Precision Time Measurement and a Novel Two-Photon Clock Scheme

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Abstract: We propose a novel two-photon excitation scheme for the ${}^{1}S_{0} \rightarrow {}^{3}P_{0}$ optical clock transition that works in the absence of hyperfine structure and applied magnetic fields. Experimental progress in Hg is discussed.

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The frontiers of precision time measurement in atomic clocks have been rapidly advancing in the last decade [1, 2, 3, 4, 5, 6]. Improvements in frequency standards advance metrology for everyday applications such as GPS and for tests of fundamental physics. Current optical-frequency atomic clocks based on the electric dipole E1 forbidden ${}^{1}S_{0} \rightarrow {}^{3}P_{0}$ transition require either hyperfine structure or an applied magnetic field to induce mixing of the ${}^{3}P_{0}$ and ${}^{3}P_{1}$ levels. This reduces the lifetime of the ${}^{3}P_{0}$ level and thus the fundamental resolution of the clock. We propose an excitation scheme utilizing a two-photon excitation coupling the ${}^{1}S_{0}$ ground state to the ${}^{3}P_{0}$ clock state through an allowed E1 and M1 excitation sequence. Utilizing a two-photon transition with counter-propagating laser beams eliminates the first-order Doppler-broadening and thus permits the use of a hot vapor cell of Hg atoms. This largely increases the number of atoms addressed compared to ion frequency standards [1, 2, 3, 4] and introduces simpler initial state preparation compared to neutral frequency standards that employ an optical lattice [5, 6].

The transition rate using this two-photon excitation scheme is given by:

$$R_{2\gamma} = N \frac{4I^2}{\hbar^4 c^4 \epsilon_0^2} \frac{\langle {}^3P_0 | \mu | {}^3P_1 \rangle_{M1}^2 \langle {}^3P_1 | d | {}^1S_0 \rangle_{E1}^2}{\Gamma \Delta^2} \approx N \frac{I^2}{\Gamma} \times 10^{-10}$$
(1)

where N is the number of atoms sampled, I is the excitation laser intensity, $\langle {}^{3}P_{0}|\mu|{}^{3}P_{1}\rangle_{M1}$ is the magnetic dipole matrix element, $\langle {}^{3}P_{1}|d|{}^{1}S_{0}\rangle_{E1}$ is the electric dipole matrix element, Γ is the experimentally broadened linewidth of the transition, and Δ is the detuning of an excitation photon's energy from the intermediate ${}^{3}P_{1}$ level. It is experimentally feasible to achieve $I > 10^{10} \left[\frac{W}{m^2}\right]$, and $N > 10^8 \left[\#\right]$. In the first phase of this experiment we anticipate $\Gamma \ll 10^8 \left[Hz\right]$ and expect to detect transitions at a rate of a few Hertz.

To drive the ${}^{1}S_{0} \rightarrow {}^{3}P_{0}$ transition in neutral Hg we have developed a high power, narrow linewidth laser system at 531 nm. A fiber amplified 1062nm laser has been doubled using a PPMgO:SLT crystal to produce CW powers of 5.5 W. Population in the metastable ${}^{3}P_{0}$ level can be detected using an NH₃ buffer gas which induces ${}^{3}P_{0}$ relaxation [7]. The presence of NH₃ broadens the transition linewidth but serves as an expedient tool to determine clock state population and demonstrate the viability of this clock scheme. An alternate, non-broadening detection scheme involves sampling the E1 allowed ${}^{3}P_{0} \rightarrow {}^{3}S_{1}$ transition.

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