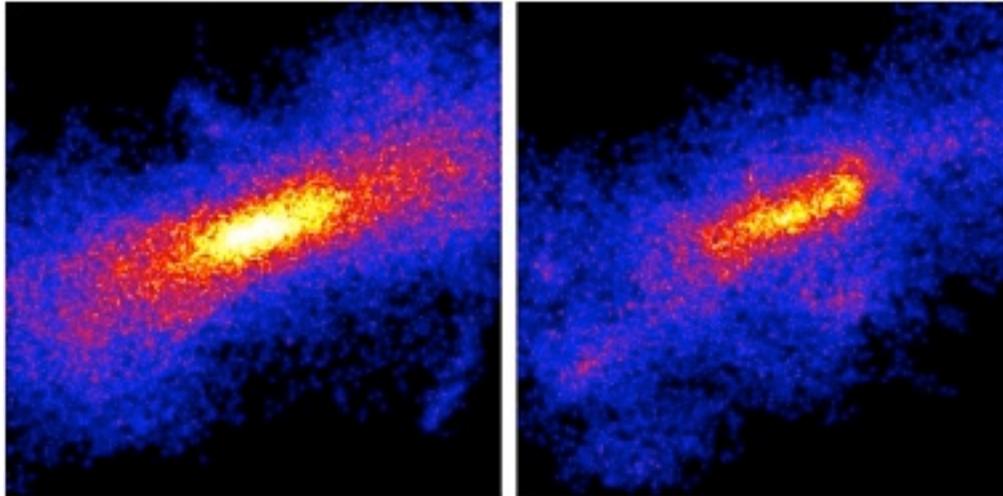
NG initial PDF  
⇒ sensitivity to counts  
“on the tail”

Lots of effort in the community to calibrate  
the **non-Gaussian mass function** -  
 $dn/d\ln M(M, z)$  - of DM halos

# A DM halo gets more massive with $f_{\text{NL}} > 0$ (and v.v.)

$f_{\text{NL}} = +5000$   
 $M = 1.2 \cdot 10^{16} M_{\odot}$

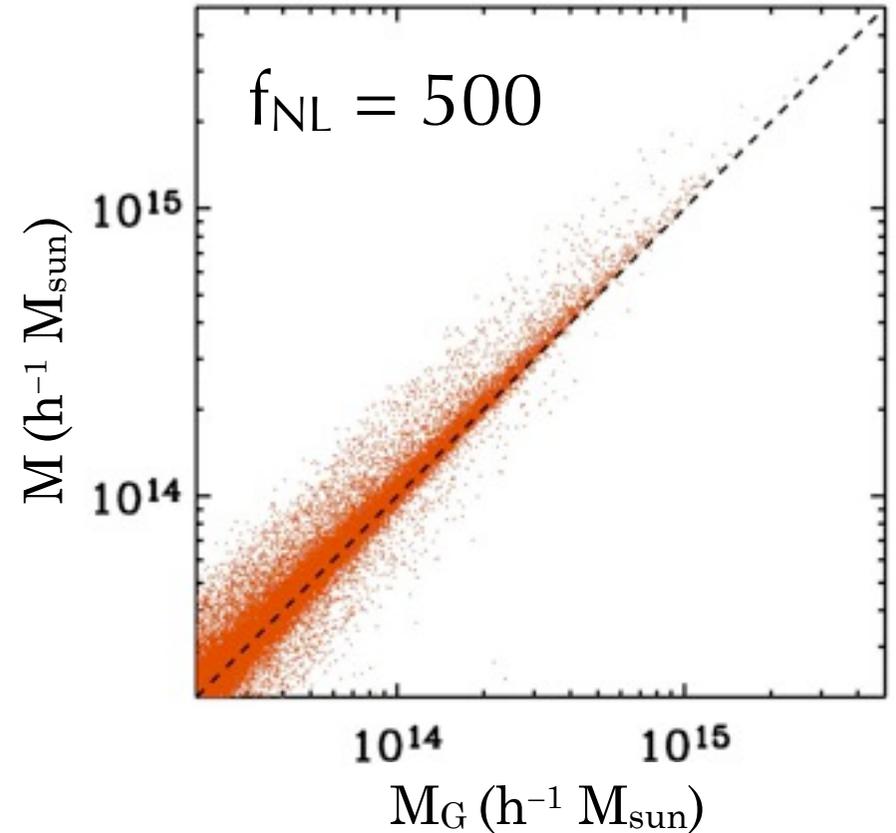
$f_{\text{NL}} = +500$   
 $M = 5.9 \cdot 10^{15} M_{\odot}$



$f_{\text{NL}} = 0$   
 $M = 5.1 \cdot 10^{15} M_{\odot}$

$f_{\text{NL}} = -500$   
 $M = 4.3 \cdot 10^{15} M_{\odot}$

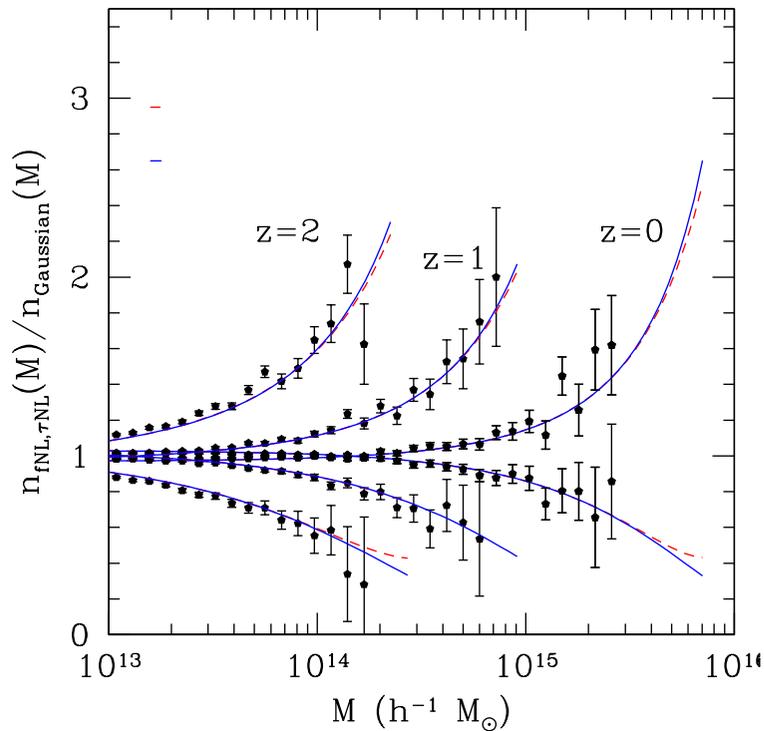
Mapping between  
 $M_{\text{G}}$  and  $M \equiv M_{\text{NG}}$ :



$\Rightarrow$  NG mass function:

$$\frac{dN}{dM} = \int \frac{dP(M|M_{\text{G}})}{dM} \frac{dN}{dM_{\text{G}}} dM_{\text{G}}$$

NG/Gaussian **mass function** ratios:  
for fixed  $M$ , more sensitivity  
at higher redshift



Smith & LoVerde 2011; many others going back to 1990s

Unfortunately, cluster counts are **weakly**  
sensitive to NG

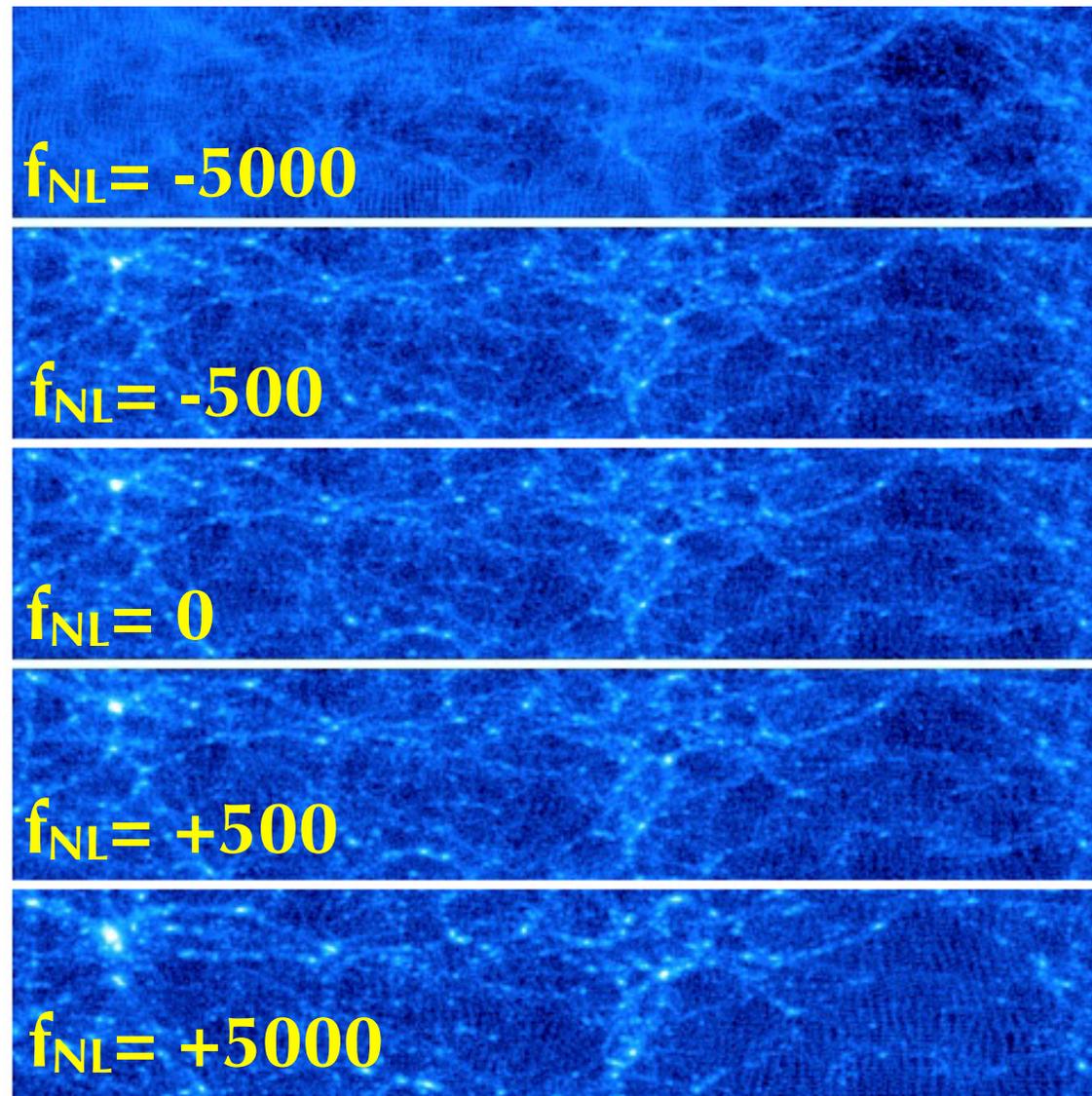
e.g.  $\sigma(f_{NL})=450$  measured from SPT (Williamson et al 2010)

Nevertheless:

- cluster abundance is sensitive to ALL non-Gaussianity

# **Effects of primordial NG on the bias of virialized objects**

# Simulations with non-Gaussianity ( $f_{\text{NL}}$ )



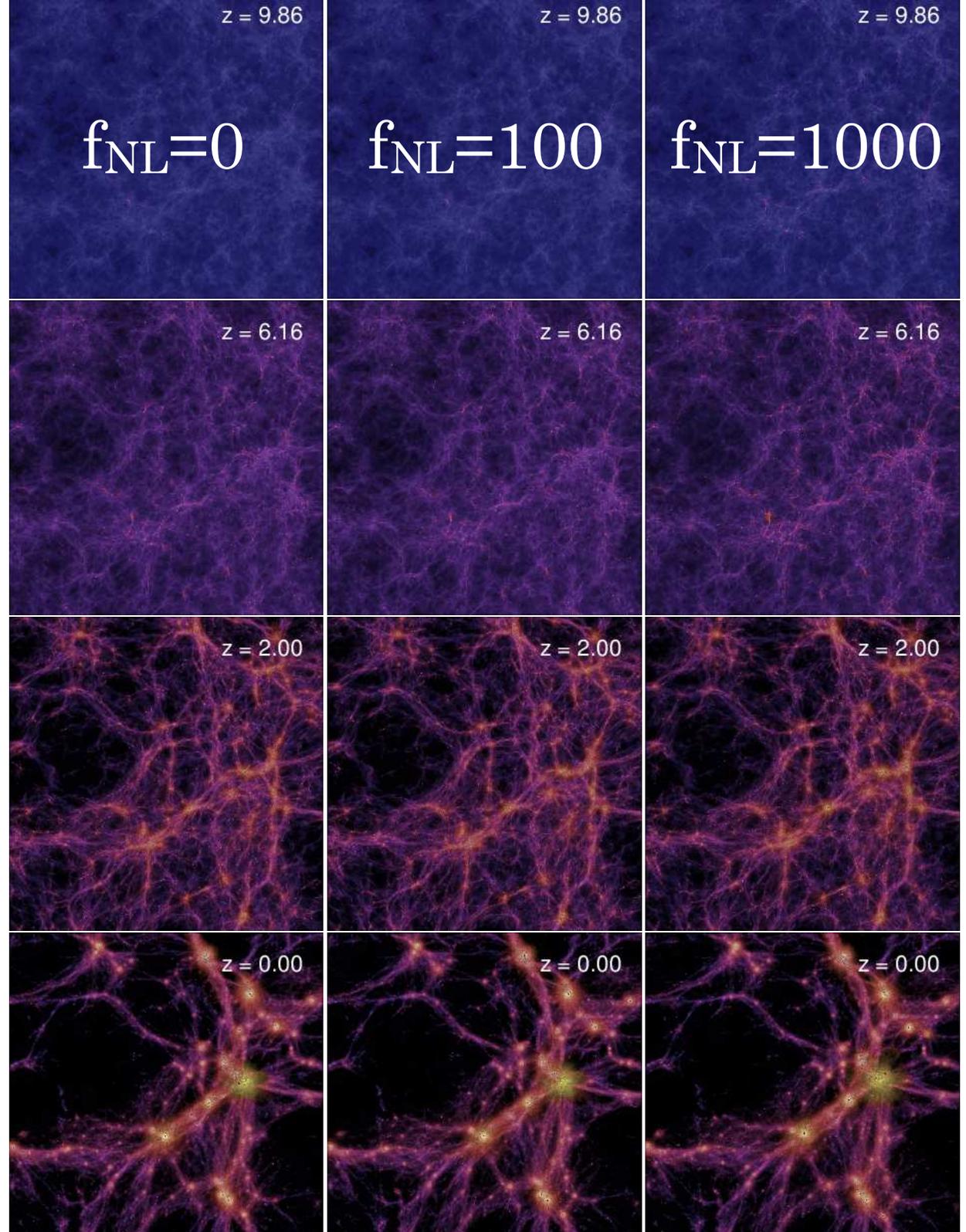
- Under-dense region evolution decrease with  $f_{\text{NL}}$
- Over-dense region evolution increase with  $f_{\text{NL}}$

80 Mpc/h

375 Mpc/h

- Same initial conditions, different  $f_{\text{NL}}$
- Slice through a box in a simulation  $N_{\text{part}}=512^3$ ,  $L=800$  Mpc/h

...and now  
with baryons!



Zhao, Li,  
Shandera & Jeong,  
arXiv:1307.5051

# Does galaxy/halo bias depend on NG?

$$\text{bias} \equiv \frac{\text{clustering of galaxies}}{\text{clustering of dark matter}} = \frac{\left(\frac{\delta\rho}{\rho}\right)_{\text{halos}}}{\left(\frac{\delta\rho}{\rho}\right)_{\text{DM}}}$$

cosmologists measure

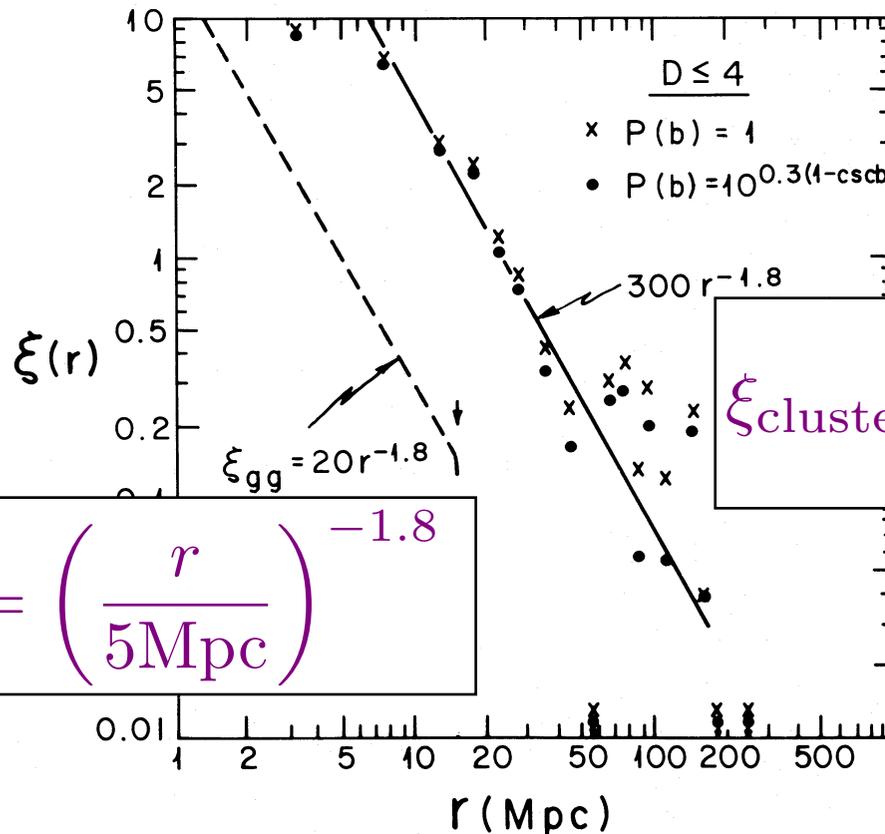
DM

↓

usually nuisance parameter(s)

↓

theory predicts

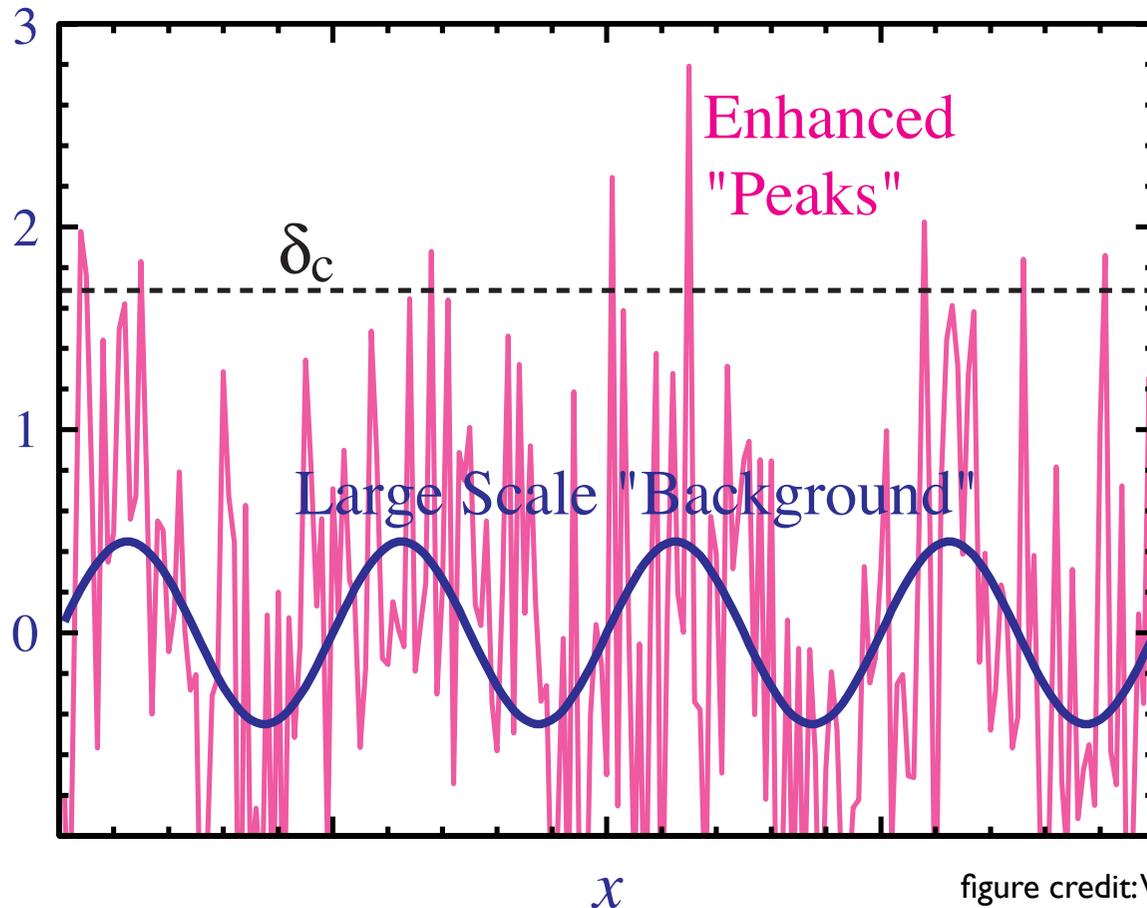


$$\xi_{\text{galaxies}}(r) = \left(\frac{r}{5\text{Mpc}}\right)^{-1.8}$$

$$\xi_{\text{clusters}}(r) = \left(\frac{r}{25\text{Mpc}}\right)^{-1.8}$$

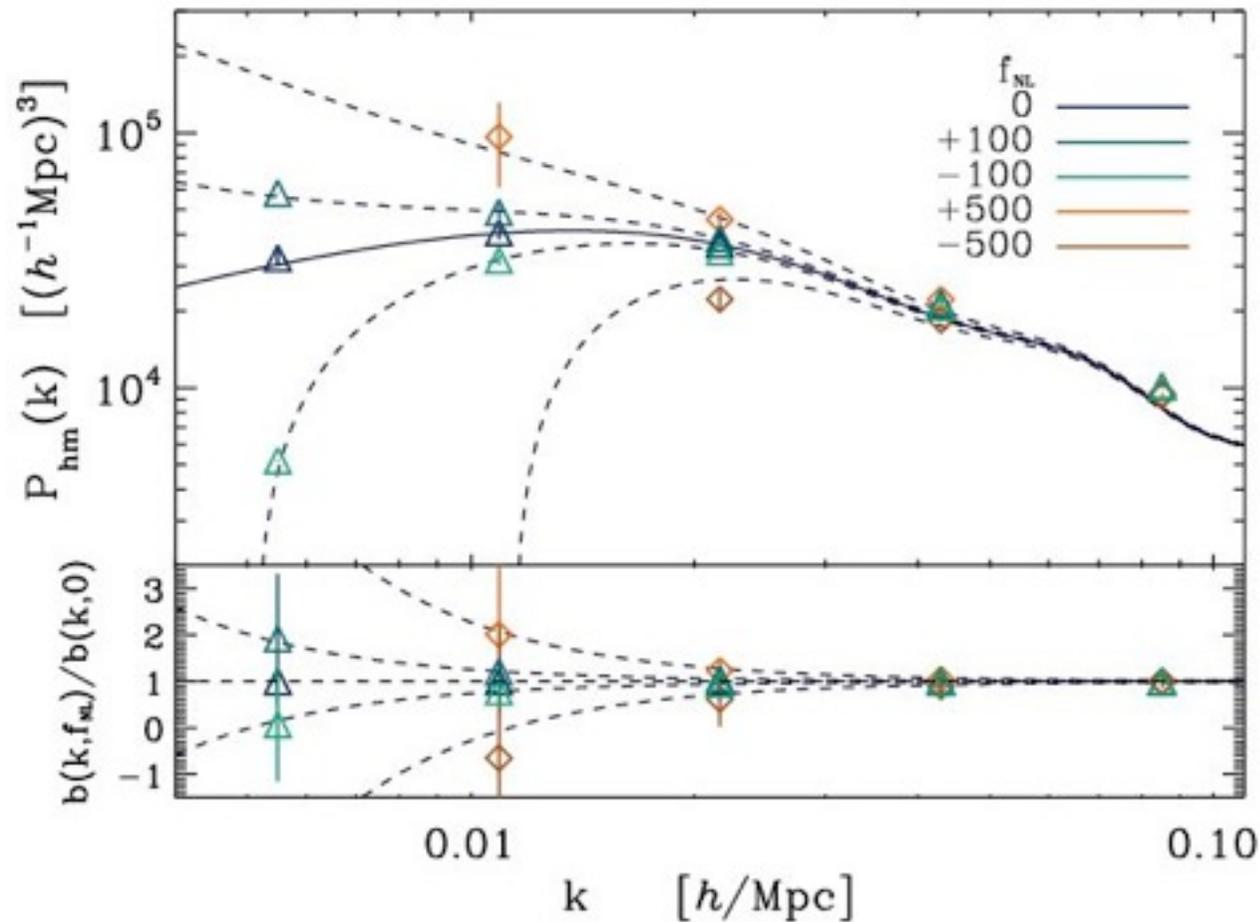
# Bias of dark matter halos

$$P_h(k, z) = b^2(k, z) P_{\text{DM}}(k, z)$$



Simulations and theory both say: **large-scale bias is scale-independent**  
(theorem if halo abundance is function of local density  
**and if the short and long modes are uncorrelated**)

# Scale dependence of NG halo bias



$$b(k) = b_{\text{G}} + f_{\text{NL}} \frac{\text{const}}{k^2}$$

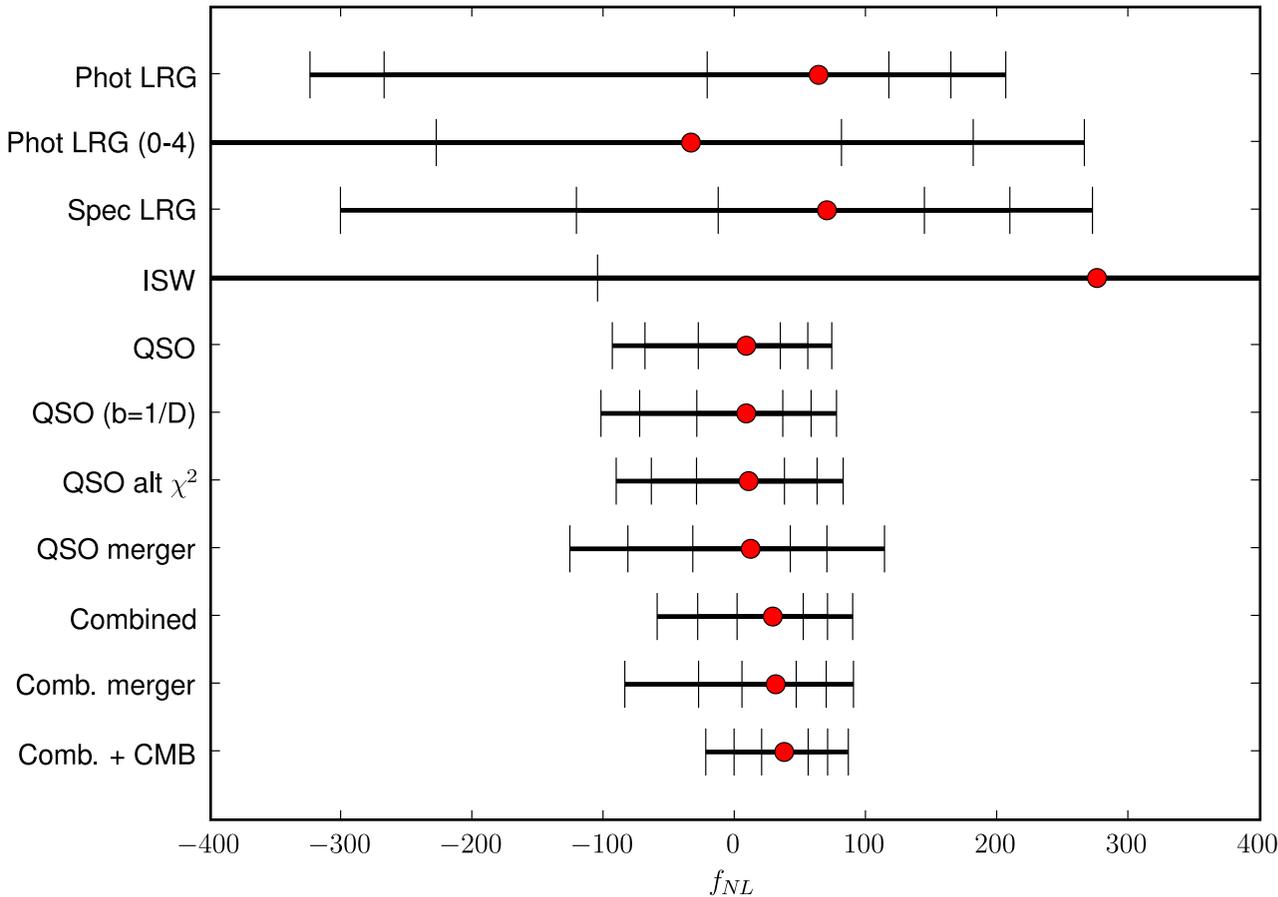
Verified using a variety of theoretical derivations and numerical simulations.

$$\Delta b(k) = f_{\text{NL}}(b_G - 1) \delta_c \frac{3 \Omega_M H_0^2}{T(k) D(a) k^2}$$

## Implications:

- ▶ Unique  $1/k^2$  scaling of bias; no free parameters
- ▶ Distinct from effect of all other cosmo parameters
- ▶ Straightforwardly measured (g-g, g-T,...)
- ▶ Extensively tested with numerical simulations; good agreement found
- ▶ In general, LSS can probe:
 
$$\Delta b(\mathbf{k}) \propto \begin{cases} k^{-2} \text{ (local)} \\ k^{-1} \text{ (folded)} \\ k^0 \text{ (equilateral)} \\ k^{-\alpha} \text{ (generic); } 0 \leq \alpha \leq 3 \end{cases}$$

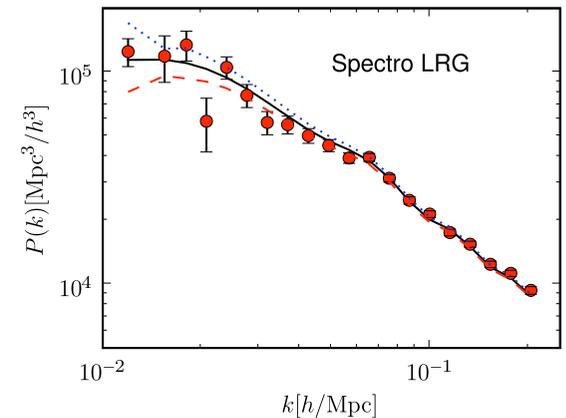
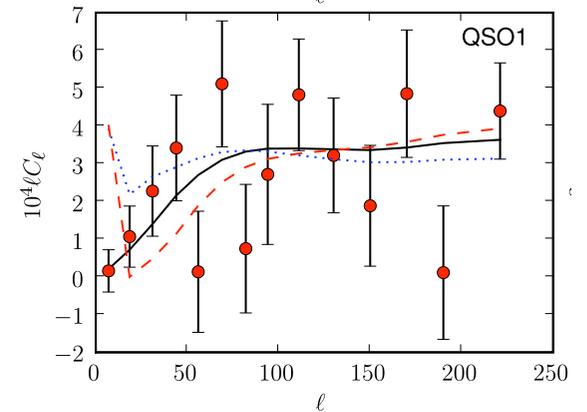
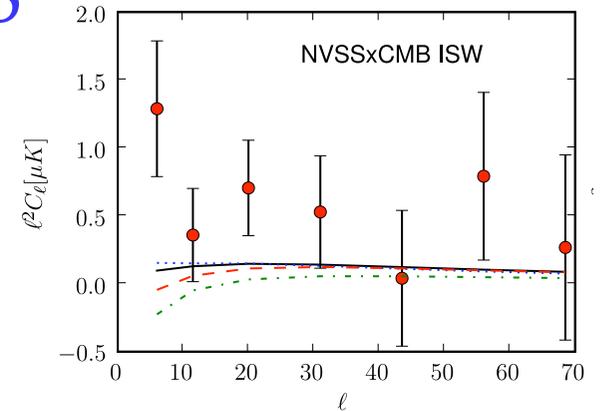
# Constraints from **current** data: SDSS



$f_{NL} = 8 \pm 30$  (68%, QSO)

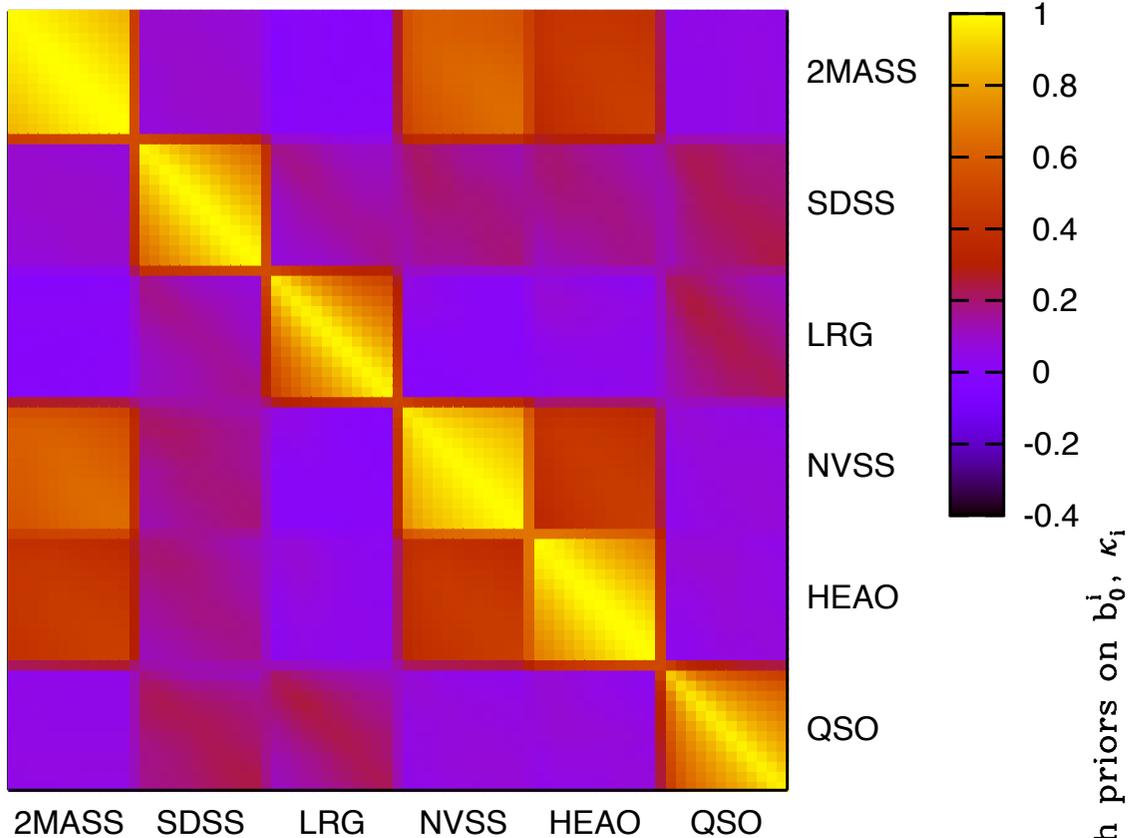
$f_{NL} = 23 \pm 23$  (68%, all)

Slosar et al. 2008  
(also Giannantonio et al 2013)



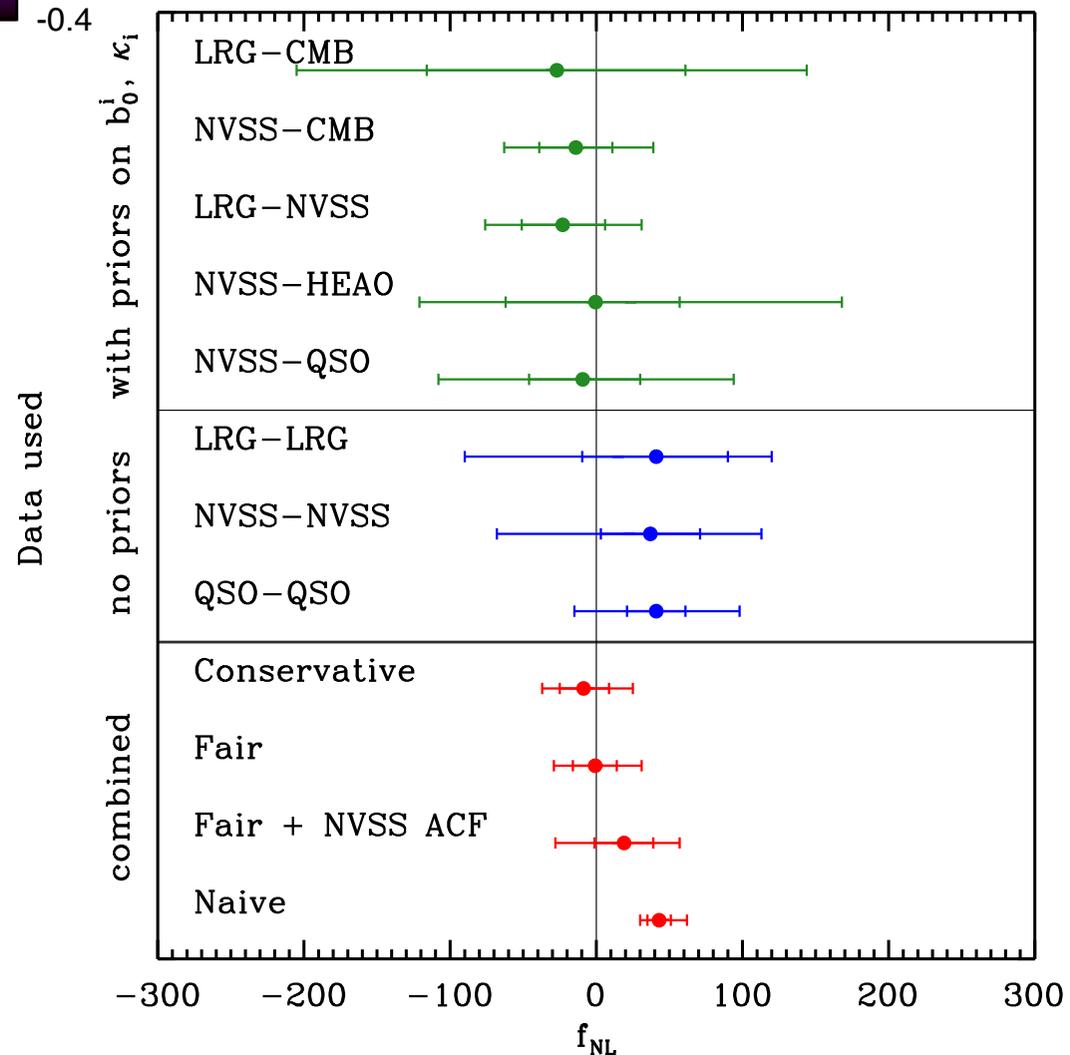
**Future** data forecasts for LSS:  $\sigma(f_{NL}) \approx \mathcal{O}(\text{few})$

(at least?) as good as, and highly complementary, to Planck CMB



Covariance matrix

# Final constraints:



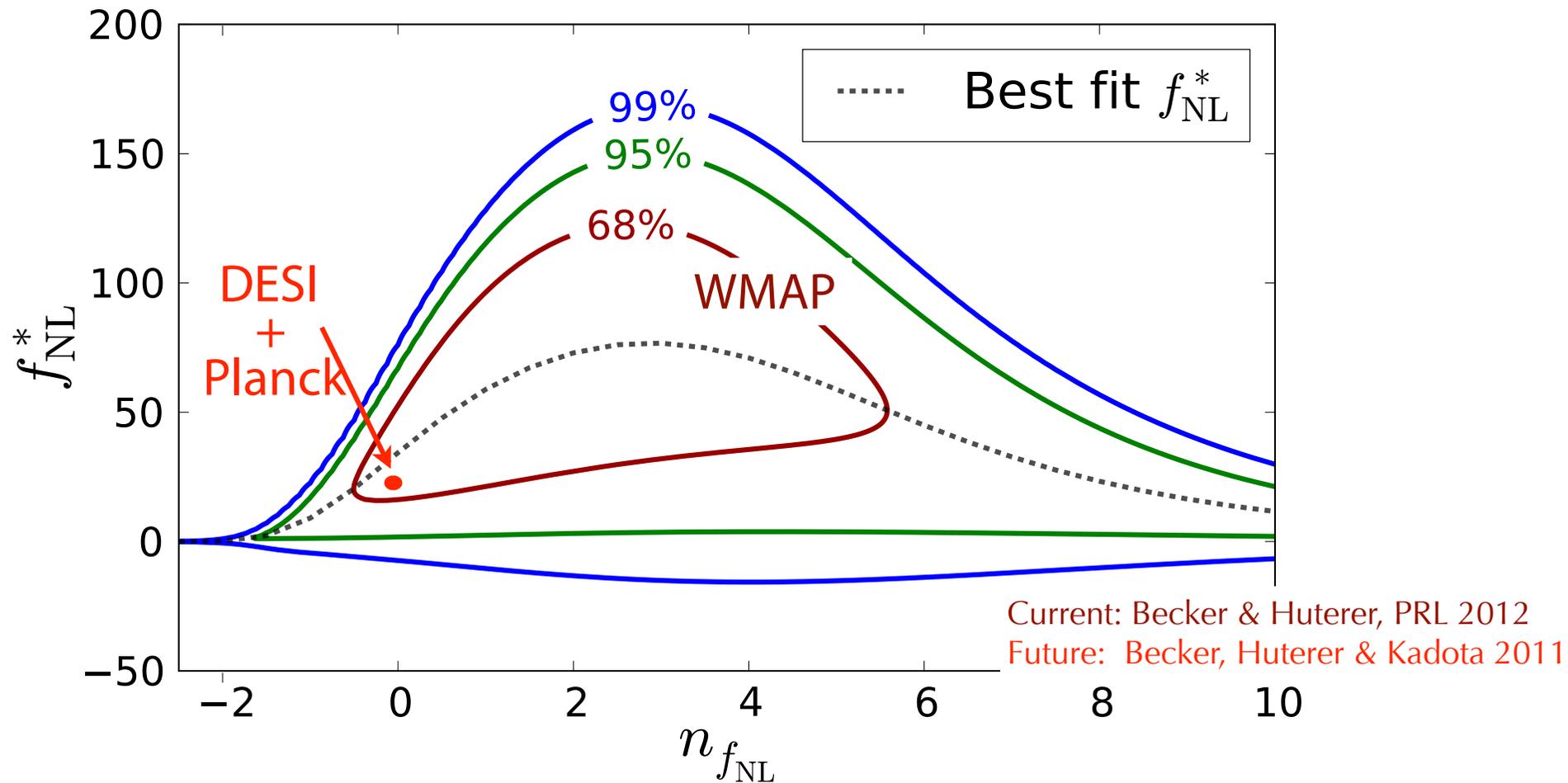
# Next Frontier: Large-Scale Structure

	CMB	LSS
dimension	2D	3D
# modes	$\propto l_{\max}^2$	$\propto k_{\max}^3$
systematics & selection func.	relatively clean	relatively messy
temporal evol.	no	yes
can slice in	$\lambda$ only	$\lambda, M, \text{bias} \dots$

**More general NG models:  
beyond  $f_{NL}$**

Current and future constraints on:

$$f_{\text{NL}}(k) = f_{\text{NL}}^* \left( \frac{k}{k_*} \right)^{n_{f_{\text{NL}}}}$$



Also: Halos of mass  $M$  probe NG on scale  $k \sim M^{-1/3}$

# Dark Energy Survey Instrument (DESI)



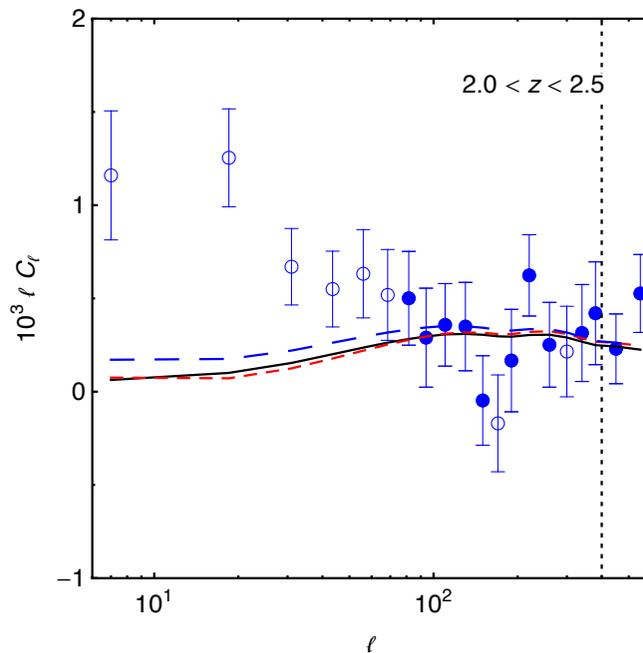
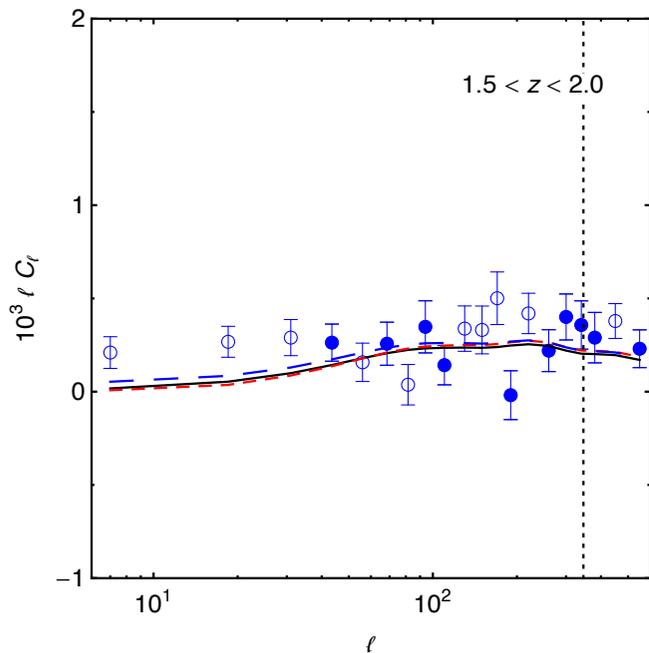
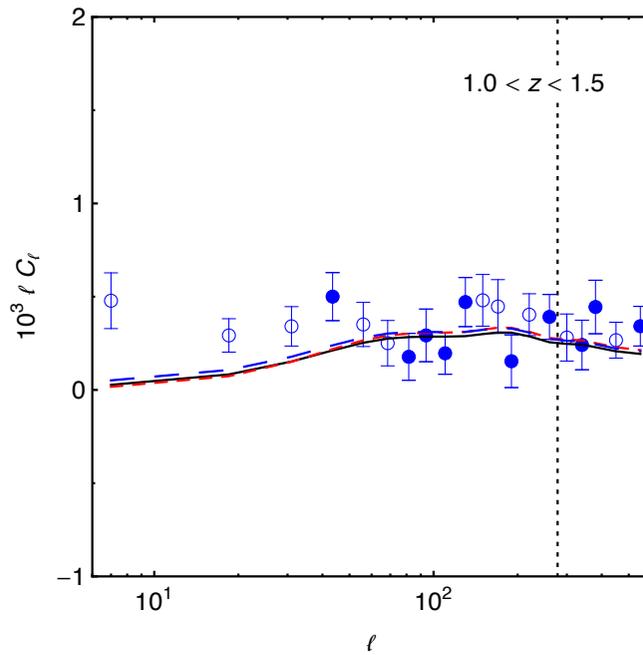
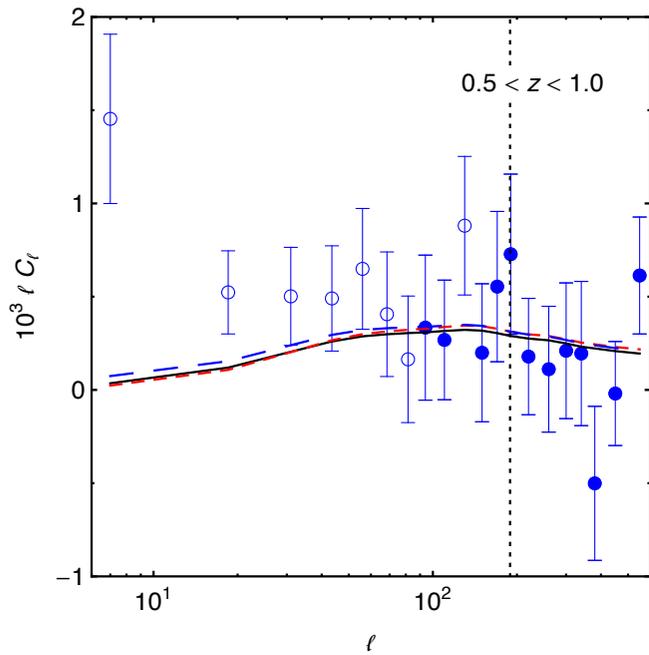
- Huge spectroscopic survey on Mayall telescope (Arizona)
- ~5000 fibres, ~15,000 sqdeg, ~20 million spectra
- LRG in  $0 < z < 1$ , ELG in  $0 < z < 1.5$ , QSO  $2.2 < z < 3.5$
- Great for dark energy (RSD, BAO)
- **Great for NG** - 3D  $P(k, z)$ , bispectrum...
- start 2018, funding DOE + institutions

**Systematic Errors:**  
**(photometric) calibration errors**

# For the NG measurements, photo-z but also: (photometric) calibration errors

- ▶ **Detector sensitivity:** sensitivity of the pixels on the camera vary along the focal plane. Sensitivity of a given pixel can change with time.
- ▶ **Observing conditions:** spatial and temporal variations.
- ▶ **Bright objects:** The light from foreground bright stars and galaxies affects the sky subtraction procedure, which impairs the surveys' completeness near bright objects.
- ▶ **Dust extinction:** Dust in the Milky Way absorbs light from the distant galaxies.
- ▶ **Star-galaxy separation:** In photometric surveys, faint stars can be erroneously included in the galaxy sample. Conversely, galaxies are sometimes misclassified as stars and culled from the sample. Remember, stars are *not* randomly distributed across the sky.
- ▶ **Deblending:** Galaxy images can overlap, and it can be difficult to cleanly separate photometric and spectroscopic measurements for the blended objects.

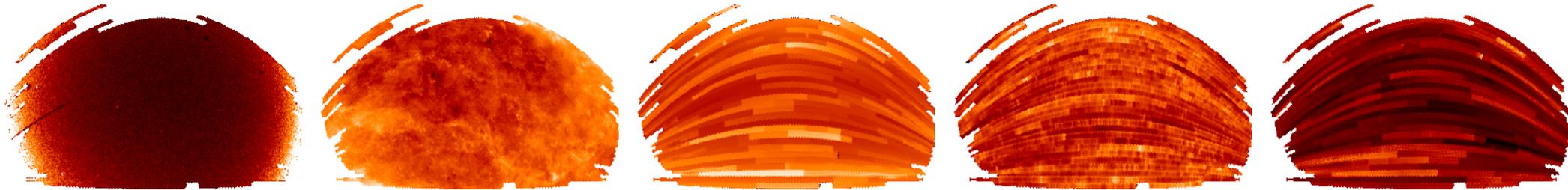
# Calibration errors in SDSS DR8 power spectra



QSO power spectra  
from SDSS;  
open circle points not  
used since they may  
be systematics-  
contaminated!

Similar results for LRGs  
(not shown)

# LSS calibration errors: example maps, power spectra



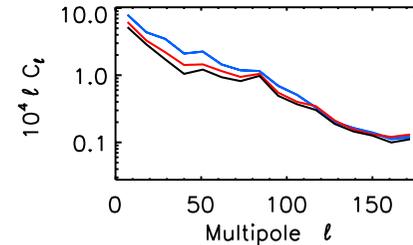
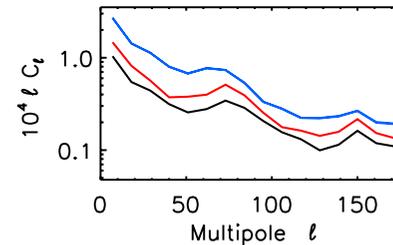
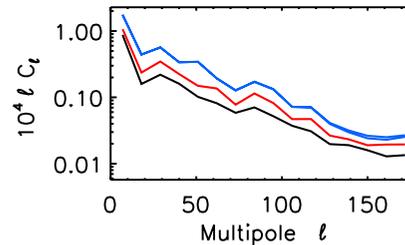
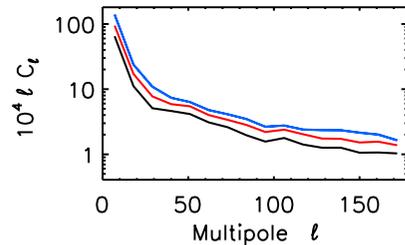
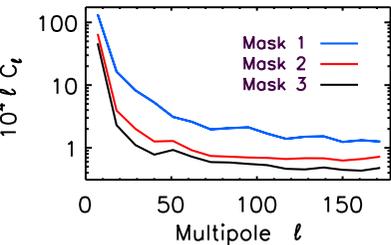
(a) Stellar density

(b) Extinction

(c) Airmass

(d) Seeing

(e) Sky brightness



Leistedt et al 2013

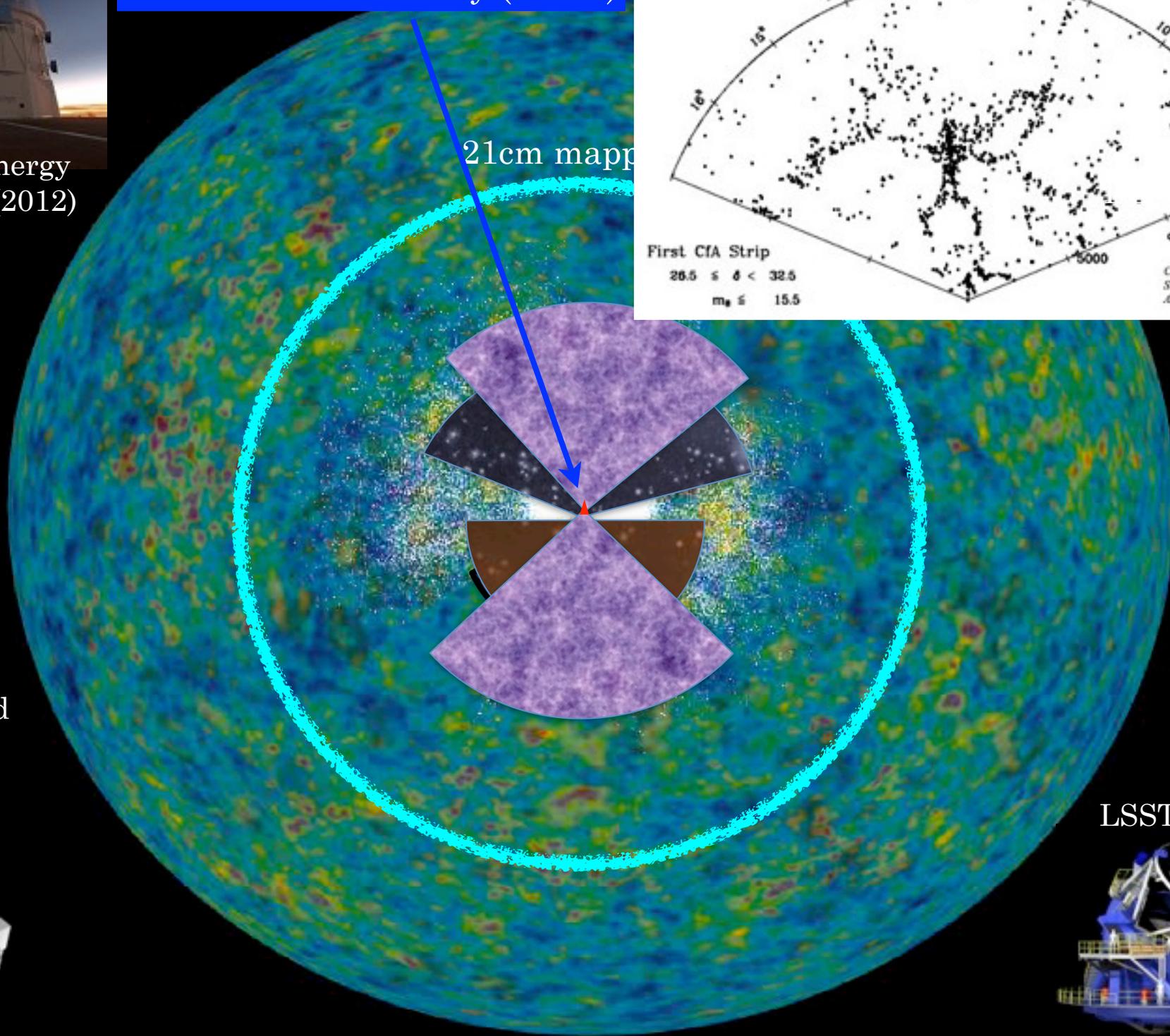
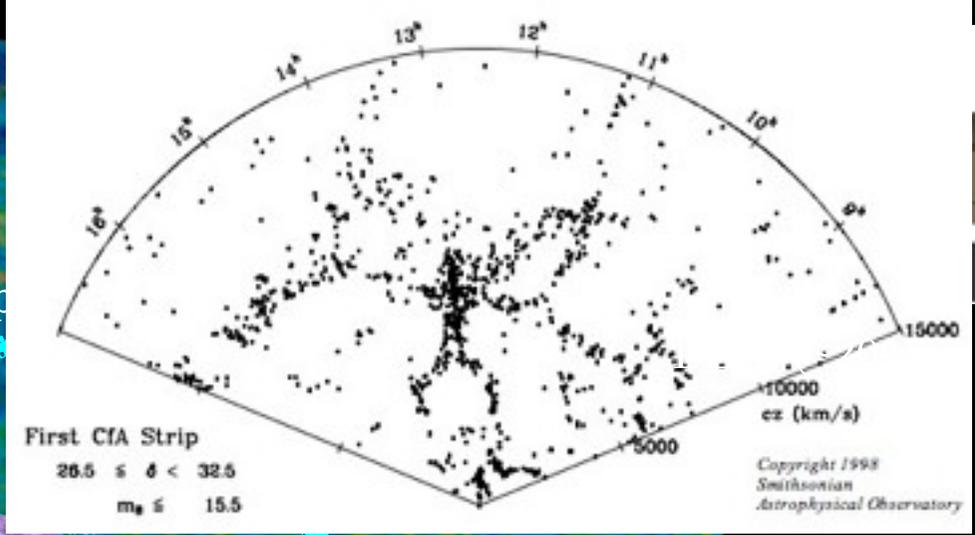
- dominate on large angular scales
- can be measured, removed using same or other data

Huterer, Cunha & Fang 2013;  
Shafer & Huterer 2015

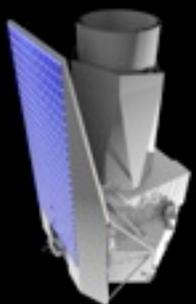
▲ Harvard-Cfa survey (1980s)



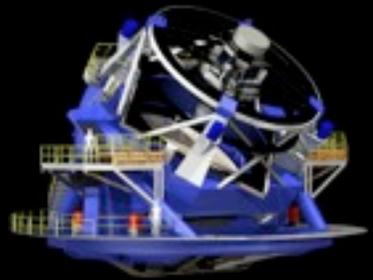
Dark Energy Survey (2012)



Euclid and WFIRST (~202X)



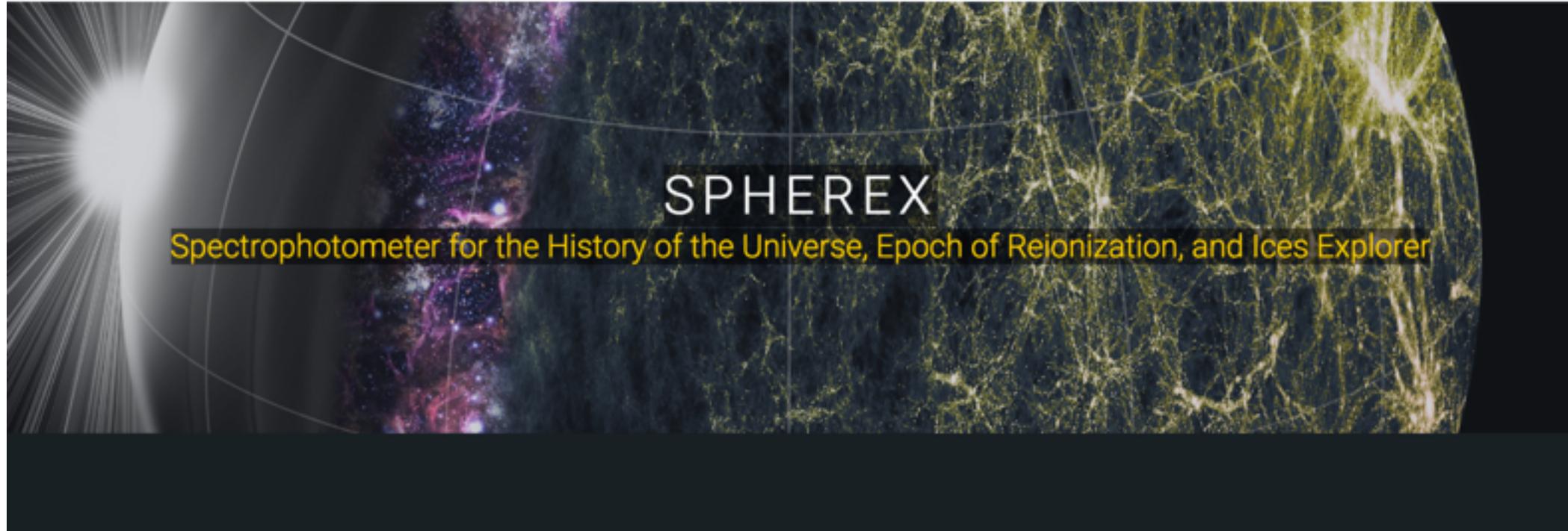
LSST (~2018)



# SPHEREx

proposal for telescope dedicated to measuring NG (and other science)

[Home](#) [Science](#) [Instrument](#) [Strategy](#) [Publications](#) [Team](#)



[spherex.caltech.edu](http://spherex.caltech.edu)

- 97 bands (!) with Linearly Variable Filters (LVF)
- $\lambda$  between 0.75 and 4  $\mu\text{m}$
- small (20cm) telescope, big field of view
- whole sky out to  $z \sim 1$
- **goal:  $\sigma(f_{\text{NL}}) \lesssim 1$**

paper: Doré, Bock et al, arXiv:1412.4872

# Conclusions:

- ▶ Primordial NG directly tests inflation:
  - ▶ How many fields
  - ▶ What interactions, couplings
- ▶ Constraints from WMAP, Planck are superb and consistent with zero NG
- ▶ Extremely good prospects for testing with galaxy surveys, at smaller scales than CMB

# Advances in Astronomy special issue on “Testing the Gaussianity and Statistical Isotropy of the Universe”

<http://www.hindawi.com/journals/aa/2010/si.gsiu/>

15 review articles (all also on arXiv)

## Testing the Gaussianity and Statistical Isotropy of the Universe

Guest Editors: Dragan Huterer, Eiichiro Komatsu, and Sarah Shandera

*Non-Gaussianity from Large-Scale Structure Surveys*, Licia Verde  
Volume 2010 (2010), Article ID 768675, 15 pages

*Non-Gaussianity and Statistical Anisotropy from Vector Field Populated Inflationary Models*, Emanuela Dimastrogiovanni, Nicola Bartolo, Sabino Matarrese, and Antonio Riotto  
Volume 2010 (2010), Article ID 752670, 21 pages

*Cosmic Strings and Their Induced Non-Gaussianities in the Cosmic Microwave Background*,

